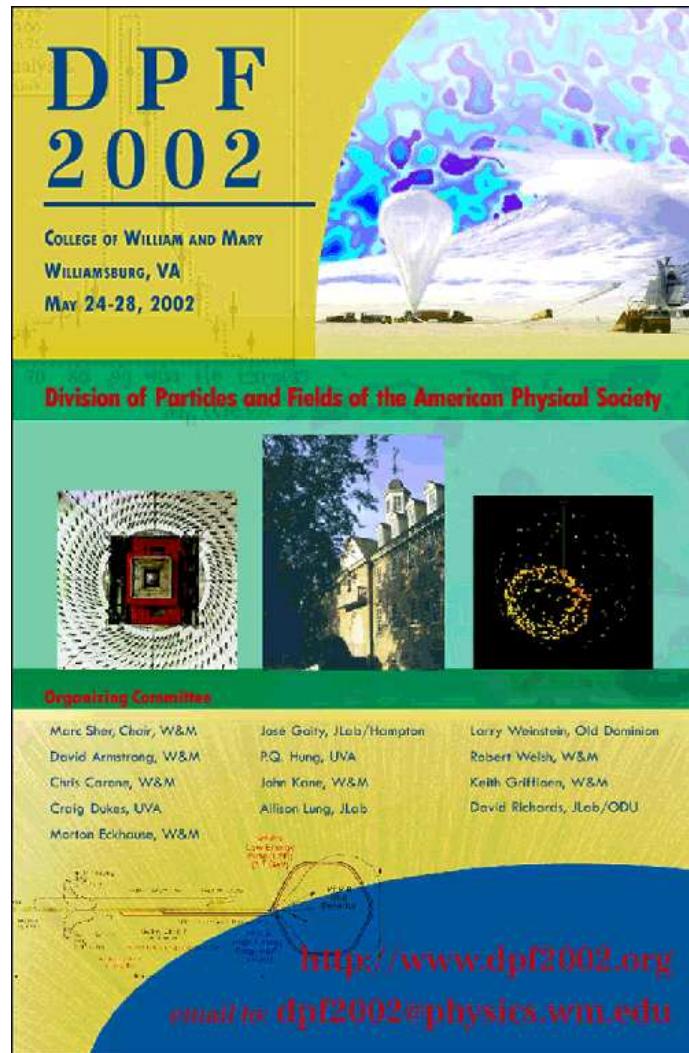


Hard diffraction at HERA



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H1 Collaboration

- Introduction
- Diffractive Structure Function F_2^D
- NLO QCD fit and diffractive pdf's
- Diffractive final states
- Summary

Preface: Why is diffraction (still) interesting ?

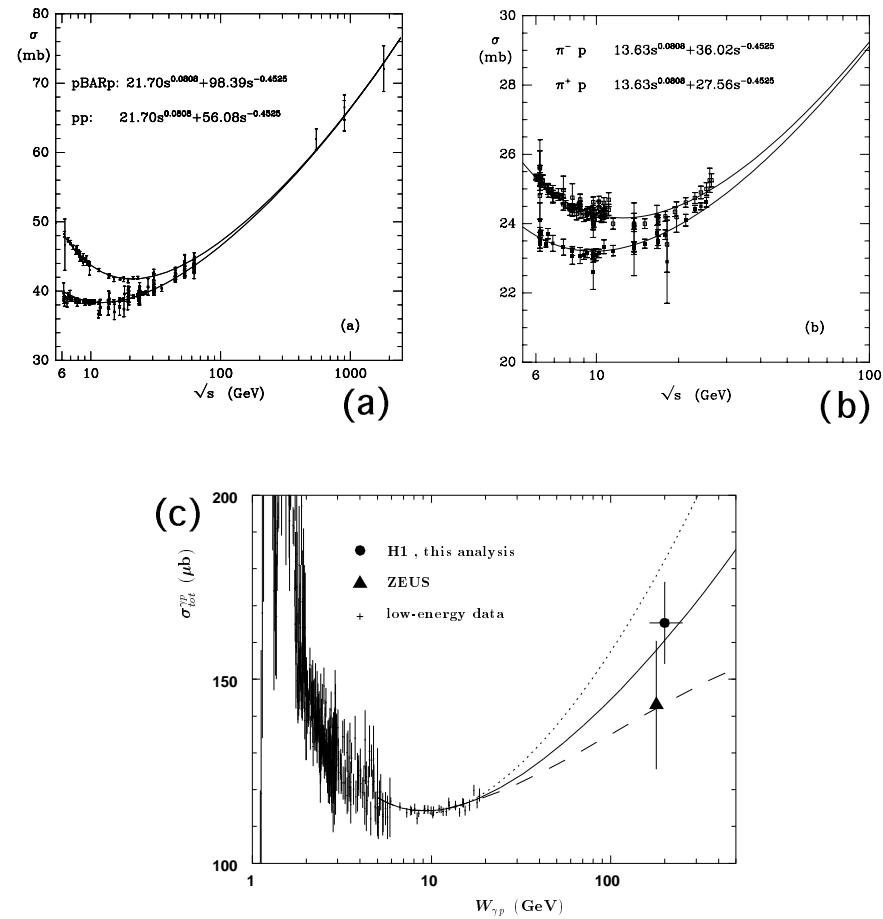
- Early observation:
total hadronic cross sections rise at high s
- σ_{tot} dominated by soft processes, where pQCD does not apply
large fraction of elastic / diffractive processes
- Parameterized in terms of
Regge phenomenology:

