

# Diffractive DIS Cross Sections and PDFs from Rapidity Gap and Leading Proton Data



Frank-Peter Schilling (CERN and DESY)

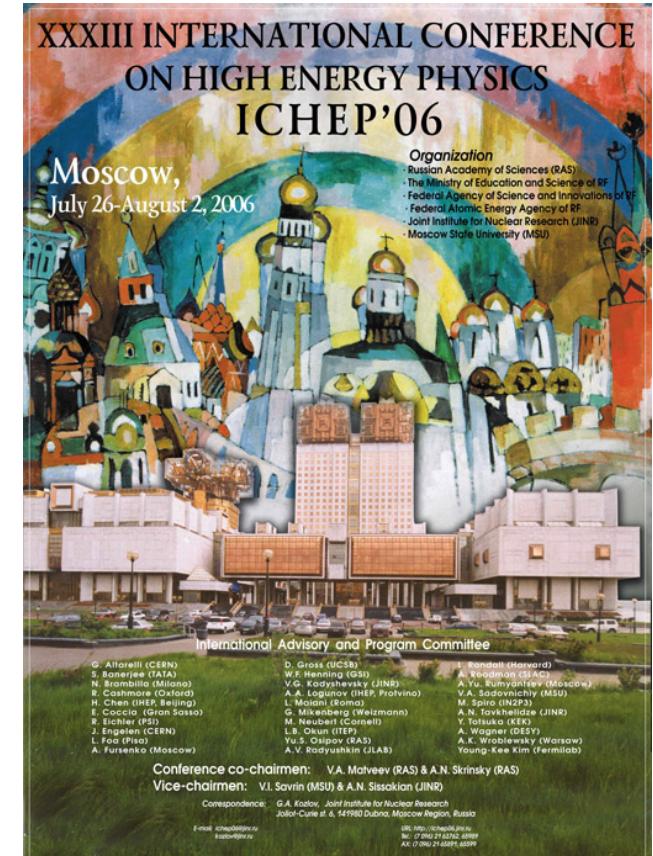


On behalf of the H1 Collaboration

ICHEP 2006  
Moscow, July 2006

Results presented (mainly) from two recent papers:

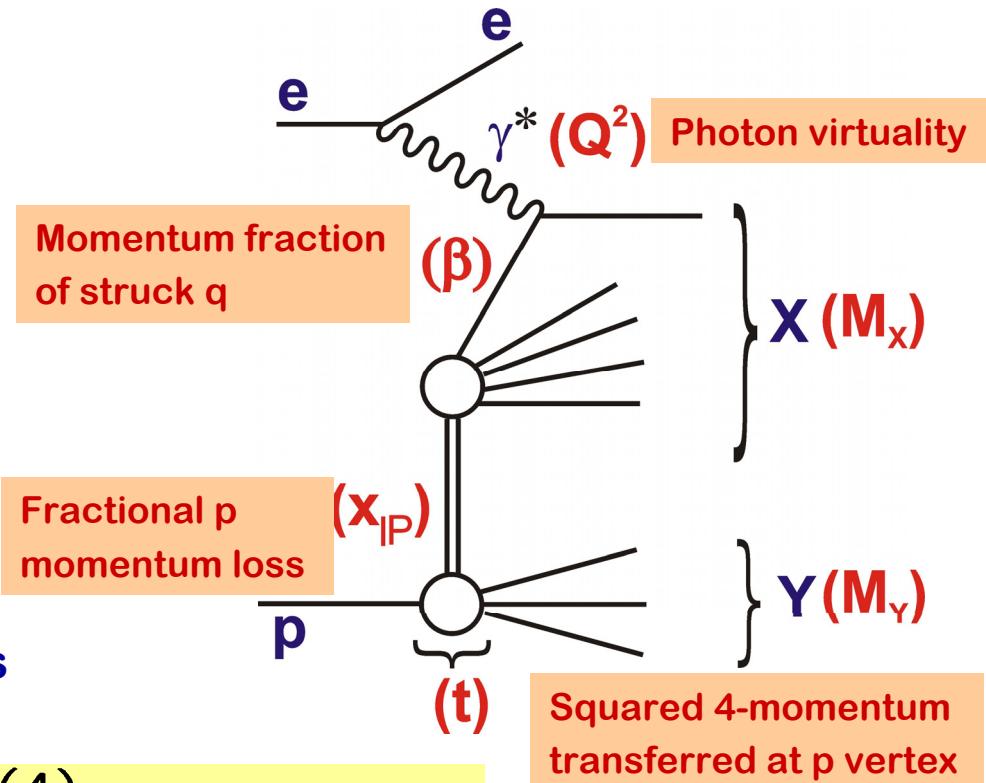
- DESY06-048, hep-ex/0606003 subm to EPJC
- DESY06-049, hep-ex/0606004 subm to EPJC



# Overview and Kinematics

**Diffractive DIS at HERA: Deep Inelastic Scattering where the proton stays intact and loses just small momentum fraction ...**

- Comparison between measurement methods:
  - leading proton (Roman Pots)
  - Large rapidity gap
- Dependences on  $t$  and  $x_{IP}$
- $Q^2$  and  $\beta$  dependences and NLO DGLAP QCD fits
- Comparison between diffractive and inclusive DIS
- Main observable: Reduced cross section  $\sigma_r^{D(3,4)}$



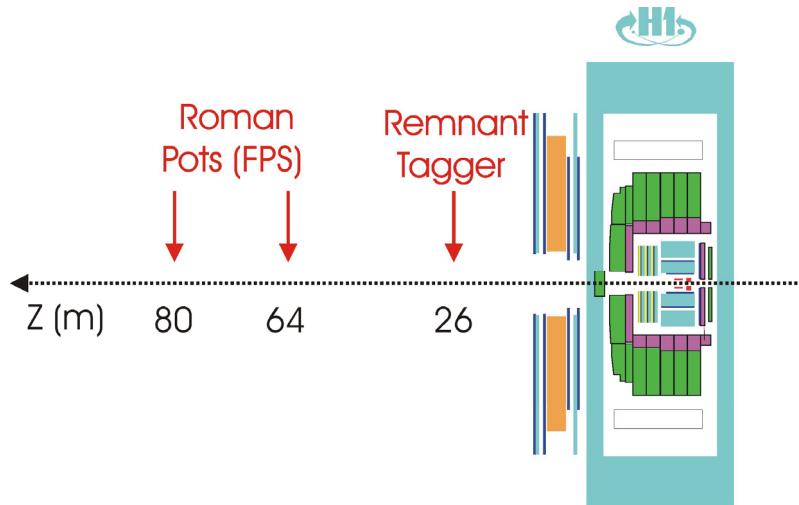
$$\frac{d^4 \sigma_{ep \rightarrow eXp}}{dx dQ^2 dx_{IP} dt} = \frac{4\pi\alpha^2}{xQ^4} Y_+ \sigma_r^{D(4)}(x, Q^2, x_{IP}, t)$$

$$\sigma_r^{D(4)}(x, Q^2, x_{IP}, t) = F_2^{D(4)} - \frac{y^2}{Y_+} F_L^{D(4)} \approx F_2^{D(4)}$$

# Event Selection Methods

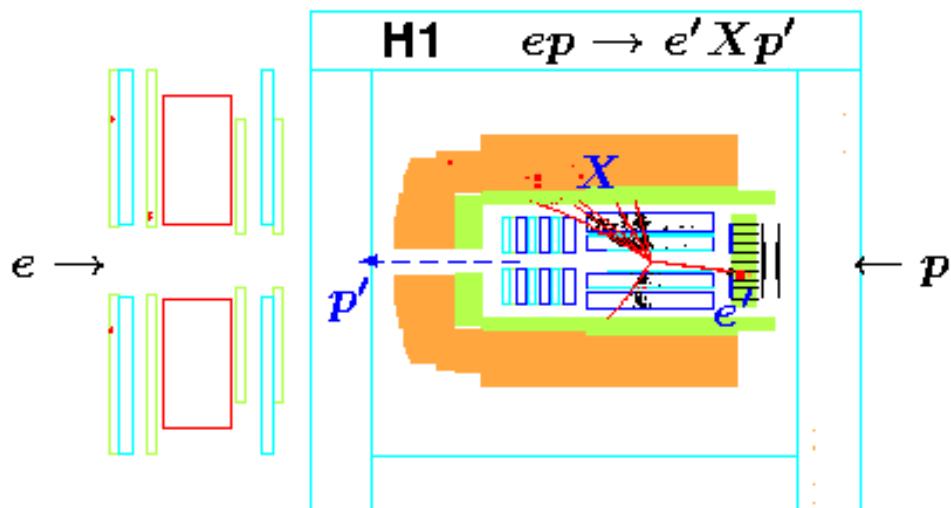
## 1. Tag and measure final state proton in Forward Proton Spectrometer (FPS method)

- No proton dissociation
- Can measure  $t$
- Acceptance at high  $x_{IP}$
- ... but low Pot acceptance



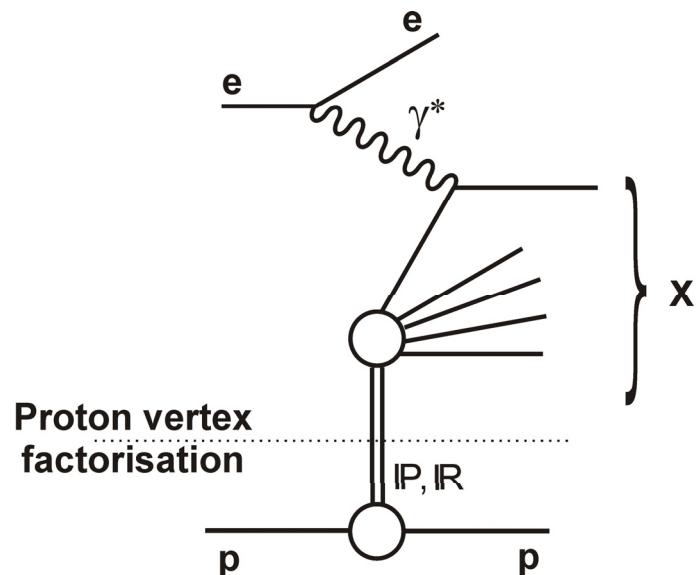
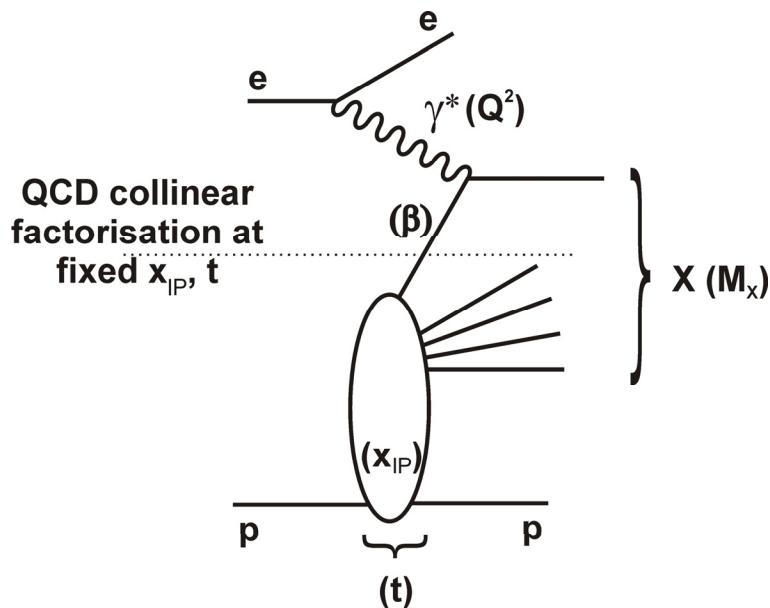
## 2. Require Large Rapidity Gap spanning at least $3.3 < \eta < 7.5$ and measure hadrons in central detector (LRG method)

- Some proton dissociation
- Correct to  $M_Y < 1.6$  GeV
- Near-perfect acceptance at low  $x_{IP}$



# Two Levels of Factorization

- QCD hard scattering collinear factorization (Collins) at fixed  $x_{IP}$  and  $t$ 
  - After integration over measured  $M_Y, t$  ranges
- “Proton vertex” factorization of  $x, Q^2$  from  $x_{IP}, t$  (and  $M_Y$ ) dependences
  - Separately for leading IP and sub-leading IR exchanges

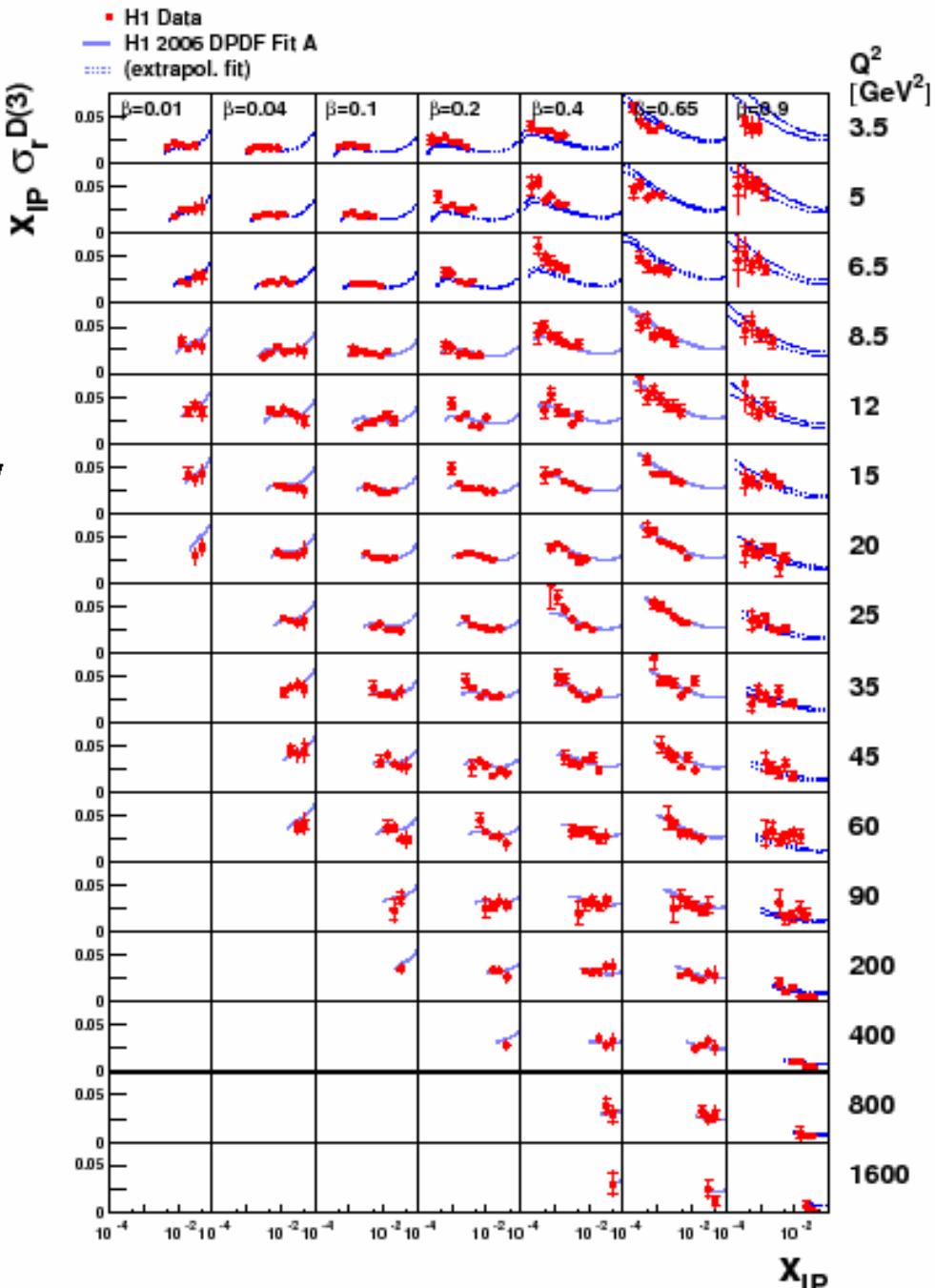
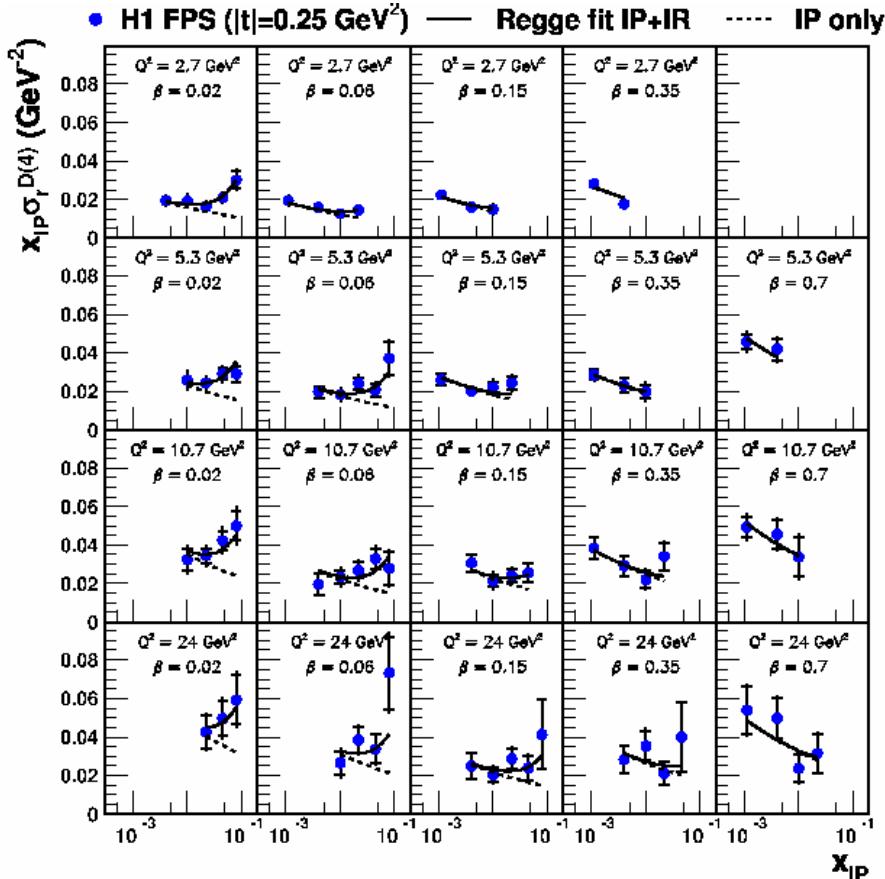


$$d\sigma_i(ep \rightarrow eXY) = f_i^D(x, Q^2, x_{IP}, t) \otimes d\hat{\sigma}^{ei}(x, Q^2)$$

$$f_i^D(x, Q^2, x_{IP}, t) = f_{IP/p}(x_{IP}, t) \times f_i^{IP}(\beta = x/x_{IP}, Q^2)$$

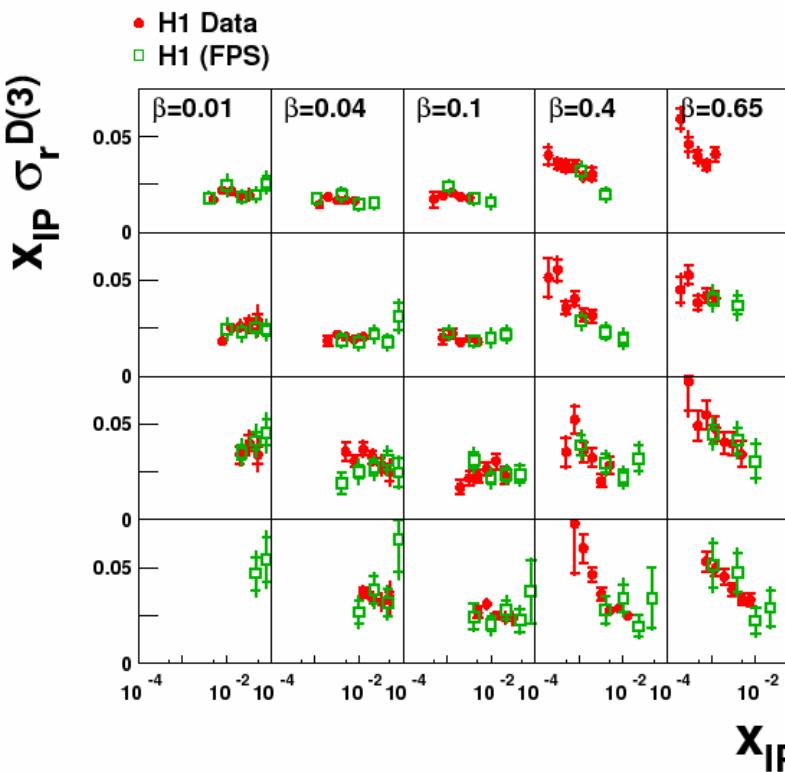
# Data Overview

- LRG:  $M_Y < 1.6 \text{ GeV}$
- $3.5 < Q^2 < 1600 \text{ GeV}^2$
- FPS:  $Y=p$
- $2.7 < Q^2 < 24 \text{ GeV}^2$

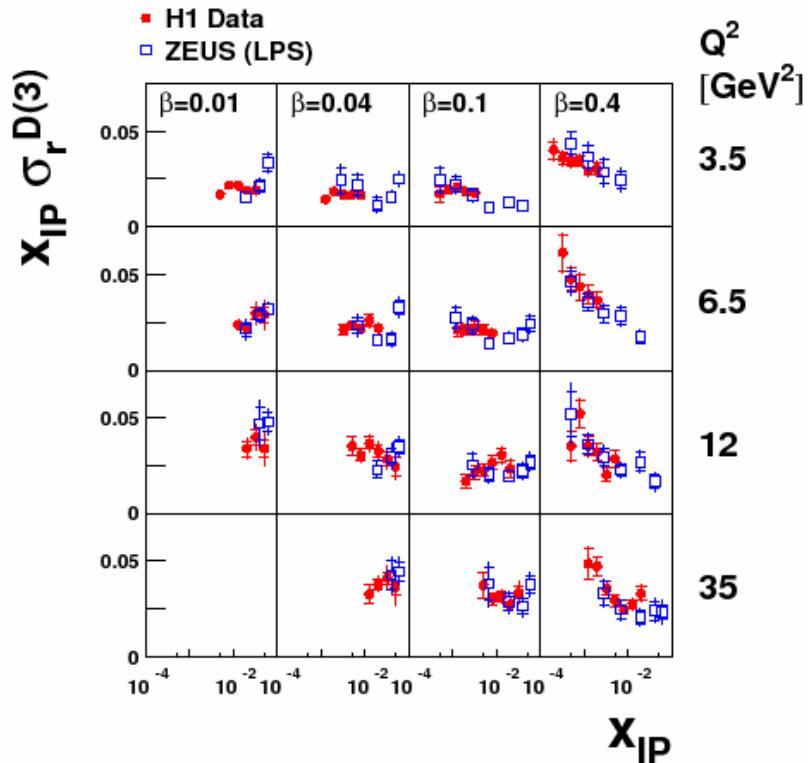


# Comparison of Rapidity Gap with Leading Proton (H1,ZEUS) data

- **LRG vs FPS data:**



- **LRG vs ZEUS LPS data:**

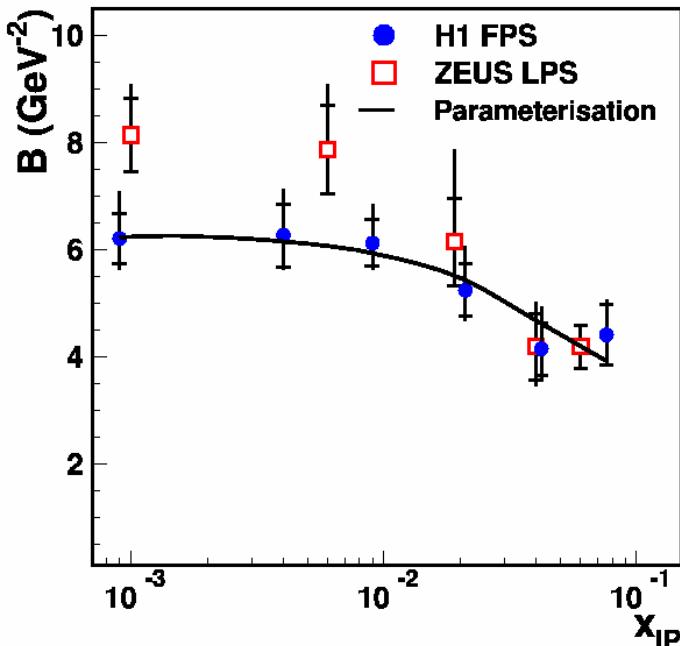


- **Agreement in detail between FPS and LRG methods: Ratio LRG/FPS independent of kinematics within errors:**  

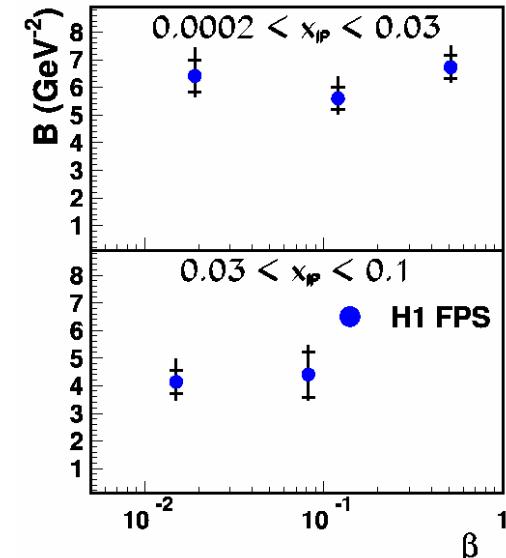
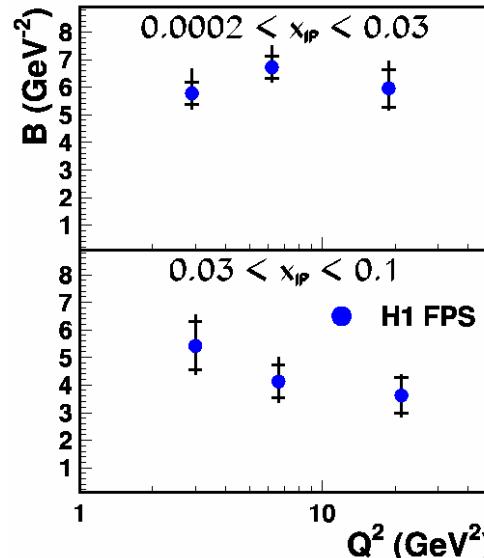
$$\frac{\sigma(M_Y < 1.6 \text{ GeV})}{\sigma(Y=p)} = 1.23 \pm 0.03(\text{stat.}) \pm 0.16(\text{syst.})$$
- **ZEUS-LPS and H1-FPS normalizations agree to 8%**
- **Verry good agreement between proton tagging and LRG methods if p dissociation is accounted for!**

# t dependence from FPS measurements

- $B(x_{IP})$  from fit  $d\sigma/dt \sim \exp(B|t|)$ :



- Independent of  $\beta$ ,  $Q^2$  within errors



- $B(x_{IP})$  data constrain IP,IR flux factors in p vertex fact. Model

- Regge motivated form:

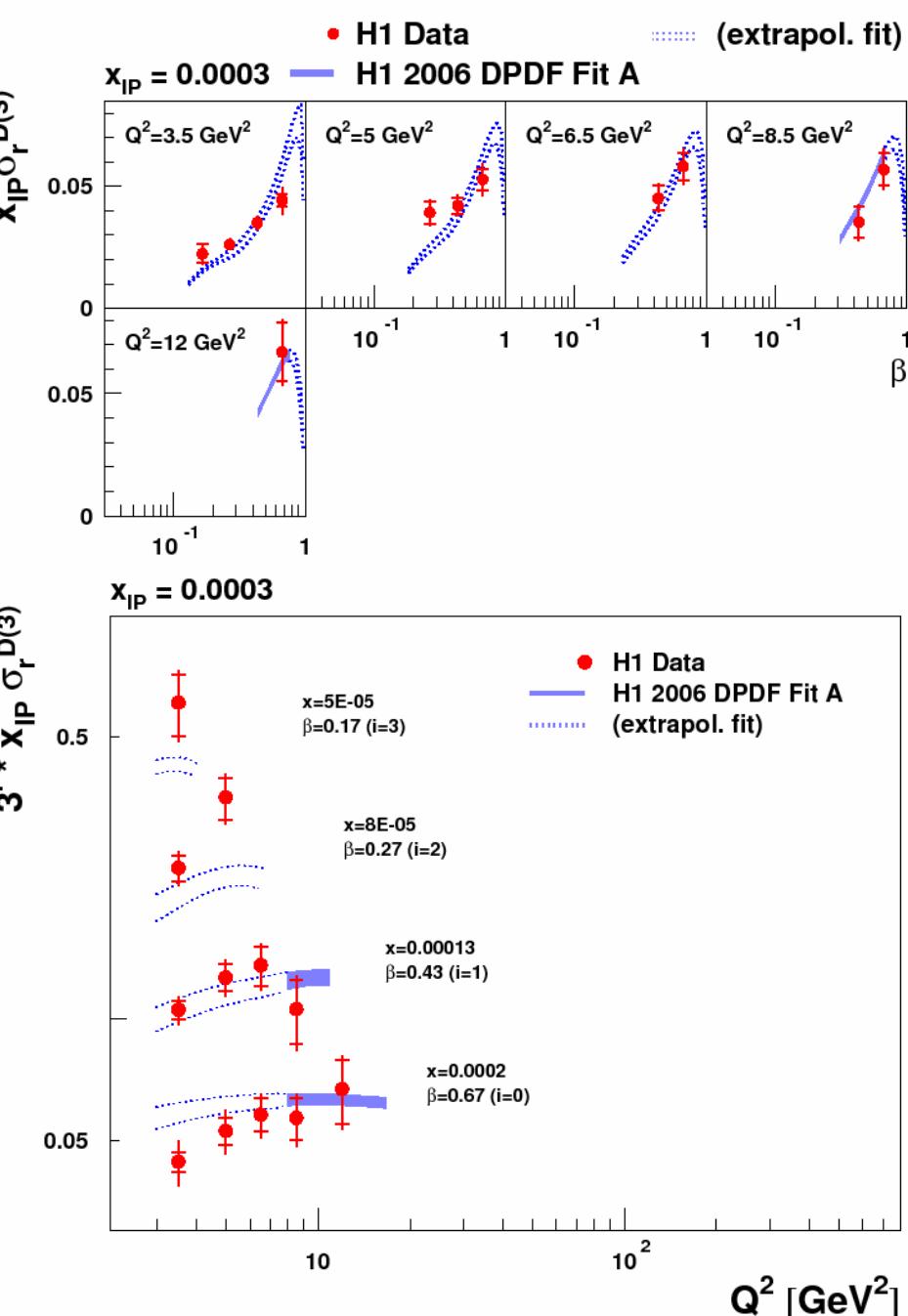
$$f_{IP/p}(x_{IP}, t) = \frac{e^{B_{IP}t}}{x_{IP}^{2\alpha_{IP}(t)-1}} \quad \alpha_{IP}(t) = \alpha_{IP}(0) + \alpha'_{IP}t$$