$$\frac{d\sigma}{dt} \sim \frac{1}{s^2} |T(s, t)|^2 = f(t) \left(\frac{s}{s_0}\right)^{2\alpha(t)-2}$$

$$\sigma_{\text{tot}} \sim \frac{1}{s} \text{Im}(T(s, t))|_{(t=0)} = s^{\alpha(0)-1}$$
- At high s : the pomeron trajectory
(vacuum quantum numbers, elastic scattering)

$$\alpha(t) = \alpha(0) + \alpha' t = 1.08 + 0.25 t$$
- Diffraction exists also in hard processes

Total cross sections at high energy:

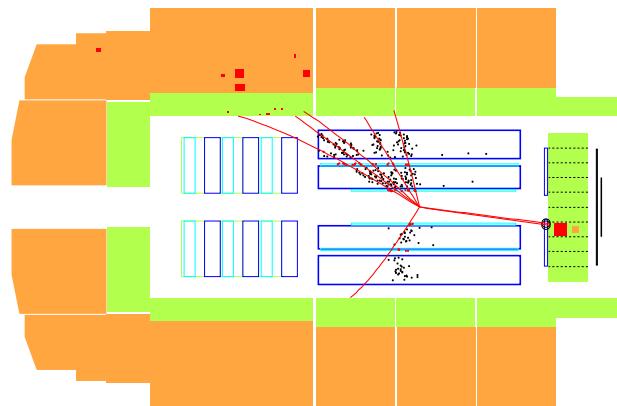


→ QCD (quark-gluon) structure of diffraction !

Diffraction in DIS at HERA

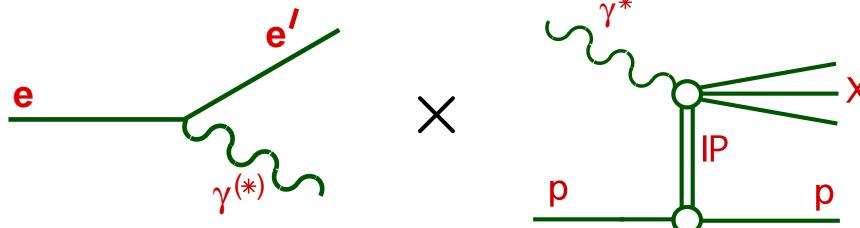
Early Observation at HERA:

10% of low- x DIS events are diffractive
 $ep \rightarrow e' p' X$

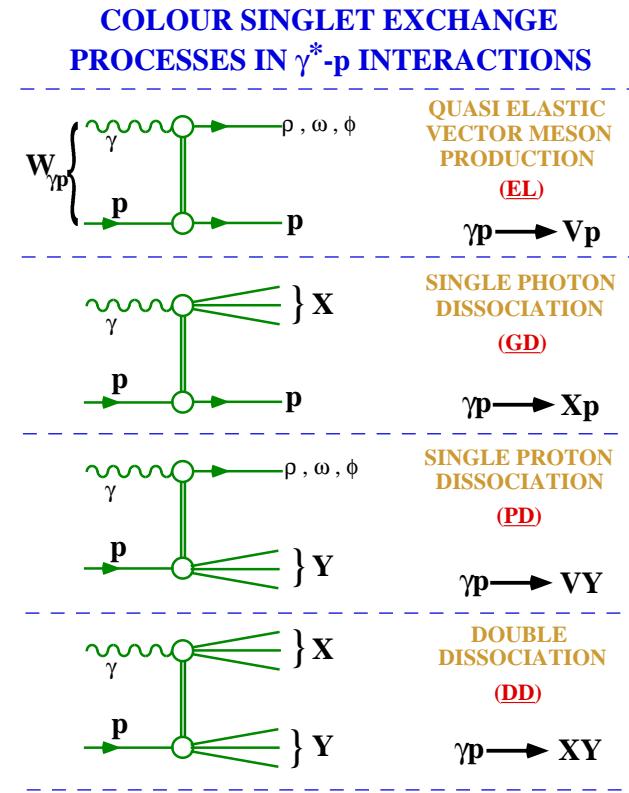


Colour-singlet or “pomeron” exchange

Can be viewed as diffractive $\gamma^* p$ interaction:



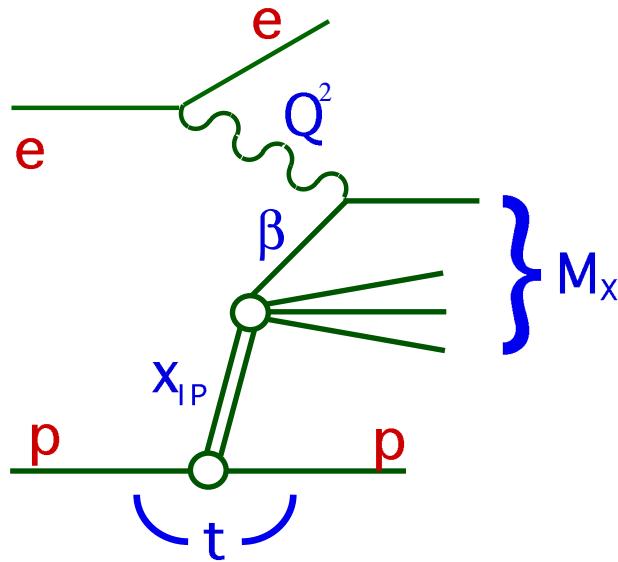
More generally: $\gamma^{(*)} p \rightarrow XY$



All can be measured by varying
 Q^2, W, t, M_X, M_Y

This talk mostly $\gamma^* p \rightarrow Xp$
 (large Q^2 , small $|t|$)

Diffractive DIS



$$x_{IP} = \xi = \frac{Q^2 + M_X^2}{Q^2 + W^2} = x_{IP}/p$$

(momentum fraction of colour singlet exchange)

$$\beta = \frac{Q^2}{Q^2 + M_X^2} = x_q/IP$$

(fraction of exchange momentum carried by q coupling to γ^* , hence $x = x_{IP}\beta$)

$$t = (p - p')^2$$

(4-momentum transfer squared at p vertex)

Diffractive reduced cross section σ_r^D :

$$\frac{d^4\sigma}{dx_{IP} dt d\beta dQ^2} = \frac{4\pi\alpha^2}{\beta Q^4} \left(1 - y + \frac{y^2}{2}\right) \sigma_r^{D(4)}(x_{IP}, t, \beta, Q^2)$$

Structure functions F_2^D and F_L^D :

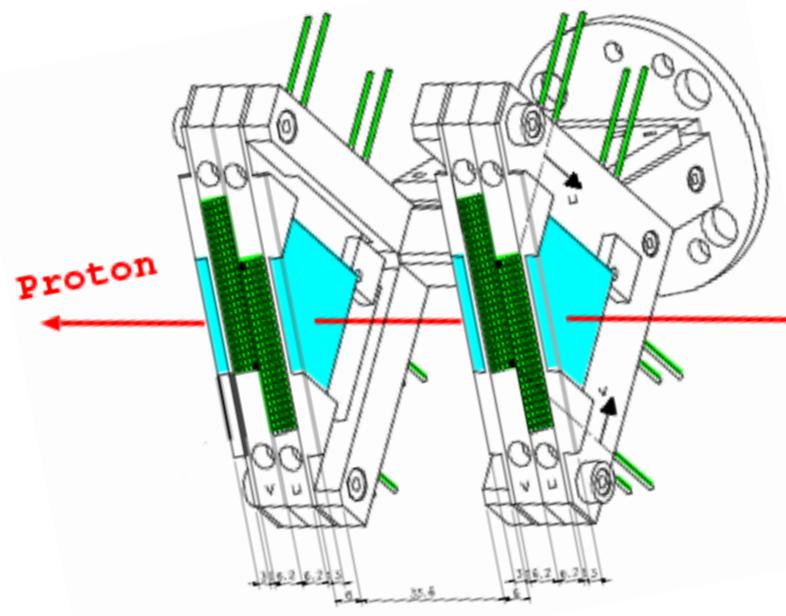
$$\sigma_r^{D(4)} = F_2^{D(4)} - \frac{y^2}{2(1-y+y^2/2)} F_L^{D(4)}$$

Integrated over t : $F_2^{D(3)} = \int dt F_2^{D(4)}$

Probe QCD (quark-gluon) structure of diffraction (Pomeron exchange) !