- Fitting low  $x_{IP}$  data to yields

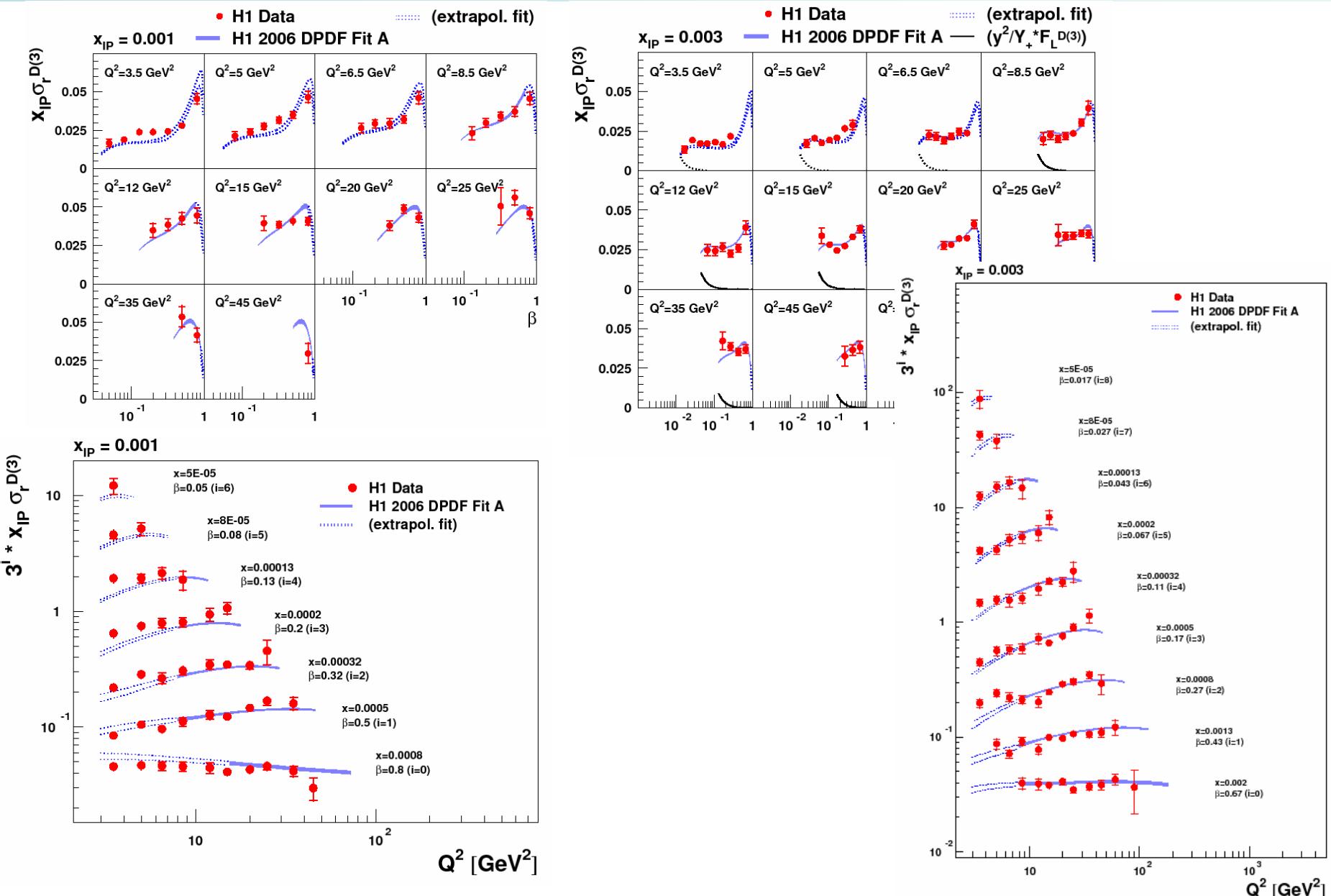
$$B = B_{IP} + 2\alpha'_{IP} \ln(1/x_{IP})$$

$$B_{IP} = 5.5_{-0.7}^{+2.0} \text{ GeV}^{-2} \quad \alpha'_{IP} = 0.06_{-0.06}^{+0.19} \text{ GeV}^{-2}$$

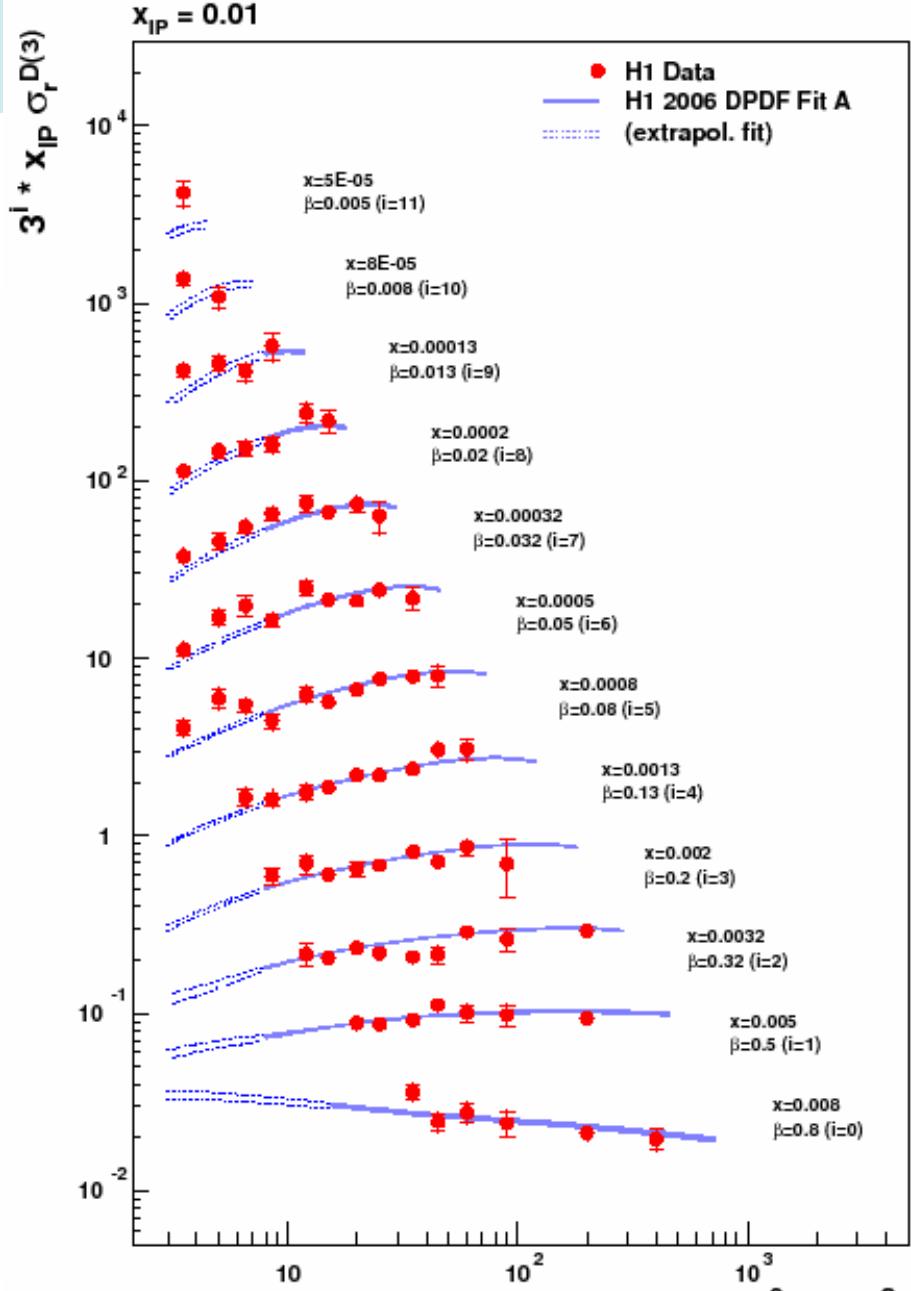
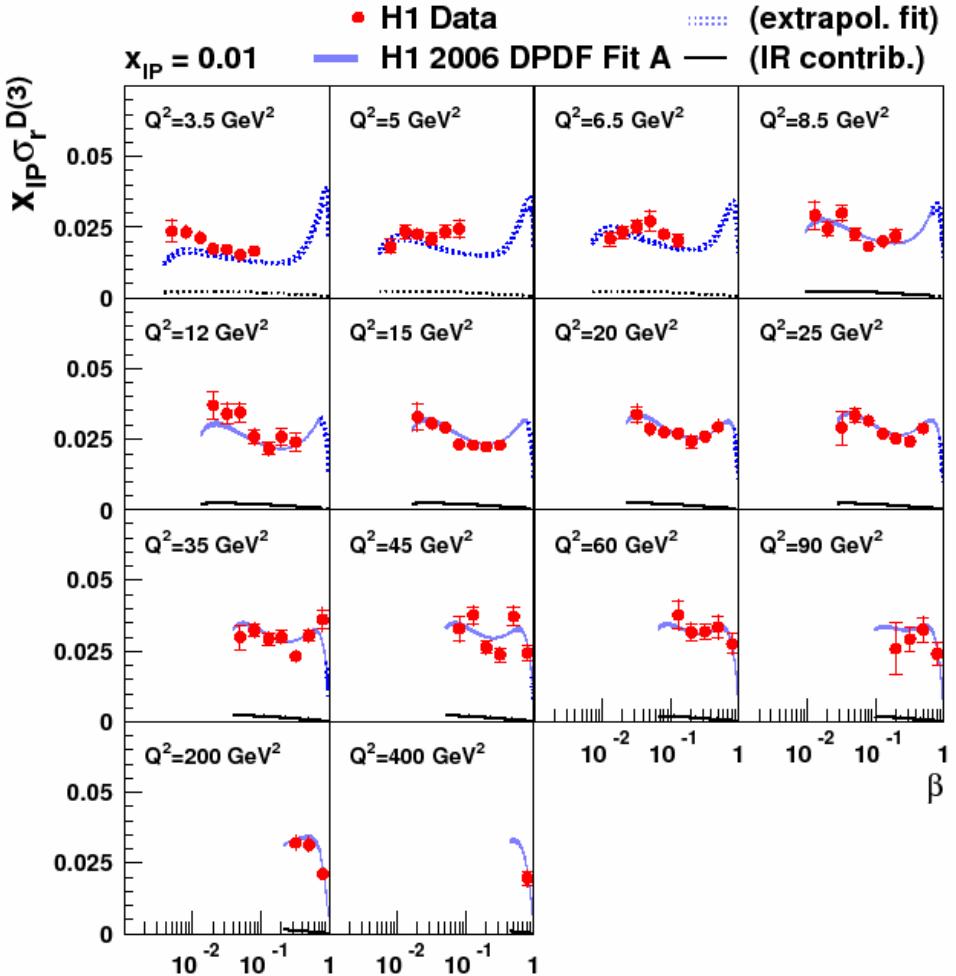
- Principal binning scheme for LRG data
- Study  $Q^2$  and  $x (= \beta \cdot x_{IP})$  dependences in detail at small number of fixed  $x_{IP}$  values
- Good precision – in best regions 5% (stat.), 5% (syst.), 6% (norm.)
- Data compared with “H1 2006 DPDF fit” and its error band (assumes p vtx factorization, see later)



# $\sigma_r^{D(3)}(\beta, Q^2, x_{IP})$ at $x_{IP}=0.001$ and $0.003$

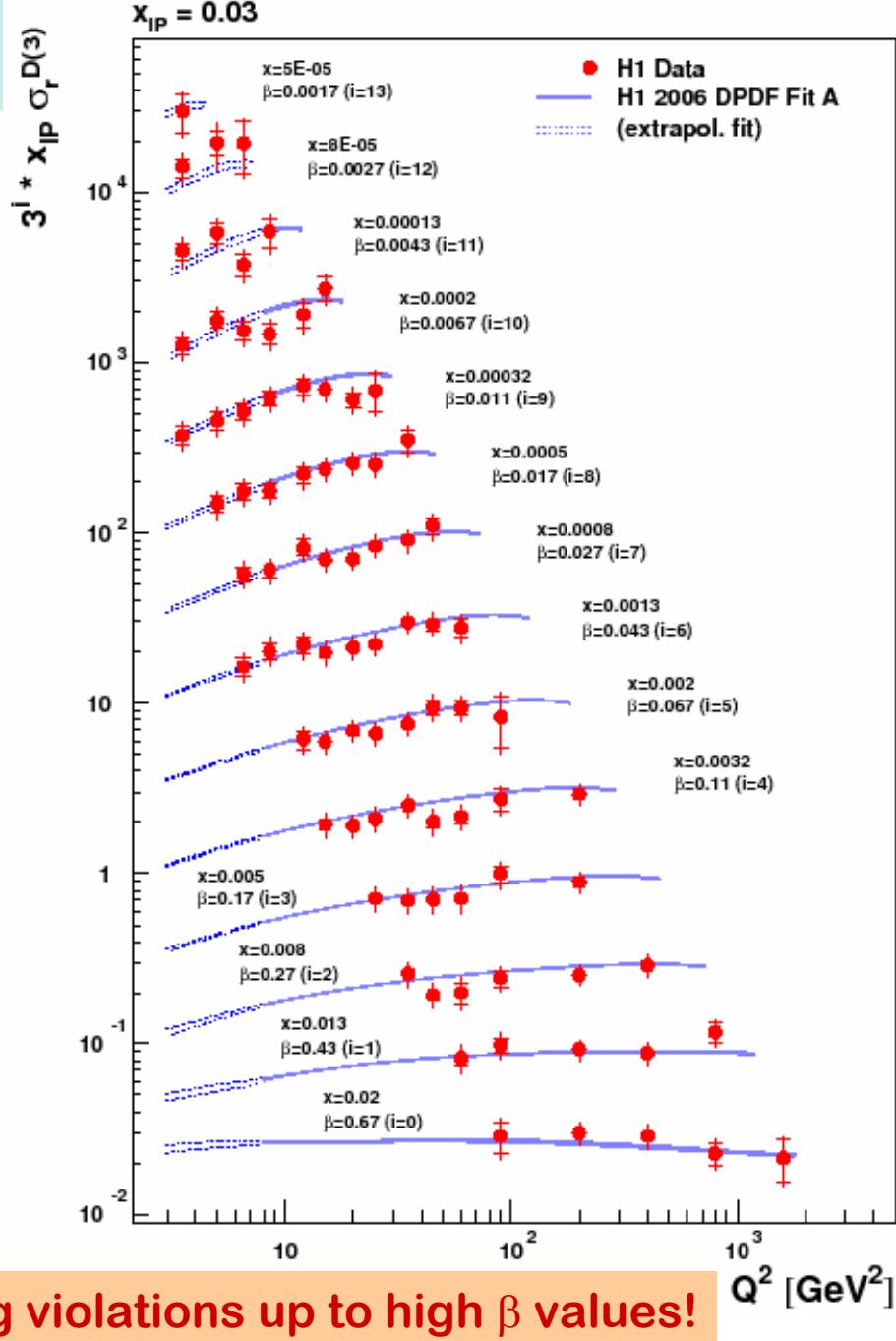
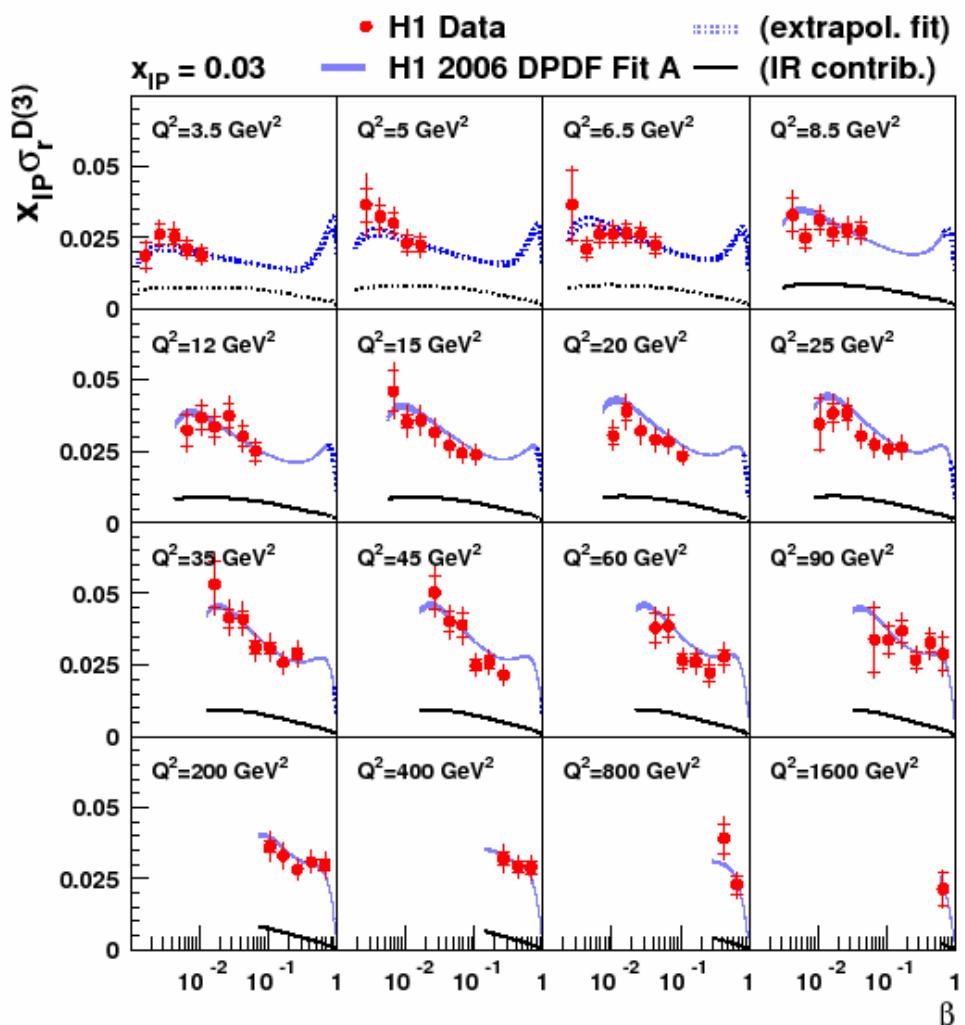


$\sigma_r^{D(3)}(\beta, Q^2, x_{IP})$  at  $x_{IP}=0.01$



Large positive  $Q^2$  scaling violations up to high  $\beta$  values!