Experimental Techniques

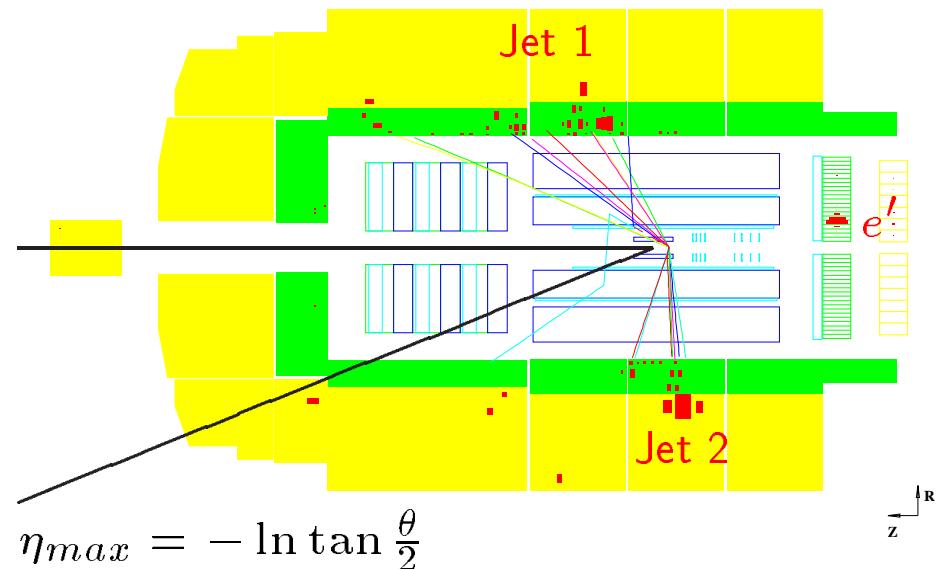
Forward Proton Spectrometer at $z = 65\ldots 90$ m



Measure leading proton

- Free of p dissociation bkgd.
 - Measure t distribution
 - low statistics (acceptance)

Rapidity Gap Selection in central detector



Require large rapidity gap

- $\Delta\eta$ large when $M_X \ll W$
 - integrate over M_Y, t
 - high statistics

Factorization Properties of F_2^D

QCD Factorization for diffractive DIS:

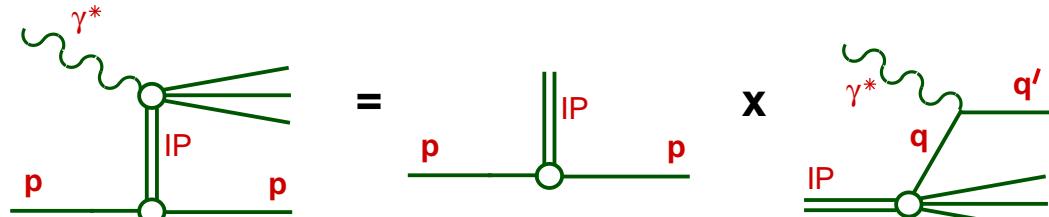
- Diffractive parton distributions (Trentadue, Veneziano, Berera, Soper, Collins, ...):

$$\frac{d^2\sigma(x, Q^2, x_{IP}, t) \gamma^* p \rightarrow p' X}{dx_{IP} dt} = \sum_i \int_x^{x_{IP}} d\xi \hat{\sigma}^{\gamma^* i}(x, Q^2, \xi) p_i^D(\xi, Q^2, x_{IP}, t)$$

- $\hat{\sigma}^{\gamma^* i}$ hard scattering part, as in incl. DIS
- p_i^D diffractive PDF's in proton, conditional probabilities, valid at fixed x_{IP}, t , obey DGLAP
- not proven for diffractive hadron-hadron scattering

Regge Factorization / resolved Pomeron model:

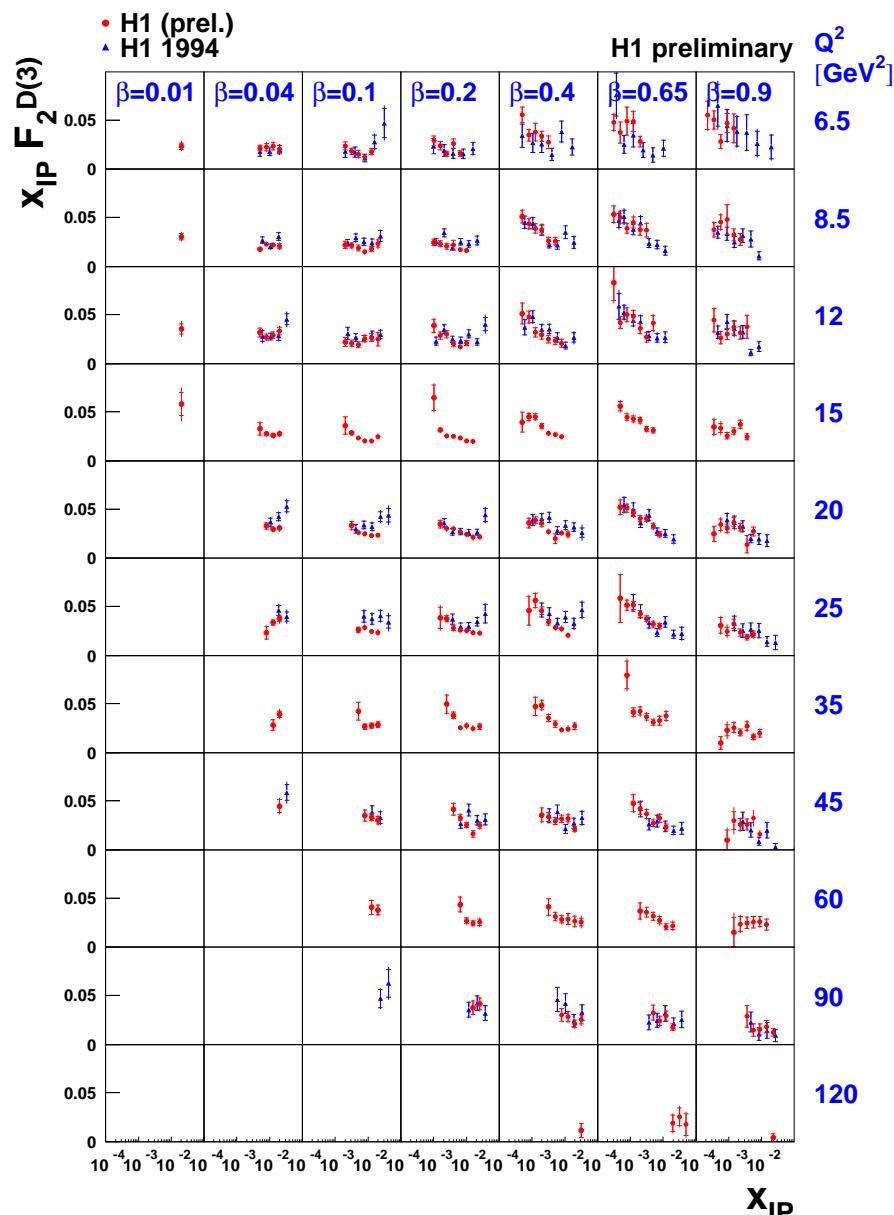
x_{IP}, t dependence factorizes out: Donnachie, Landshoff, Ingelman, Schlein, ...)



$$F_2^D(x_{IP}, t, \beta, Q^2) = f_{IP/p}(x_{IP}, t) F_2^{IP}(\beta, Q^2)$$

- additional assumption, no proof !
- consistent with present data if sub-leading \mathbf{R} included

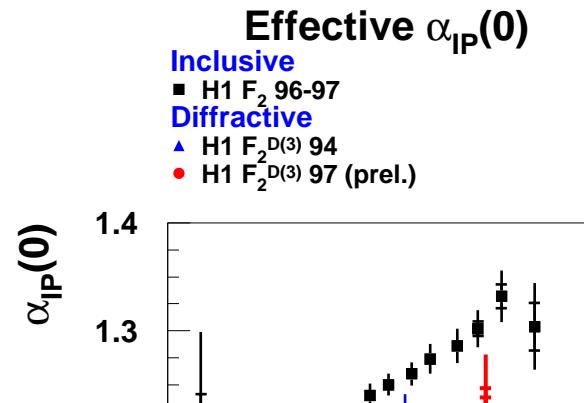
New $F_2^{D(3)}$ measurement (Rapidity gap)



5 times more statistics than previous data

Fit to x_{IP} dependence: effective $\alpha_{IP}(0)$

$$F_2^D(x_{IP}, \beta, Q^2) \sim B(\beta, Q^2) \left(\frac{1}{x_{IP}} \right)^{2\alpha_{IP}-1}$$

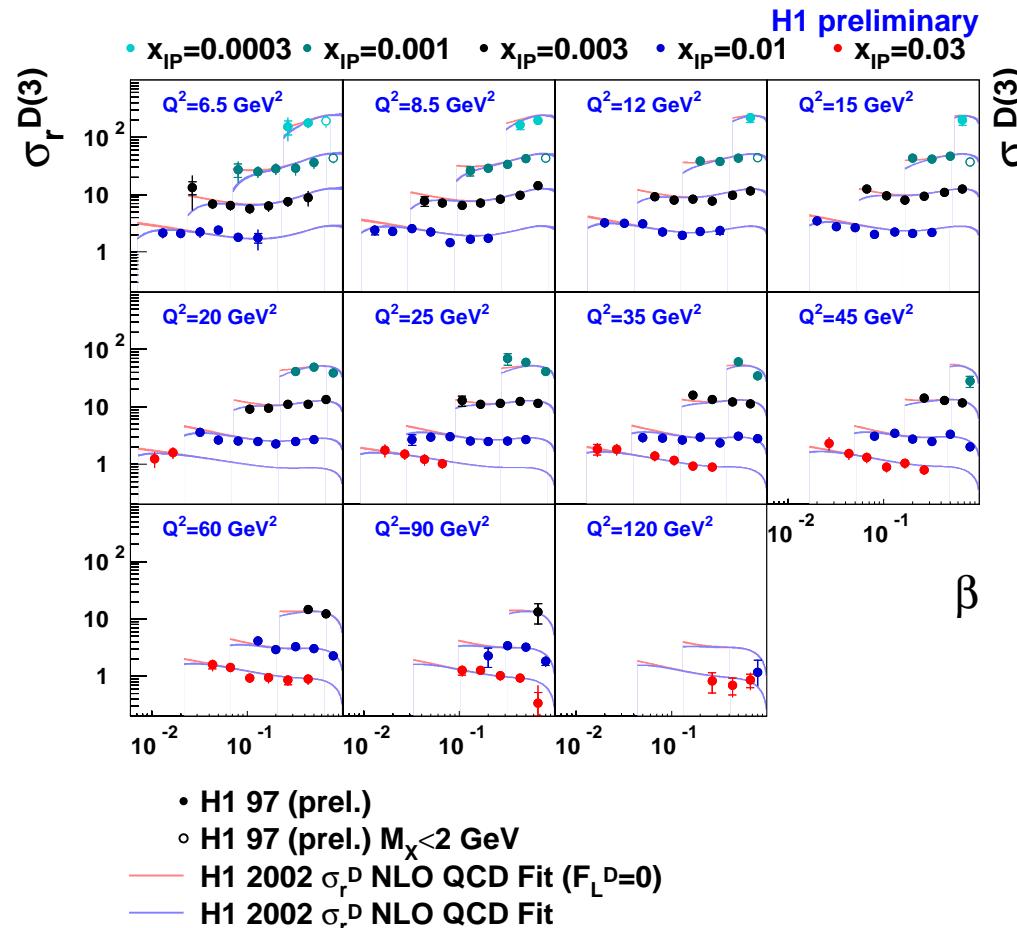


$$\alpha_{IP}(0) = 1.173 \pm 0.02 \pm 0.02^{+0.06}_{-0.03}$$

Growth with Q^2 slower in diffractive case?