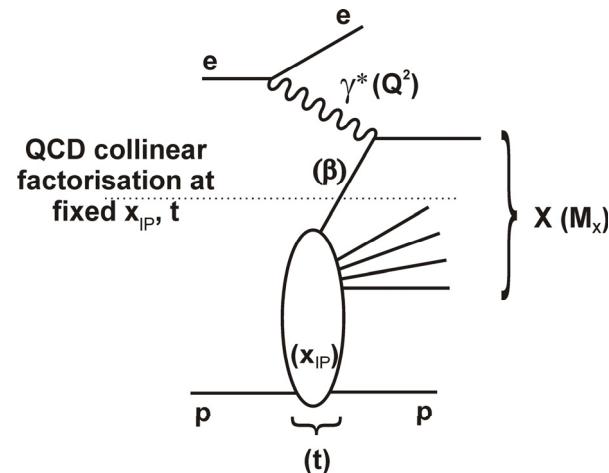
$\sigma_r^{D(3)}(\beta, Q^2, x_{IP})$  at  $x_{IP}=0.03$



Large positive  $Q^2$  scaling violations up to high  $\beta$  values!

# H1 2006 DPDF fit: Overview

- Fit LRG data from fixed  $x_{IP}$  binning, using NLO DGLAP evolution of PDFs (massive scheme) to describe  $x, Q^2$  dependences
- Proton vtx factorisation framework (supported by data)
  - relate data from different  $x_{IP}$  values with complementary  $x, Q^2$  coverage
- For IP exchange, free parameters are
  - $\alpha_{IP}(0)$  (describes  $x_{IP}$  dependence)
  - PDF parameters at evolution starting scale  $Q_0^2$



$$f_{IP/p}(x_{IP}, t) = \frac{e^{B_{IP}t}}{x_{IP}^{2\alpha_{IP}(t)-1}}$$

- For sub-leading IR
  - all flux parameters taken from previous data
  - PDFs taken from Owens- $\pi$ ; single free param for normalization

# Kinematic range and DPDF parameterization

- To ensure data fitted are compatible with chosen framework, test sensitivity of fit results to variations of kinematic boundaries
  - Results stable for most variations ( $\beta_{\max}$ ,  $\beta_{\min}$ ,  $M_{X,\min}$ ,  $x_{IP,\max}$ )
  - Systematic variation of gluon density with minimum  $Q^2$  of data included in fit for  $Q^2 < 8.5 \text{ GeV}^2$ ; stable for larger  $Q^2\text{-min}$
- Fit all LRG data with  $Q^2 \geq 8.5 \text{ GeV}^2$ ,  $M_x > 2 \text{ GeV}$ ,  $\beta \leq 0.8$
- Parameterize
  - quark singlet  $z\Sigma(z, Q_0^2)$
  - gluon  $zg(z, Q_0^2)$  density

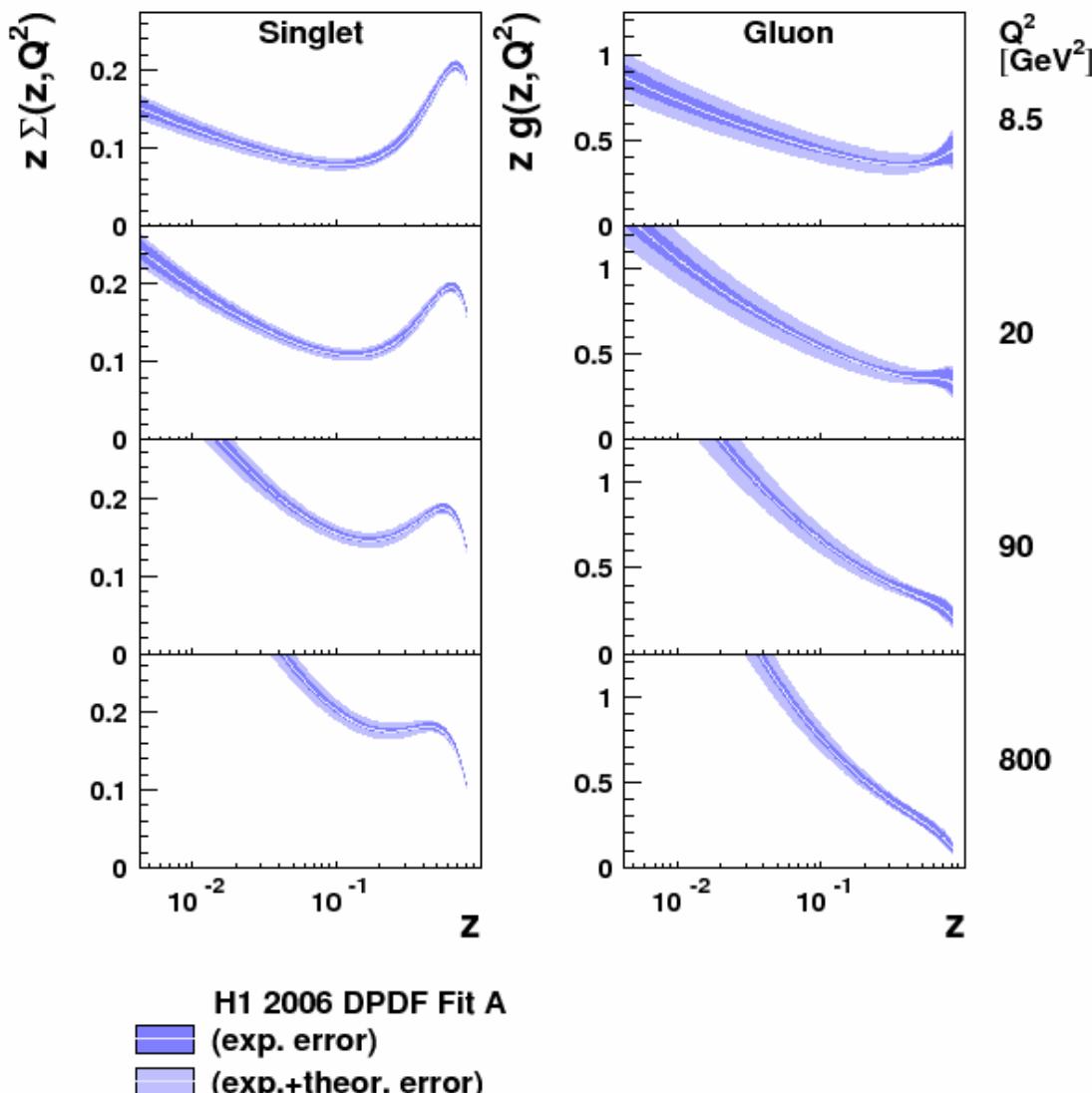
$$z\Sigma(z, Q_0^2) = A_q z^{B_q} (1 - z)^{C_q}$$

$$zg(z, Q_0^2) = A_g (1 - z)^{C_g}$$

- Gluon insensitive to  $B_g$
- Small number of parameters
  - need to optimize  $Q_0^2$  wrt  $\chi^2$
- Using world average value for  $\alpha_s(M_z) = 0.118$
- Results reproducible with Chebychev polynomials

# H1 2006 DPDF fit results (log z scale)

- $Q_0^2 = 1.75 \text{ GeV}^2$
- $\chi^2 \sim 158 / 183 \text{ dof}$
- Experimental uncertainty obtained by propagating errors on data (c.f. incl. fits,  $\Delta\chi^2=1$ )
- Theoretical uncertainty from varying fixed parameters of fit (flux params,  $\alpha_s$ ,  $m_c$ ,  $m_b$  etc.) and  $Q_0^2$  ( $\Delta\chi^2=1$ )
- Singlet constrained to ~5%, gluon to ~15% at low  $z$ ; error blowing up at highest  $z$



# A closer look at high z region

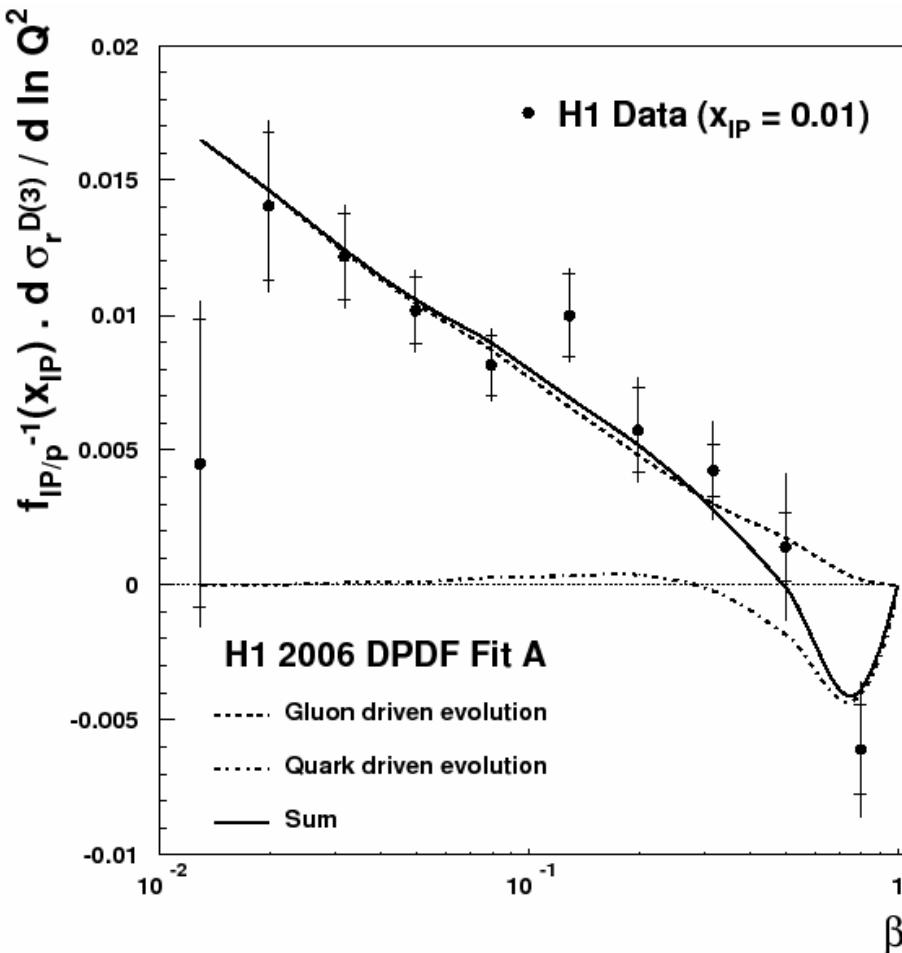
- As there are only singlet quarks, the evolution eq. for  $F_2^D$  is

$$\frac{dF_2^D}{d \ln Q^2} \sim \frac{\alpha_s}{2\pi} [P_{qg} \otimes g + P_{qq} \otimes \Sigma]$$



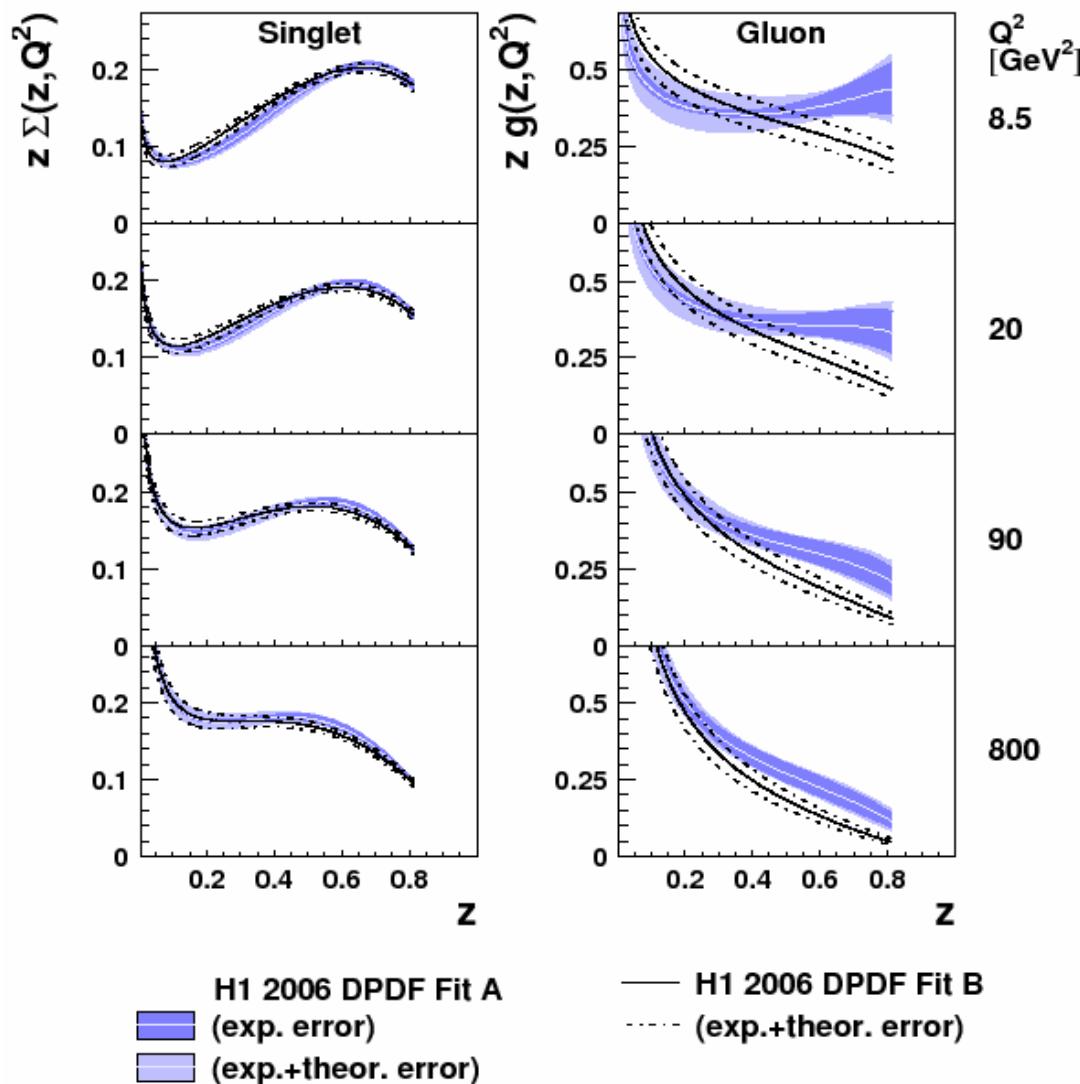
- At low  $\beta$ , evolution driven by  $g \rightarrow qq$ :
  - strong sensitivity to gluon
- At high  $\beta$  relative error on derivative grows,  $q \rightarrow qg$  contribution becomes important:
  - sensitivity to gluon is lost

Log. Derivative wrt  $Q^2$ :



# H1 2006 DPDF fit results (lin. z scale)

- Lack of sensitivity to high  $z$  gluon confirmed by dropping  $C_g$  parameter, so gluon is simple constant at  $Q_0^2$ :
- Fit B,  $\chi^2 \sim 164/184$  dof
  - Singlet very stable
  - Gluon similar at low  $z$
  - Substantial change to gluon at high  $z$



# Effective Pomeron Intercept

- From QCD fit to LRG data

$$\alpha_{IP}(0) = 1.118 \pm 0.008(\text{exp.})^{+0.029}_{-0.010}(\text{th.})$$

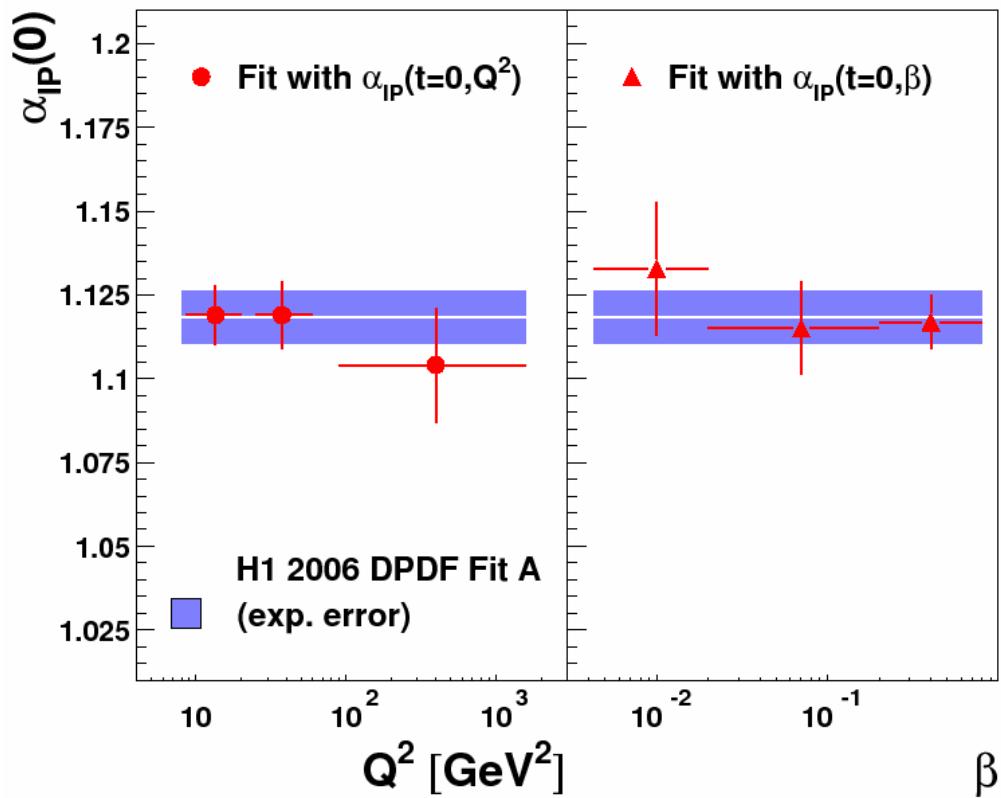
- Dominant uncertainty from strong correlation with  $\alpha'_{IP}$ : taking  $\alpha'_{IP} = 0.25$  instead of  $0.06 \text{ GeV}^{-2}$  yields  $\alpha_{IP}(0) \sim 1.15$