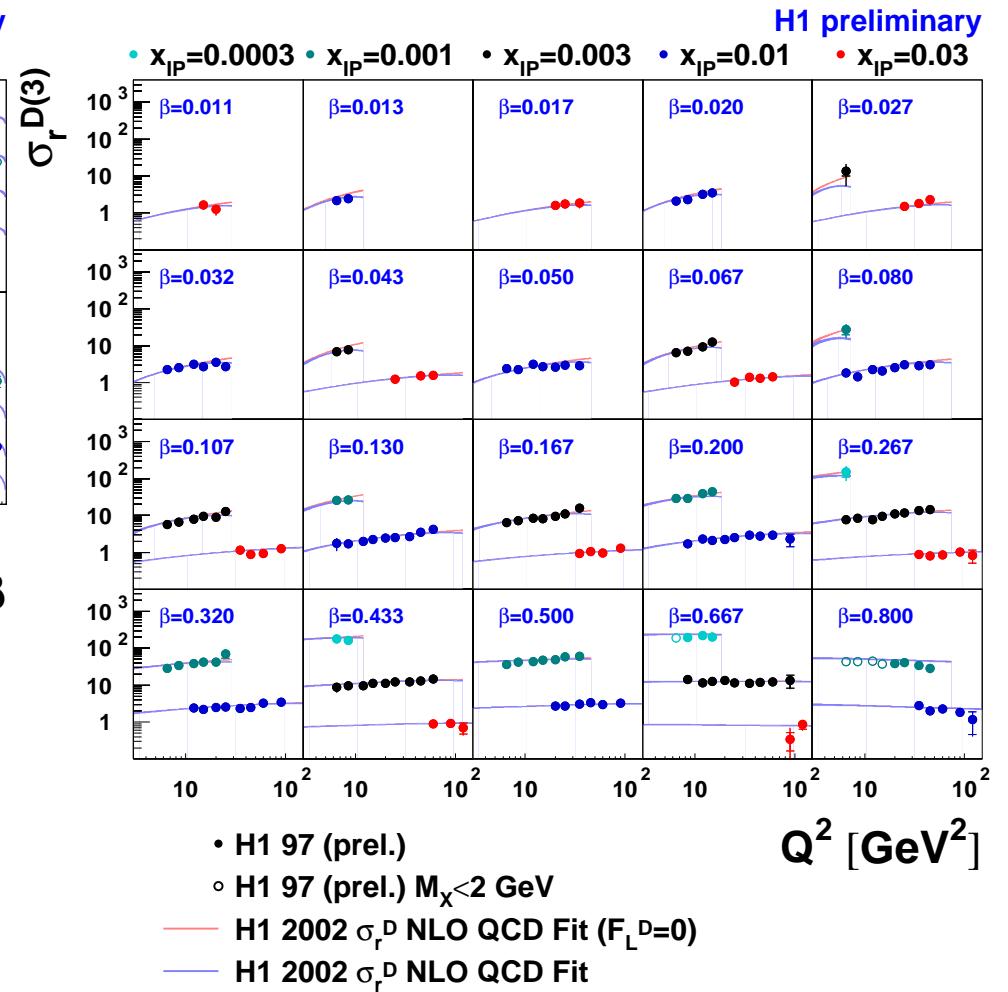
$F_2^{D(3)}$: β and Q^2 dependence overview

β dependence at fixed Q^2 :



sensitive to diffr. pdf's integrated over t

Q^2 dependence at fixed β :

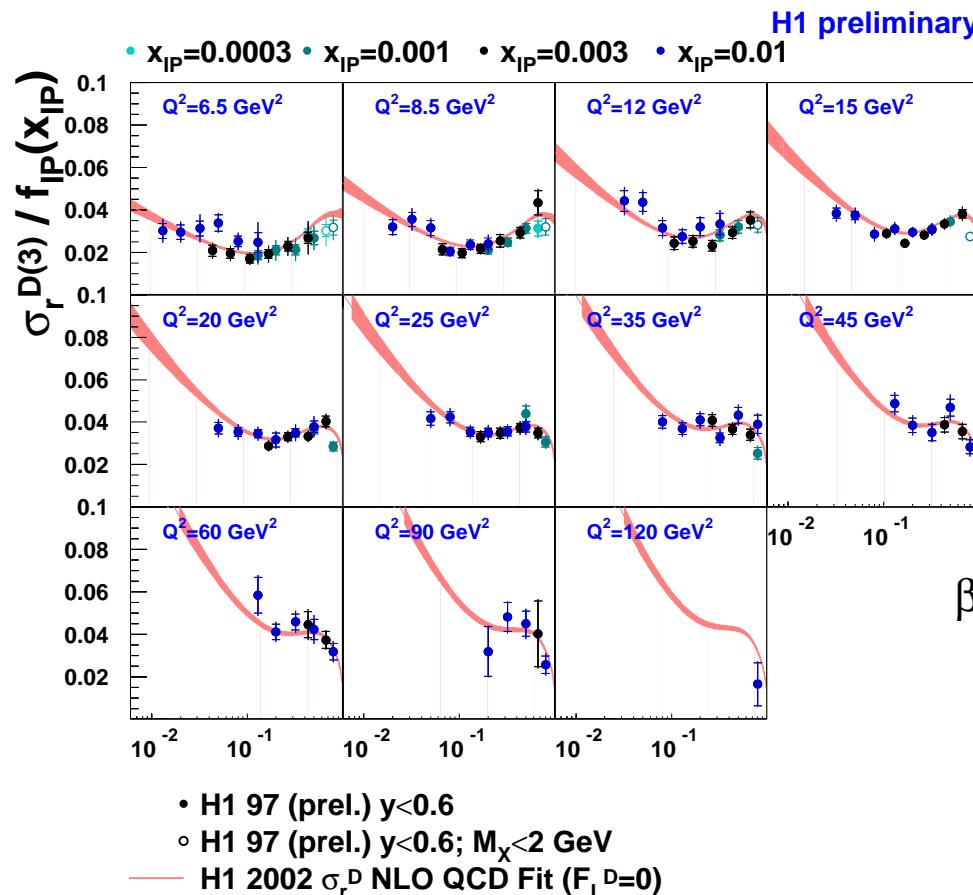


Different behaviour than for proton !