- No significant variation in  $Q^2$  or  $\beta$  (consistent with p vtx factorization)

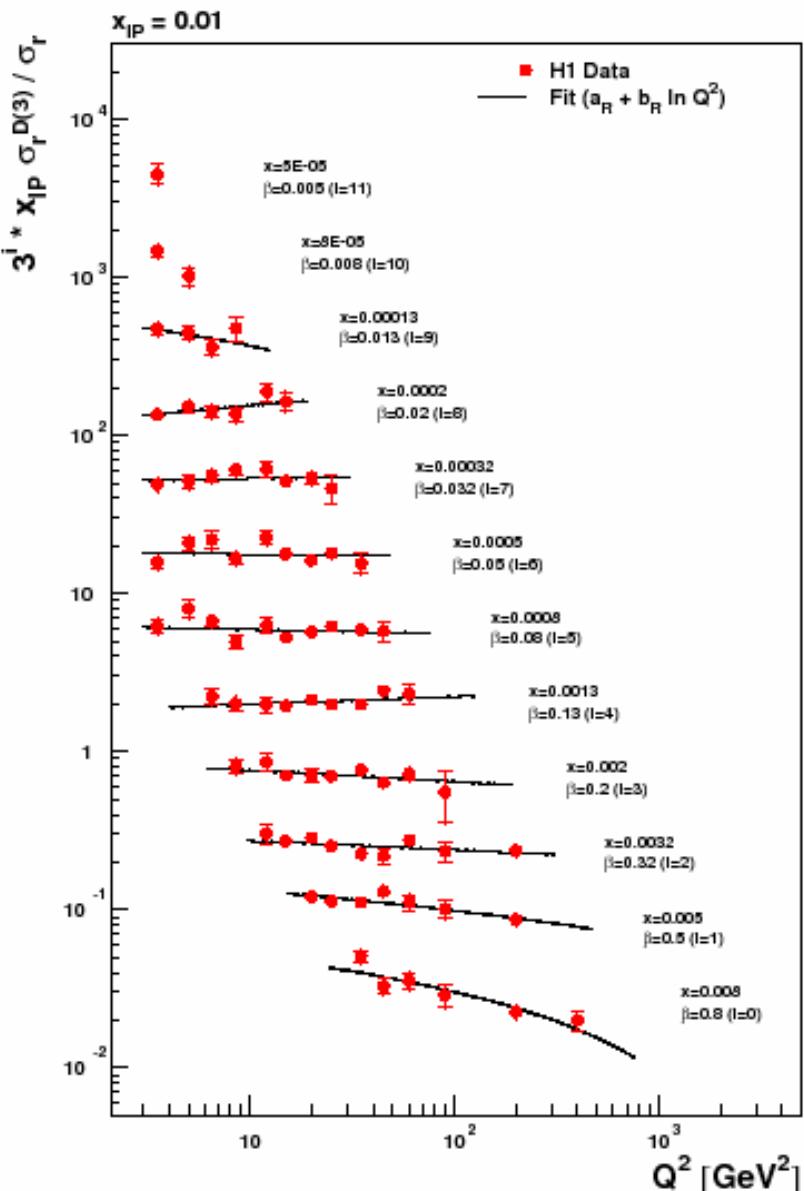
- Don't confirm  $\alpha_{IP}(Q^2)$  seen by ZEUS

- Consistent result from fits to FPS data:

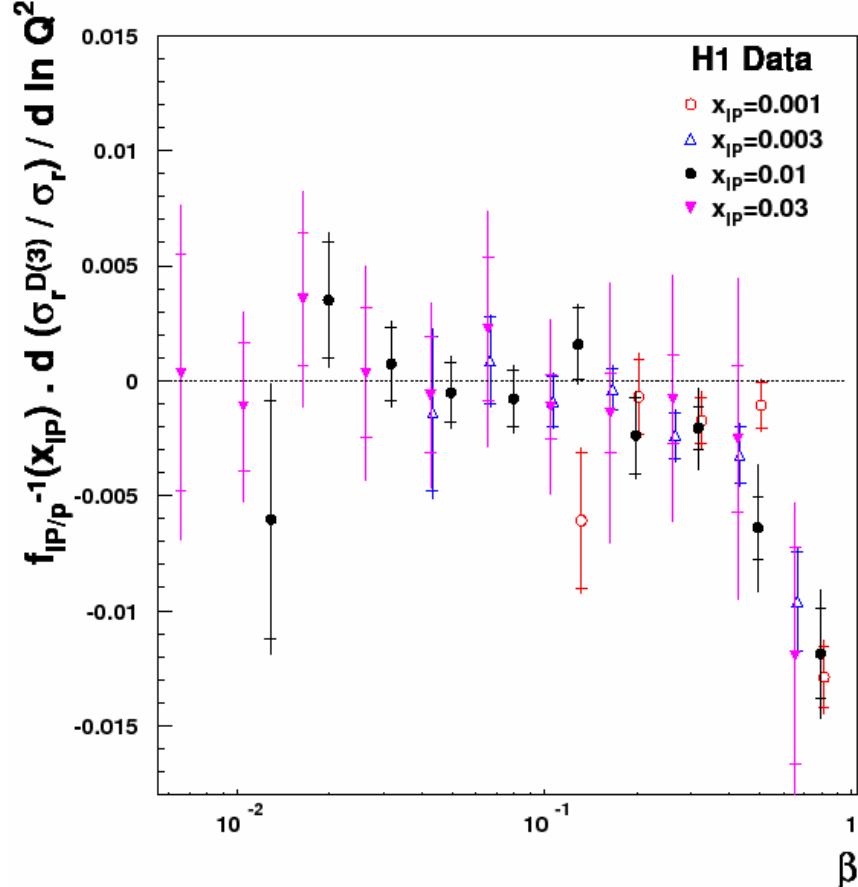
$$\alpha_{IP}(0) = 1.114 \pm 0.018(\text{stat.}) \pm 0.012(\text{syst.})^{+0.040}_{-0.020}(\text{th.})$$



# Ratio diffractive / inclusive DIS



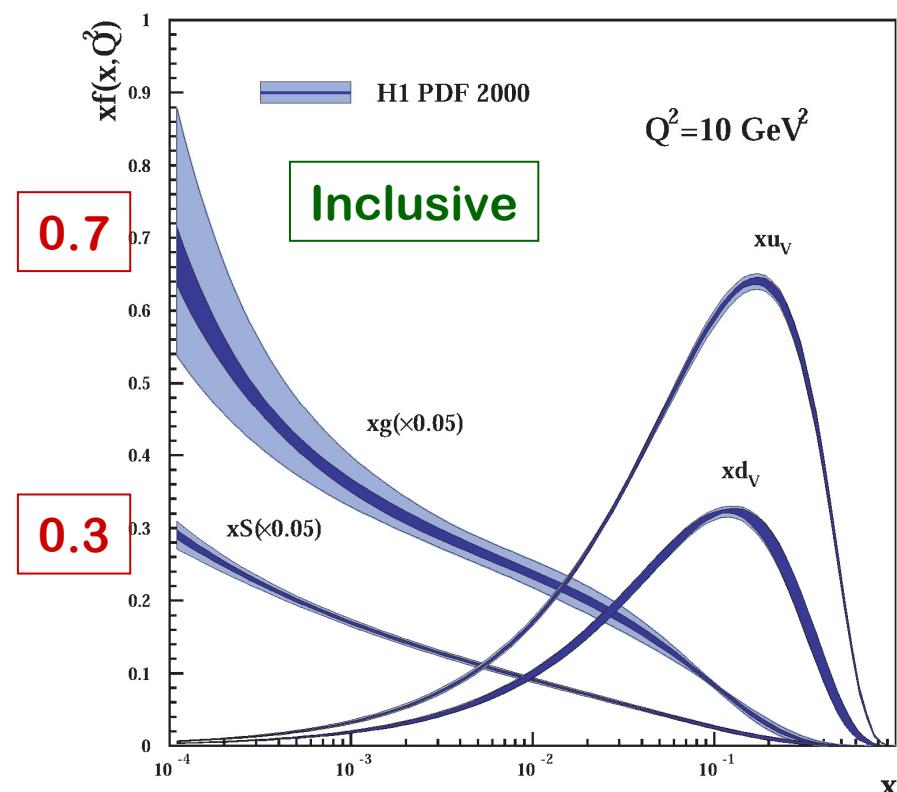
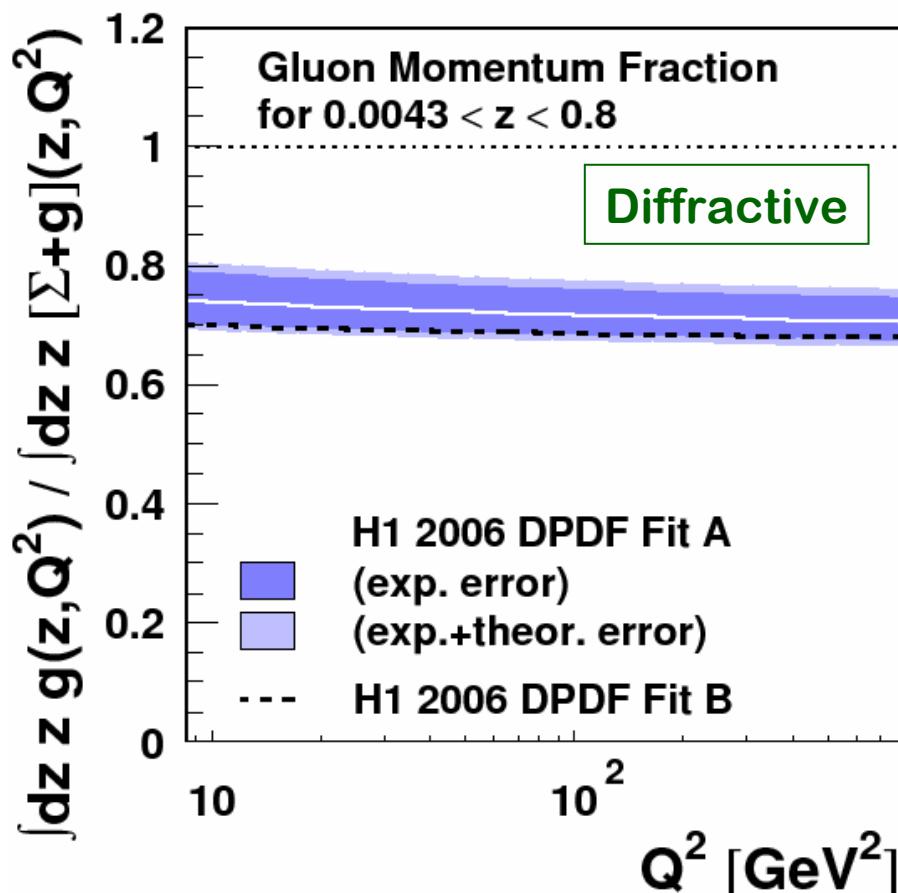
- Fit ratio at fixed ( $x_{IP}, x$ ) to  $A+B \ln Q^2$



- Ratio remarkably flat (derivative~0) except at high  $\beta$
- Ratio also measured vs  $x$  (not shown): approx. flat as well

# $Q^2$ derivative and gluon/quark ratios

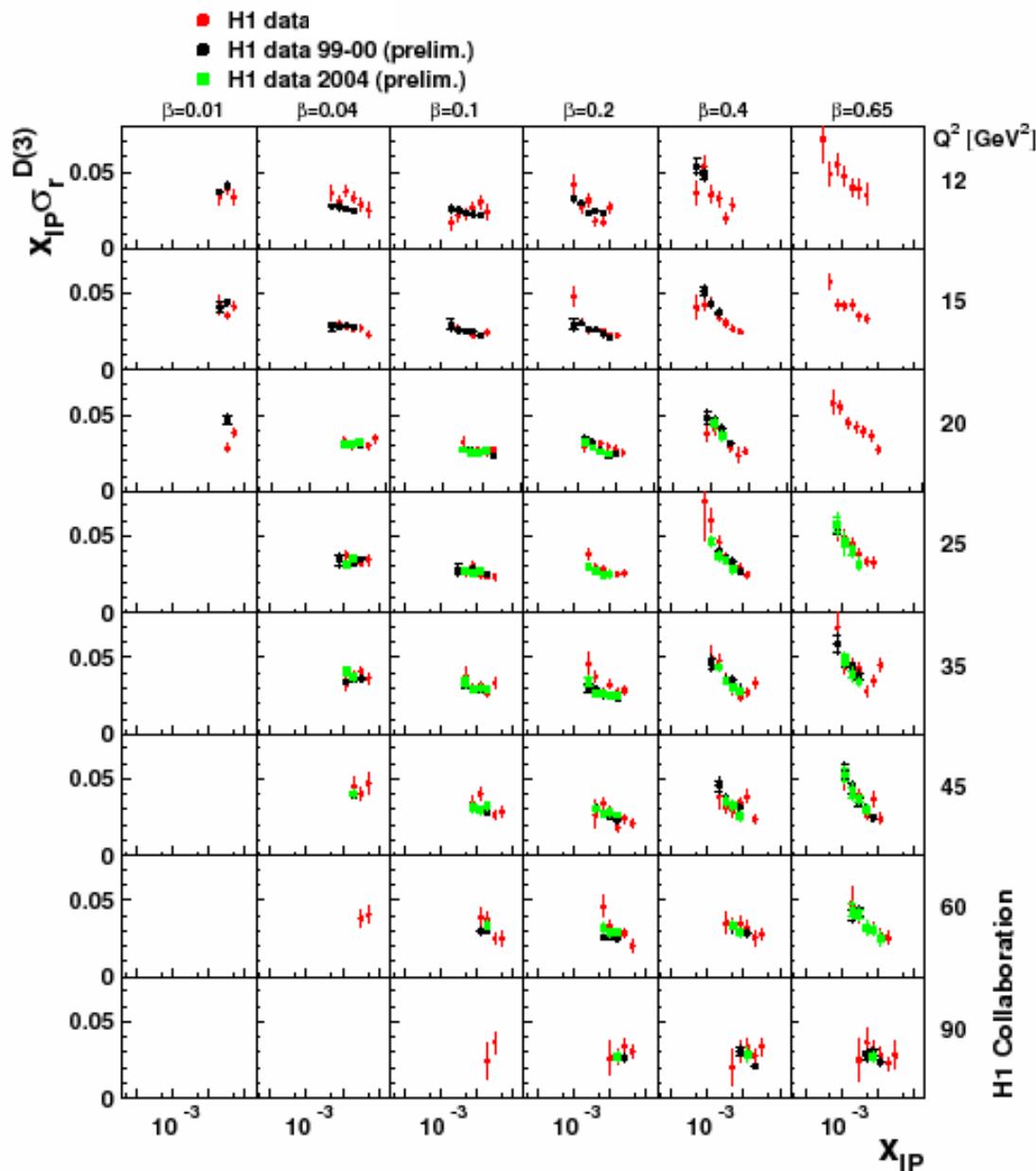
- If  $\frac{d(\sigma_r^D / \sigma_r)}{d \ln Q^2} \sim 0$  then  $\frac{1}{\sigma_r^D} \frac{d\sigma_r^D}{d \ln Q^2} \approx \frac{1}{\sigma_r} \frac{d\sigma_r}{d \ln Q^2} \rightarrow \frac{g^D}{q^D} \sim \frac{g}{q}$



At low  $x$ , quark:gluon ratio  $\sim 70\% / 30\%$ , common to diffractive and inclusive

# New Data using Rapidity Gap Method

- Published data
- Prel. 99-00 data,  $34 \text{ pb}^{-1}$ 
  - $10 < Q^2 < 105 \text{ GeV}^2$
- Prel. 2004 data,  $34 \text{ pb}^{-1}$ 
  - $17.5 < Q^2 < 105 \text{ GeV}^2$
- Large increase in statistics
- Consistent with published data
- For details, see contributed paper



# Summary

- H1 diffractive measurements using FPS and LRG methods published
  - hep-ex/0606003 and hep-ex/0606004 (both subm. to EPJC)
  - Data from two methods agree in detail! Also agreement with ZEUS-LPS
- New preliminary H1 data with large statistics
- Proton vertex factorization with reggeon exchanges at high  $x_{IP}$  provides good model for  $x_{IP}$  dependence:  $\alpha_{IP}(t) \sim 0.118 + 0.06 t$
- Ratio diffractive/inclusive DIS measured
  - ~flat with  $Q^2$  (fixed  $x, x_{IP}$ ), also with  $W$  (fixed  $Q^2, M_X$ )
- Diffractive PDFs extracted from NLO QCD fits to  $\beta, Q^2$  dependences for  $Q^2 \geq 8.5 \text{ GeV}^2$  (H1 2006 DPDF Fits A+B)
  - Quark singlet very well constrained (~5%)
  - Gluon constrained to ~15%, but poorly known at high  $z$
  - New DPDFs basis for prediction of diffractive cross sections at HERA, TEVATRON and LHC!
  - Use diffractive dijets to constrain gluon at high  $z$  (see talk by M.Kapishin)

# Backup / Extra Figures

# Data Sets and Observables

- **FPS data sample**
  - 1999-2000 data ( $28 \text{ pb}^{-1}$ )
  - Study of t dependence:
- **LRG data sample**
  - 1997 data ( $2 \text{ pb}^{-1}$ ,  $Q^2 < 13.5 \text{ GeV}^2$ )
  - 1997 data ( $11 \text{ pb}^{-1}$ ,  $13.5 < Q^2 < 105 \text{ GeV}^2$ )
  - 99-00 data ( $62 \text{ pb}^{-1}$ ,  $Q^2 > 133 \text{ GeV}^2$ )

$$x_{IP} \frac{d^2\sigma^{ep \rightarrow eXp}}{dx_{IP} dt}$$

- The Diffractive reduced cross section:

$$\frac{d^4\sigma^{ep \rightarrow eXp}}{dx dQ^2 dx_{IP} dt} = \frac{4\pi\alpha^2}{xQ^4} Y_+ \sigma_r^{D(4)}(x, Q^2, x_{IP}, t)$$

- Relates to the structure functions  $F_2^D$  and  $F_L^D$  as:

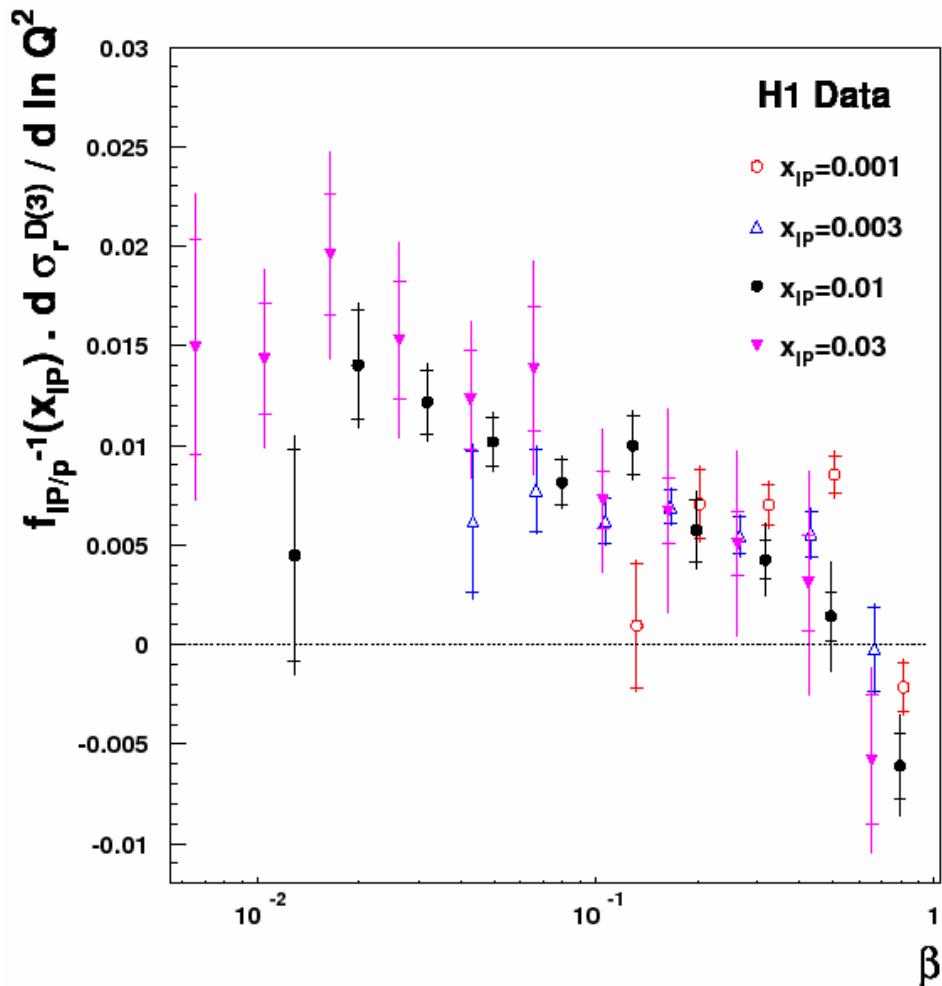
$$\sigma_r^{D(4)}(x, Q^2, x_{IP}, t) = F_2^{D(4)} - \frac{y^2}{Y_+} F_L^{D(4)} \approx F_2^{D(4)}$$

- Integrated over t:

$$\sigma_r^{D(3)}(x, Q^2, x_{IP}) = \int_{-1}^{t_{min}} \sigma_r^{D(4)}(x, Q^2, x_{IP}, t) dt$$

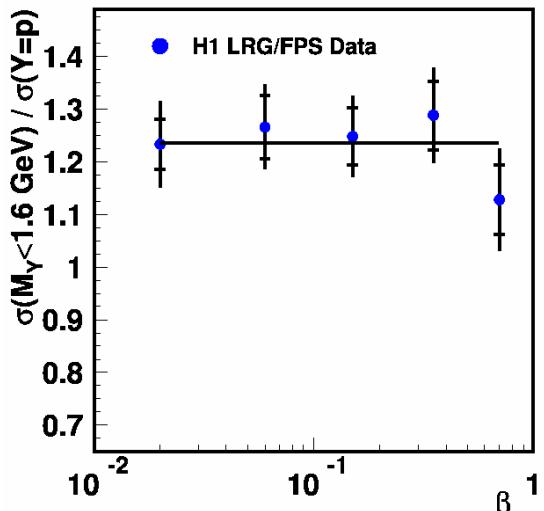
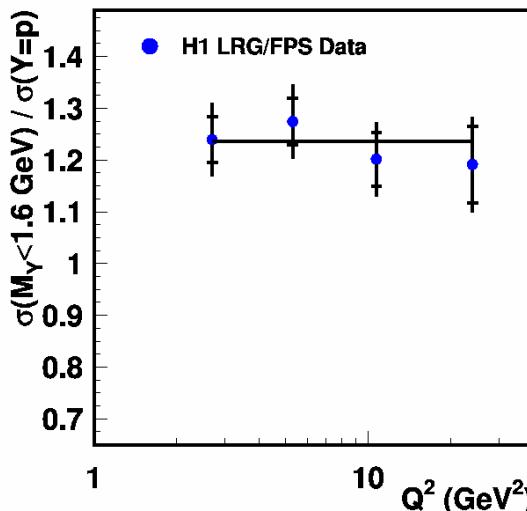
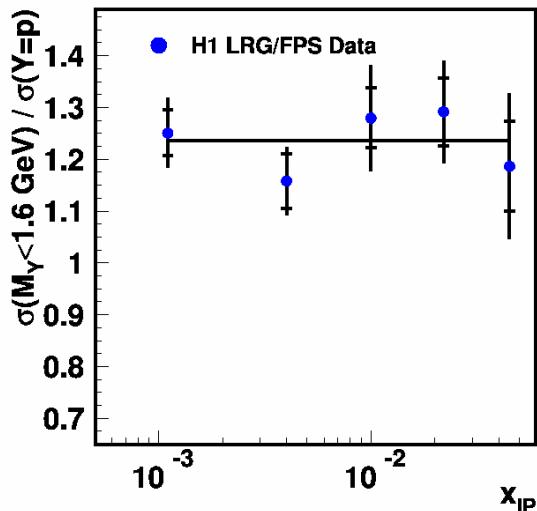
# Logarithmic $Q^2$ derivative

- $\sigma_r^{D(3)}$  measures diffractive quark density
- Its dependence on  $Q^2$  is sensitive to diffractive gluon density
- Fit data at fixed  $(x, x_{IP})$  to  $\sigma_r^D = A + B \ln Q^2$ , so that  $B = d\sigma_r^D / d \ln Q^2$
- Divide results by  $f_{IP/p}(x_{IP})$  to compare different  $x_{IP}$  values
- Derivatives large and positive at low  $\beta$
- Suggests large gluon density (independent of  $x_{IP}$  within errors)



# Comparison LRG vs FPS Data

- Form ratio LRG/FPS of measurements as function of  $x_{IP}$ ,  $\beta$  or  $Q^2$  after integration over others



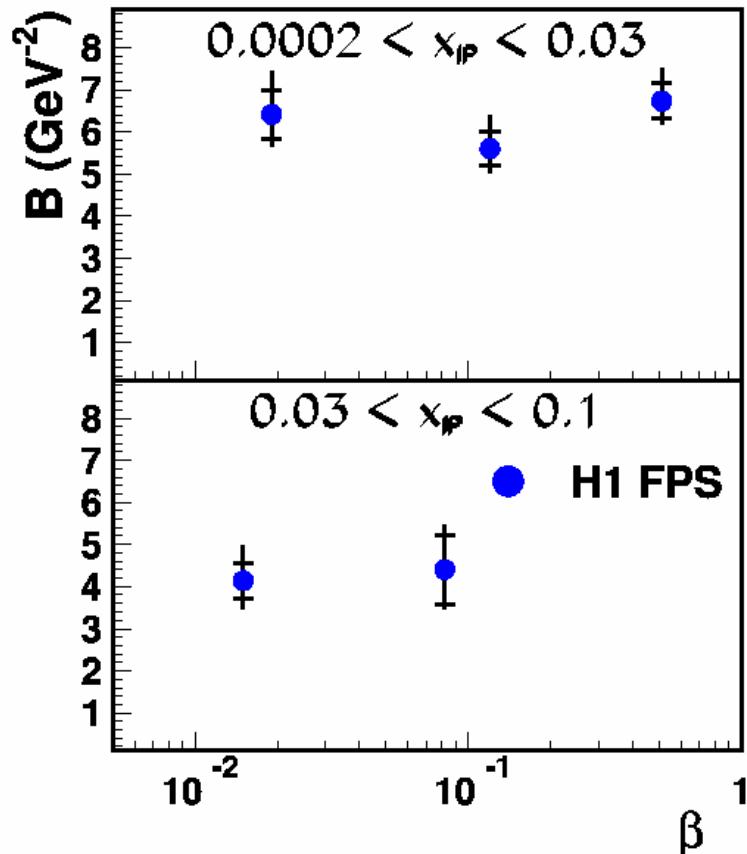
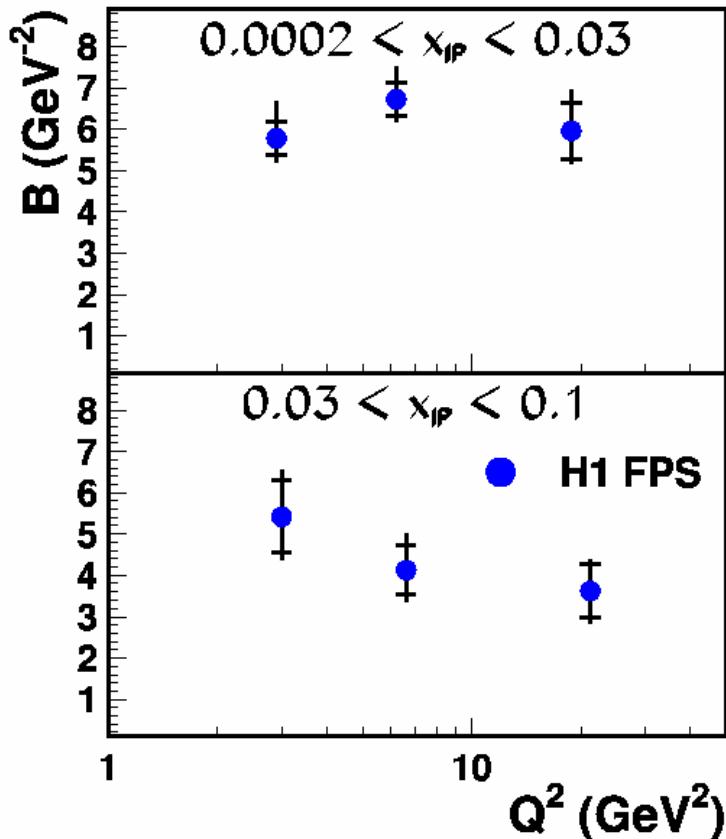
- Independent of kinematics within errors

$$\frac{\sigma(M_Y < 1.6 \text{ GeV})}{\sigma(Y=p)} = 1.23 \pm 0.03(\text{stat.}) \pm 0.16(\text{syst.})$$

- Agreement in detail between methods
- $M_Y$  dependence factorizes within (10% non-normalization) errors!

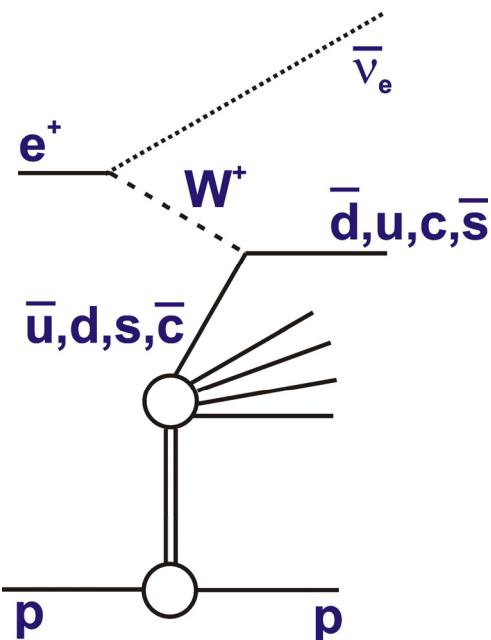
# $t$ -slope dependence on $\beta$ or $Q^2$ ?

- $B$  measured double differentially in ( $\beta$  or  $Q^2$ ) and  $x_{IP}$

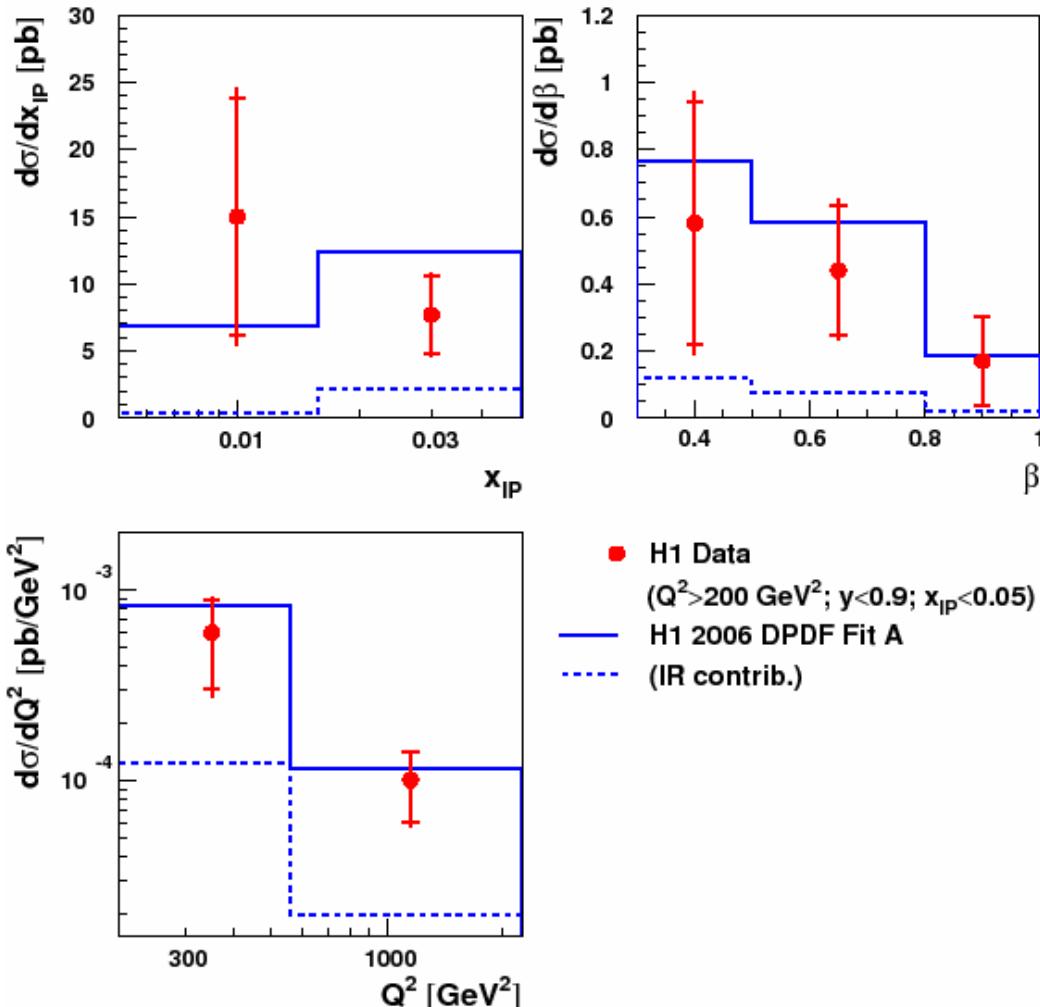


- No change of  $t$  dependence with  $\beta$  or  $Q^2$  at fixed  $x_{IP}$
- Proton vertex factorization for  $t$  dependence working within errors

# Diffractive Charged Current



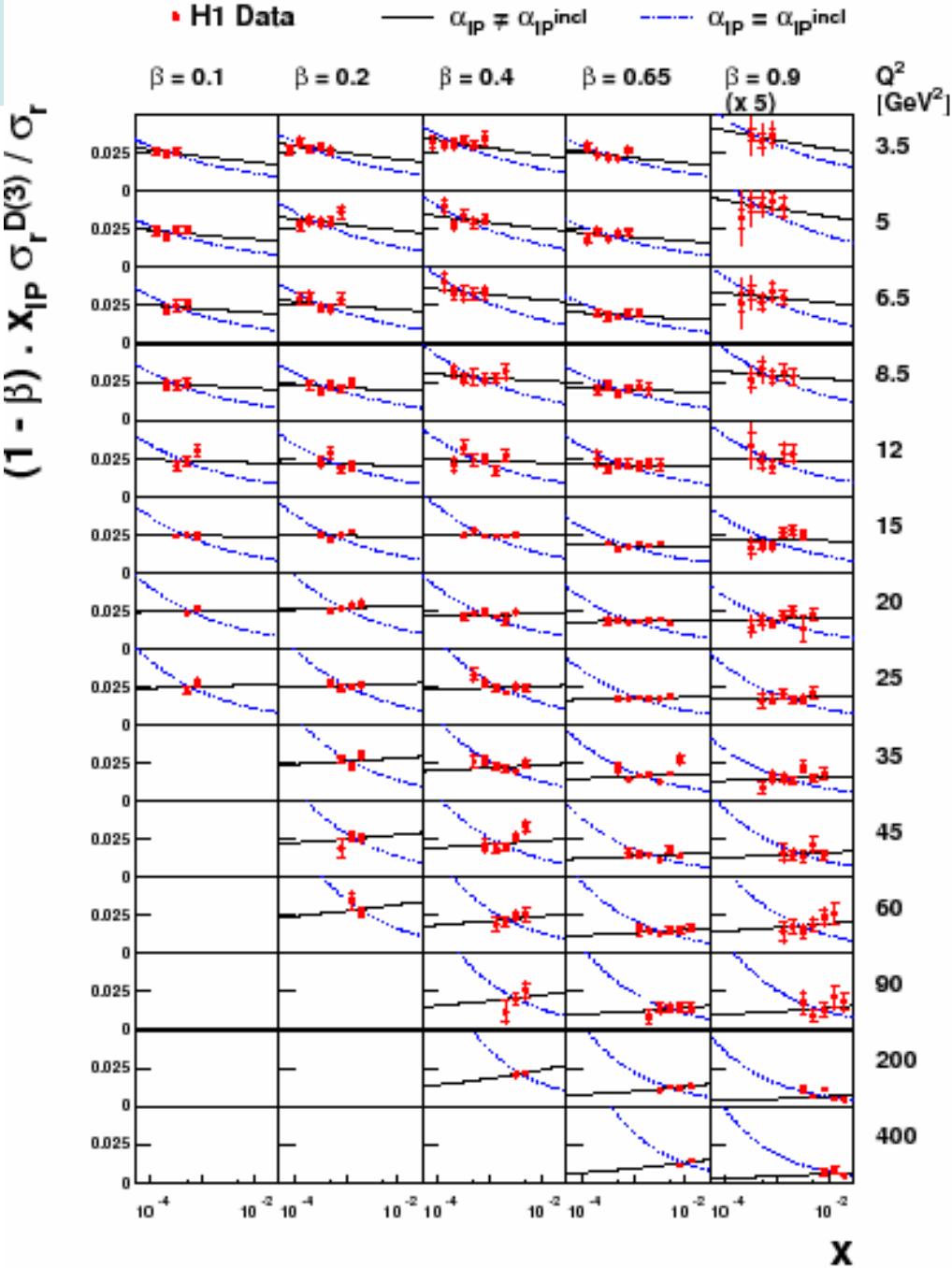
- Sensitive to flavour decomposition of singlet (completely unconstrained by NC data)
- Good agreement with 2006 DPDF fit (assumes  $u=d=s=u\bar{d}=d\bar{u}=s\bar{u}$ ,  $c$  from BGF), though statistics very limited so far



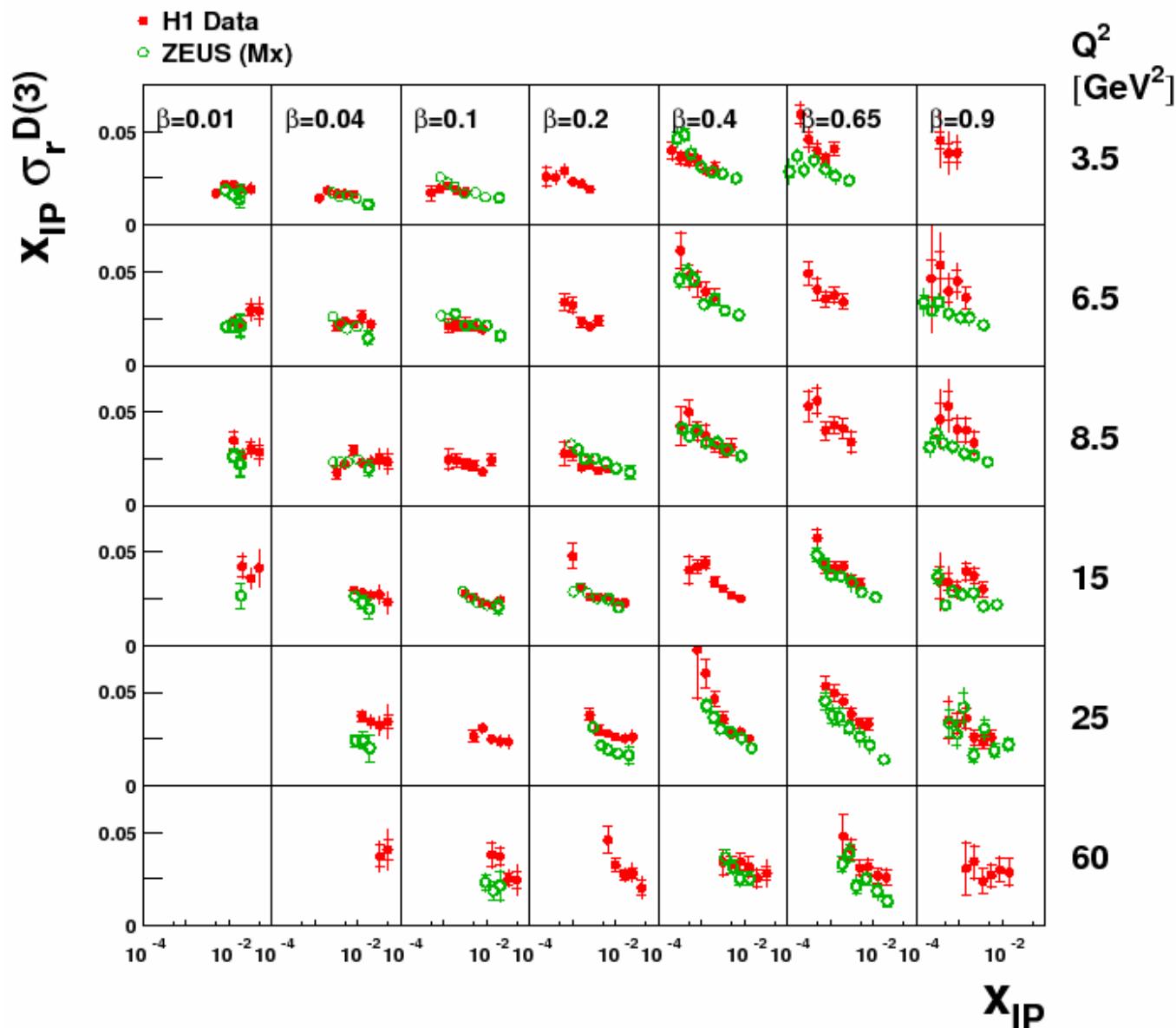
## Ratio diffractive/inclusive: x dependence

- Plot  $\sigma_r^D/\sigma_r$  at fixed  $\beta, Q^2$  (hence fixed  $M_X$ ) vs  $x$  ( $\sim 1/W^2$ )
- Corresponds to  

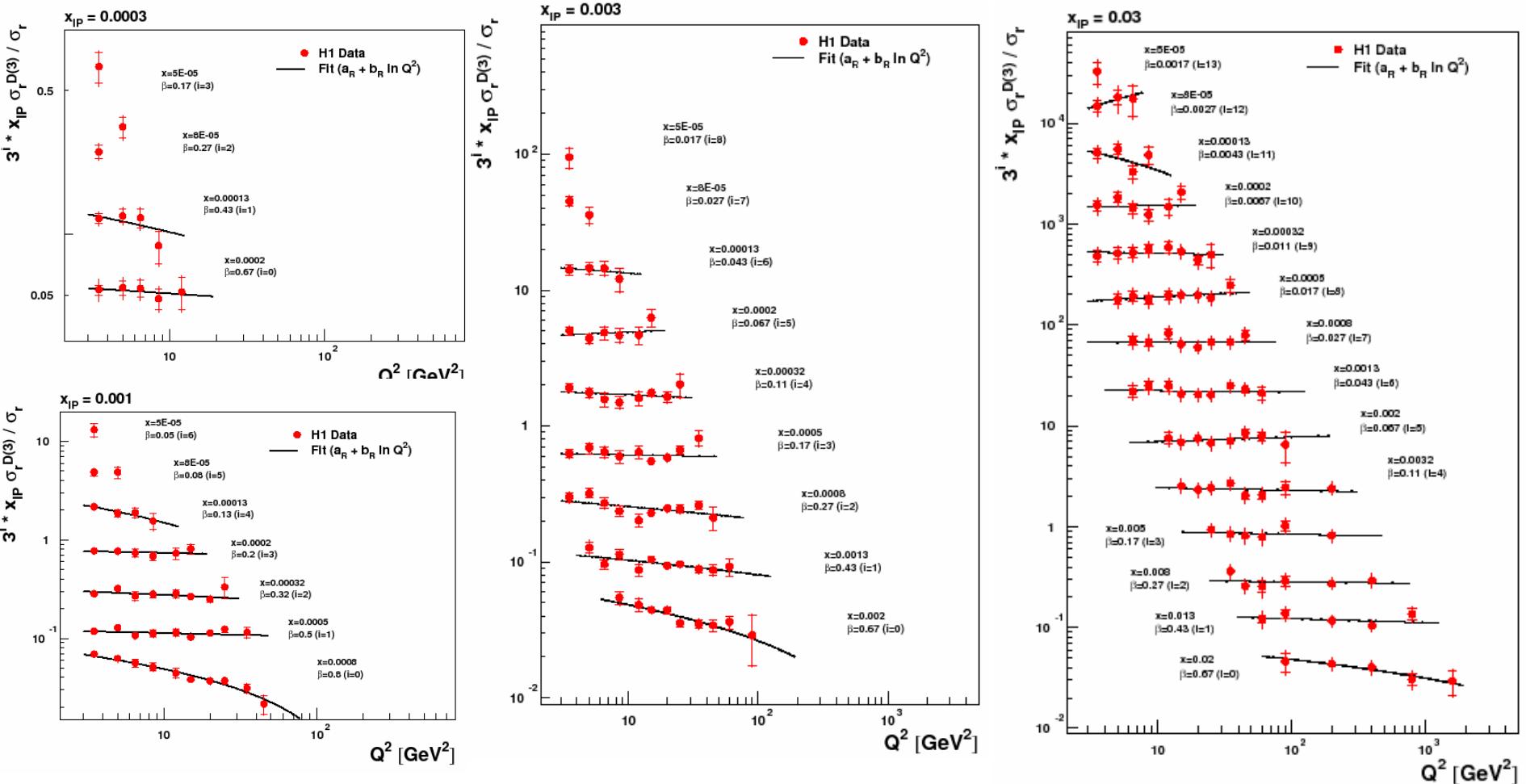
$$M_X^2 \cdot \frac{d\sigma_r^D}{dM_X^2} / \sigma_{tot}$$
- Remarkably flat vs  $x$  over most of kinematic range (bins with large  $F_L$  or IR contrib not shown)
- Diffractive and inclusive cross sections cannot be described with the same  $\alpha_{IP}(0)$ , even if it is  $Q^2$  dependent



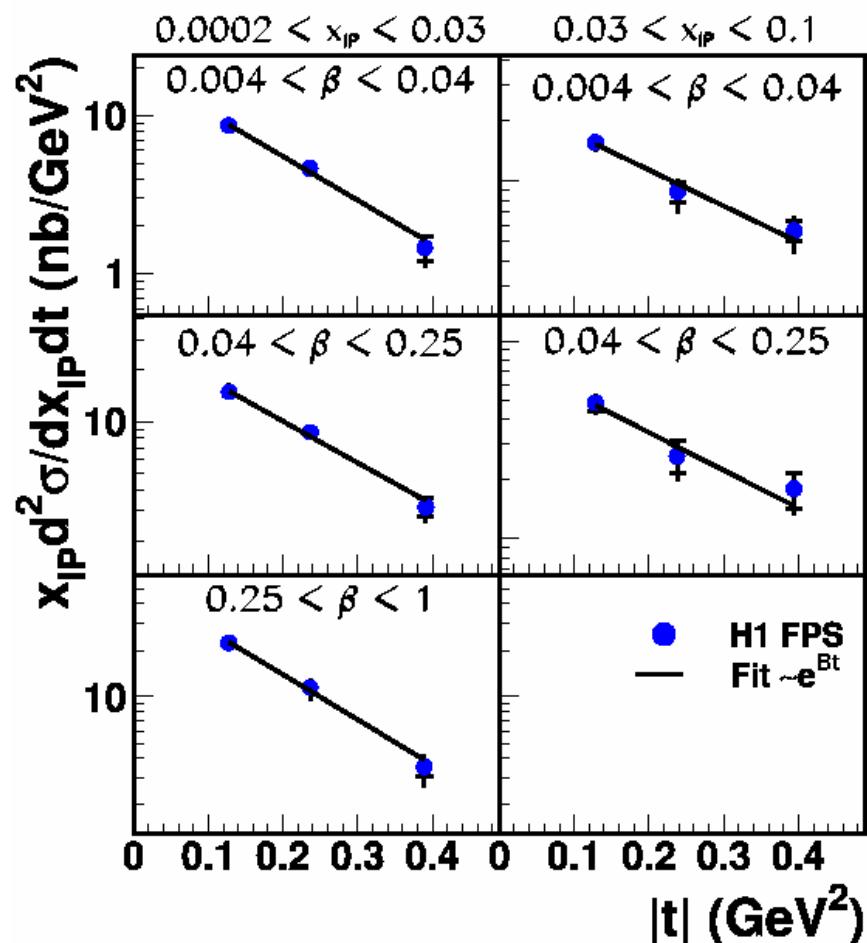
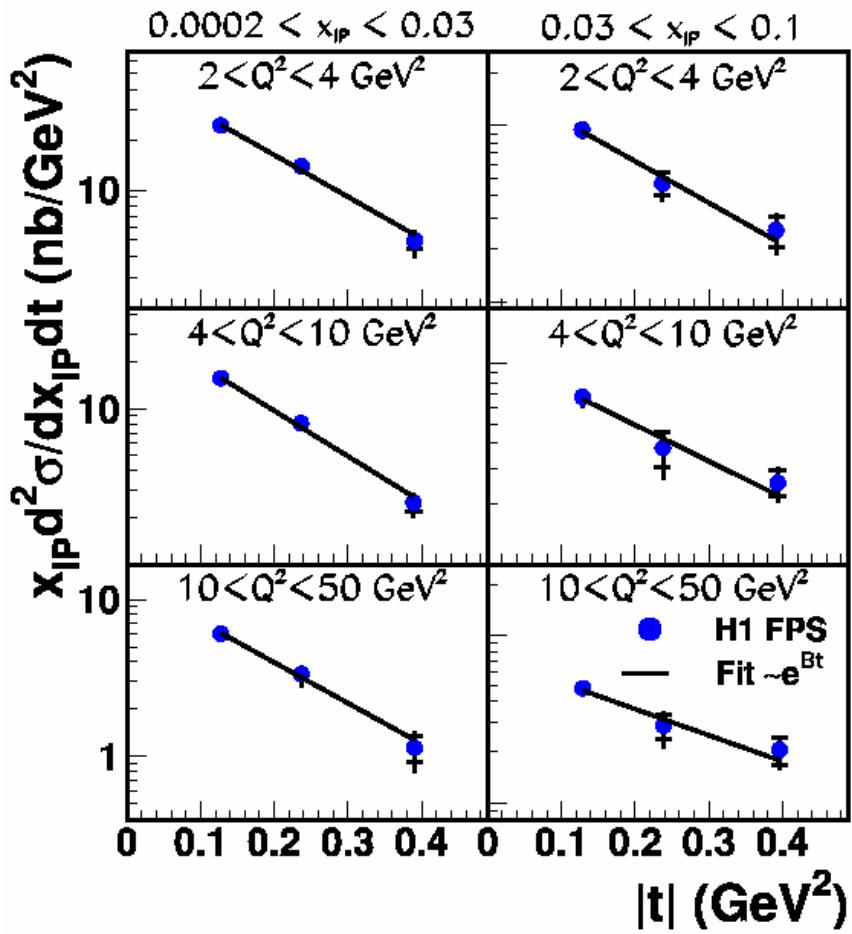
# Comparison H1-LRG and ZEUS-Mx data



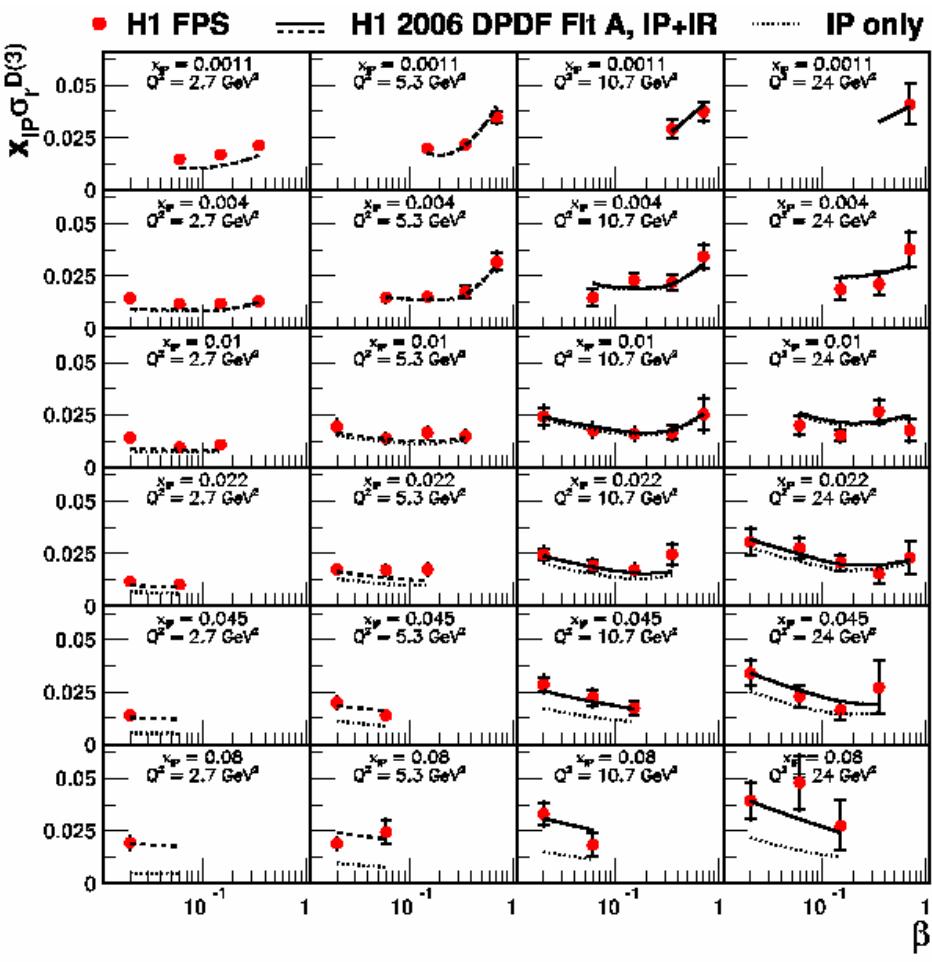
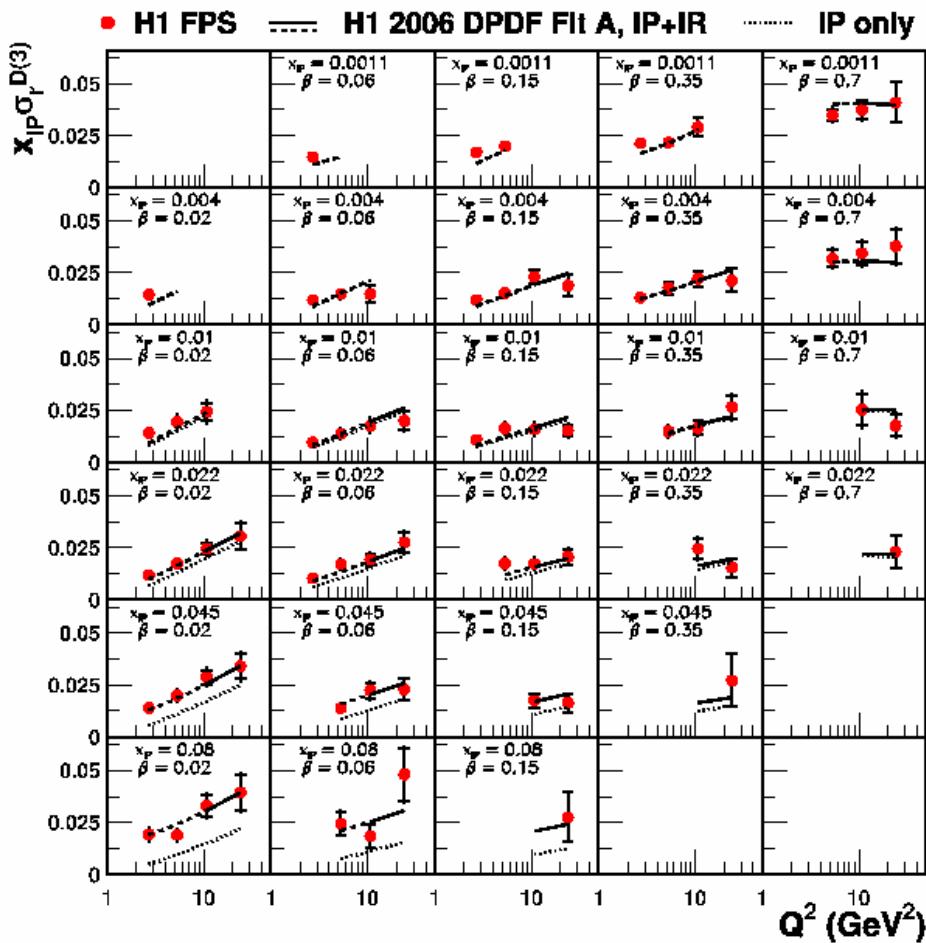
# Ratio diffractive/inclusive vs $Q^2$ extra plots



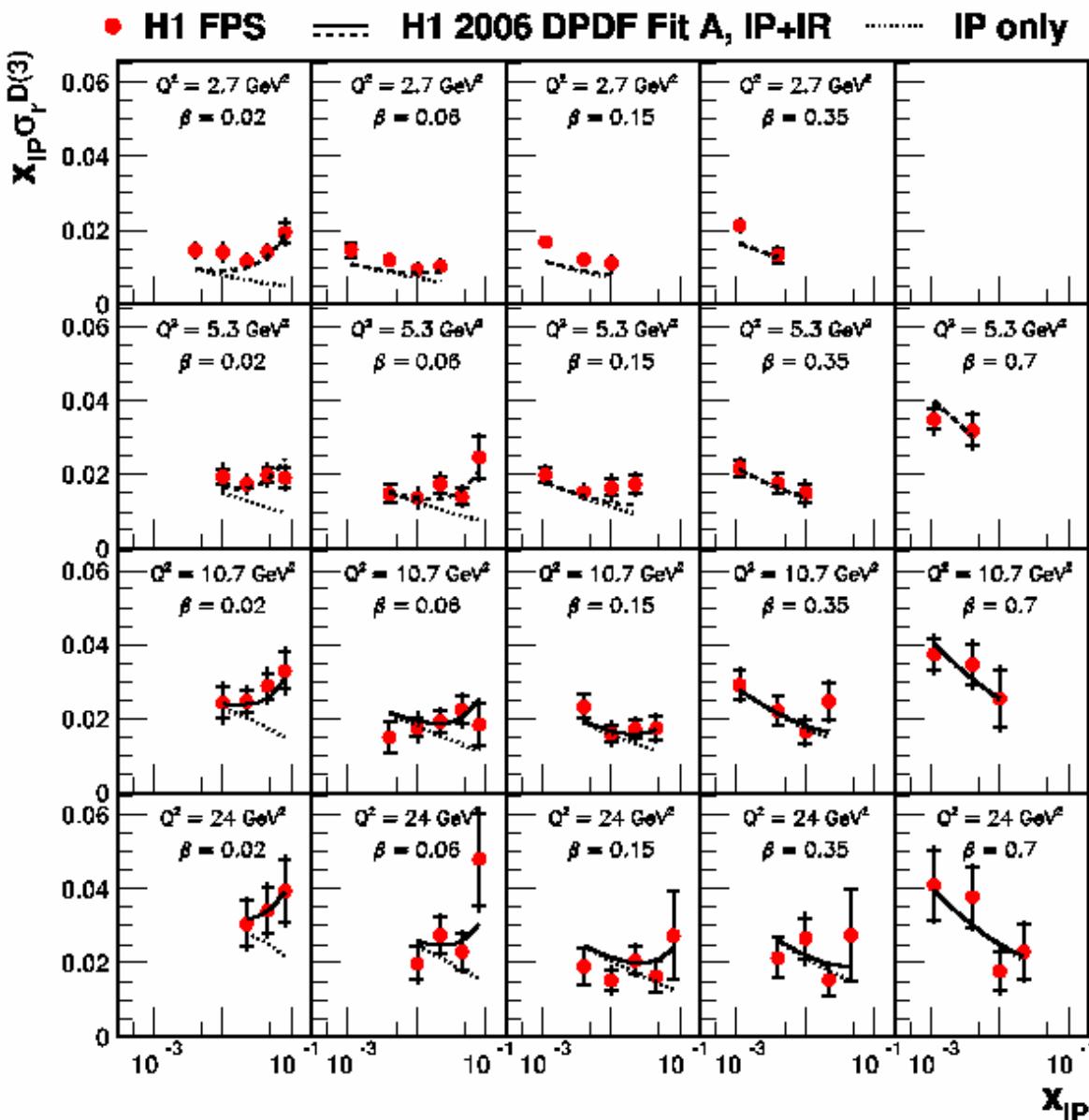
# $t$ dependence in bins of $\beta$ or $Q^2$



# $Q^2$ and $\beta$ dependences of FPS data



# $x_{IP}$ dependence of FPS data



# Final Results from two Publications

arXiv:hep-ex/0606004 v1 1 Jun 2006

DESY 06-049  
May 2006

ISSN 0418-9833

## Measurement and QCD Analysis of the Diffractive Deep-Inelastic Scattering Cross Section at HERA

H1 Collaboration

**DESY 06-049, hep-ex/0606004**

### Abstract

A detailed analysis is presented of the diffractive deep-inelastic scattering process  $e p \rightarrow e X Y$ , where  $Y$  is a proton or a low mass proton excitation carrying a fraction  $1 - x_p > 0.95$  of the incident proton longitudinal momentum and the squared four-momentum transfer at the proton vertex satisfies  $|t| < 1 \text{ GeV}^2$ . Using data taken by the H1 experiment, the cross section is measured for photon virtualities in the range  $3.5 \leq Q^2 \leq 1600 \text{ GeV}^2$ , triple differentially in  $x_p$ ,  $Q^2$  and  $\beta = x/x_p$ , where  $x$  is the Bjorken scaling variable. At low  $x_p$ , the data are consistent with a factorisable  $x_p$  dependence, which can be described by the exchange of an effective pomeron trajectory with intercept  $\alpha_p(0) = 1.118 \pm 0.008$  (exp.)  $^{+0.029}_{-0.011}$  (model). Diffractive parton distribution functions and their uncertainties are determined from a next-to-leading order DGLAP QCD analysis of the  $Q^2$  and  $\beta$  dependences of the cross section. The resulting gluon distribution carries an integrated fraction of around 70% of the exchanged momentum in the  $Q^2$  range studied. Total and differential cross sections are also measured for the diffractive charged current process  $e^+ p \rightarrow \nu_e X Y$  and are found to be well described by predictions based on the diffractive parton distributions. The ratio of the diffractive to the inclusive neutral current  $e p$  cross sections is studied. Over most of the kinematic range, this ratio shows no significant dependence on  $Q^2$  at fixed  $x_p$  and  $x$  or on  $x$  at fixed  $Q^2$  and  $\beta$ .

Submitted to *Eur. Phys. J. C*

DESY 06-048  
May 2006

ISSN 0418-9833

## Diffractive Deep-Inelastic Scattering with a Leading Proton at HERA

H1 Collaboration

**DESY 06-048, hep-ex/0606003**

### Abstract

The cross section for the diffractive deep-inelastic scattering process  $e p \rightarrow e X p$  is measured, with the leading final state proton detected in the H1 Forward Proton Spectrometer. The data analysed cover the range  $x_p < 0.1$  in fractional proton longitudinal momentum loss,  $0.08 < |t| < 0.5 \text{ GeV}^{-2}$  in squared four-momentum transfer at the proton vertex,  $2 < Q^2 < 50 \text{ GeV}^2$  in photon virtuality and  $0.004 < \beta = x/x_p < 1$ , where  $x$  is the Bjorken scaling variable. For  $x_p \lesssim 10^{-2}$ , the differential cross section has a dependence of approximately  $d\sigma/dt \propto e^{\beta t}$ , independently of  $x_p$ ,  $\beta$  and  $Q^2$  within uncertainties. The cross section is also measured triple differentially in  $x_p$ ,  $\beta$  and  $Q^2$ . The  $x_p$  dependence is interpreted in terms of an effective pomeron trajectory with intercept  $\alpha_p(0) = 1.114 \pm 0.018$  (stat.)  $\pm 0.012$  (syst.)  $^{+0.040}_{-0.020}$  (model) and a sub-leading exchange. The data are in good agreement with an H1 measurement for which the event selection is based on a large gap in the rapidity distribution of the final state hadrons, after accounting for proton dissociation contributions in the latter. Within uncertainties, the dependence of the cross section on  $x$  and  $Q^2$  can thus be factorised from the dependences on all studied variables which characterise the proton vertex, for both the pomeron and the sub-leading exchange.

Submitted to *Eur. Phys. J. C*