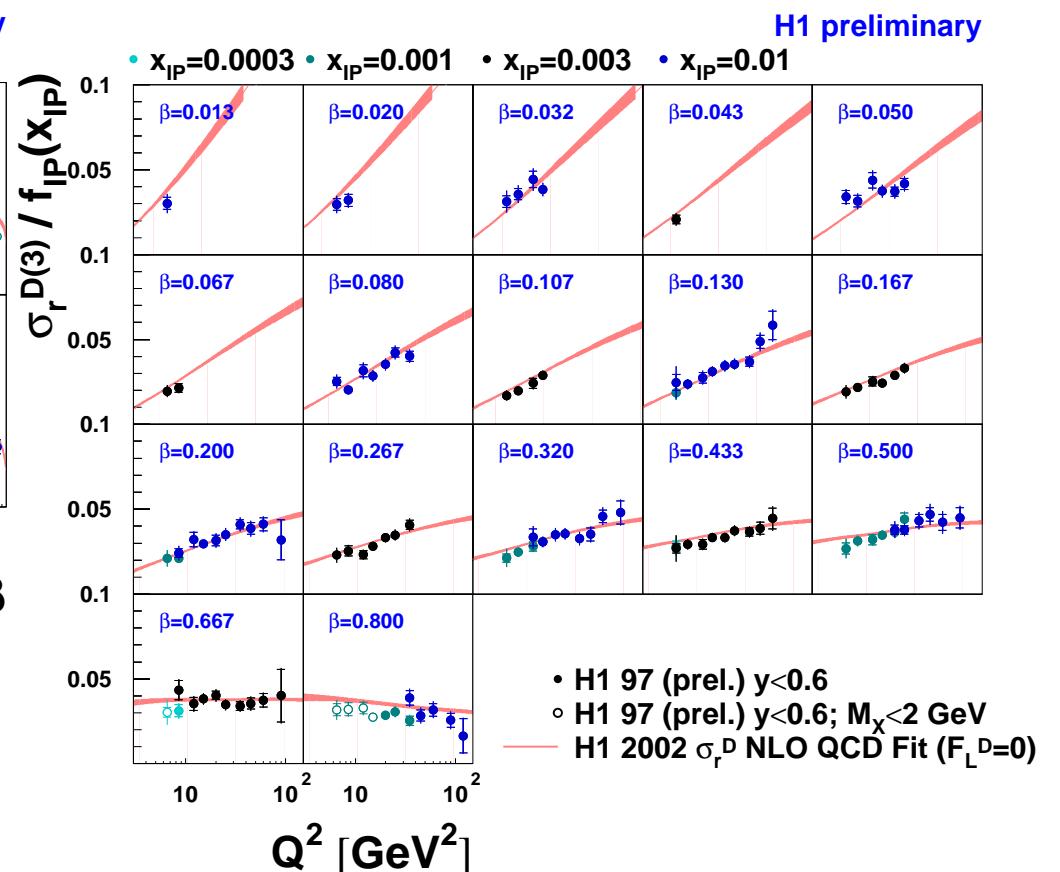
Taking out the x_{IP} dependence ...

Data divided by flux factor $f_{IP}(x_{IP})$

β dependence at fixed Q^2 :



Q^2 dependence at fixed β :



$$\beta \text{ dep.: } F_2^D = \sum e_i^2 (\mathbf{q}_i + \bar{\mathbf{q}}_i)$$

Scaling violations: gluon

Data consistent with Regge factorization

Modelling of $\sigma_r^{D(3)}$:

NLO DGLAP QCD Fit

- Shape of Q^2, β dep. of σ_r^D observed to be largely independent of x_{IP} :

$$\sigma_r^{D(4)}(x_{IP}, t, \beta, Q^2) = f_{IP}(x_{IP}, t) * \sigma_r^{D(2)}(\beta, Q^2)$$

- x_{IP} dependence conveniently parameterized as

$$f_{IP}(x_{IP}) = \int dt x_{IP}^{1-2\alpha_{IP}(t)} e^B t$$

using $\alpha_{IP}(0) = 1.173 \pm 0.018$ (determined from data)

- Small contribution from sub-leading exchange at large $x_{IP} > 0.01$ required

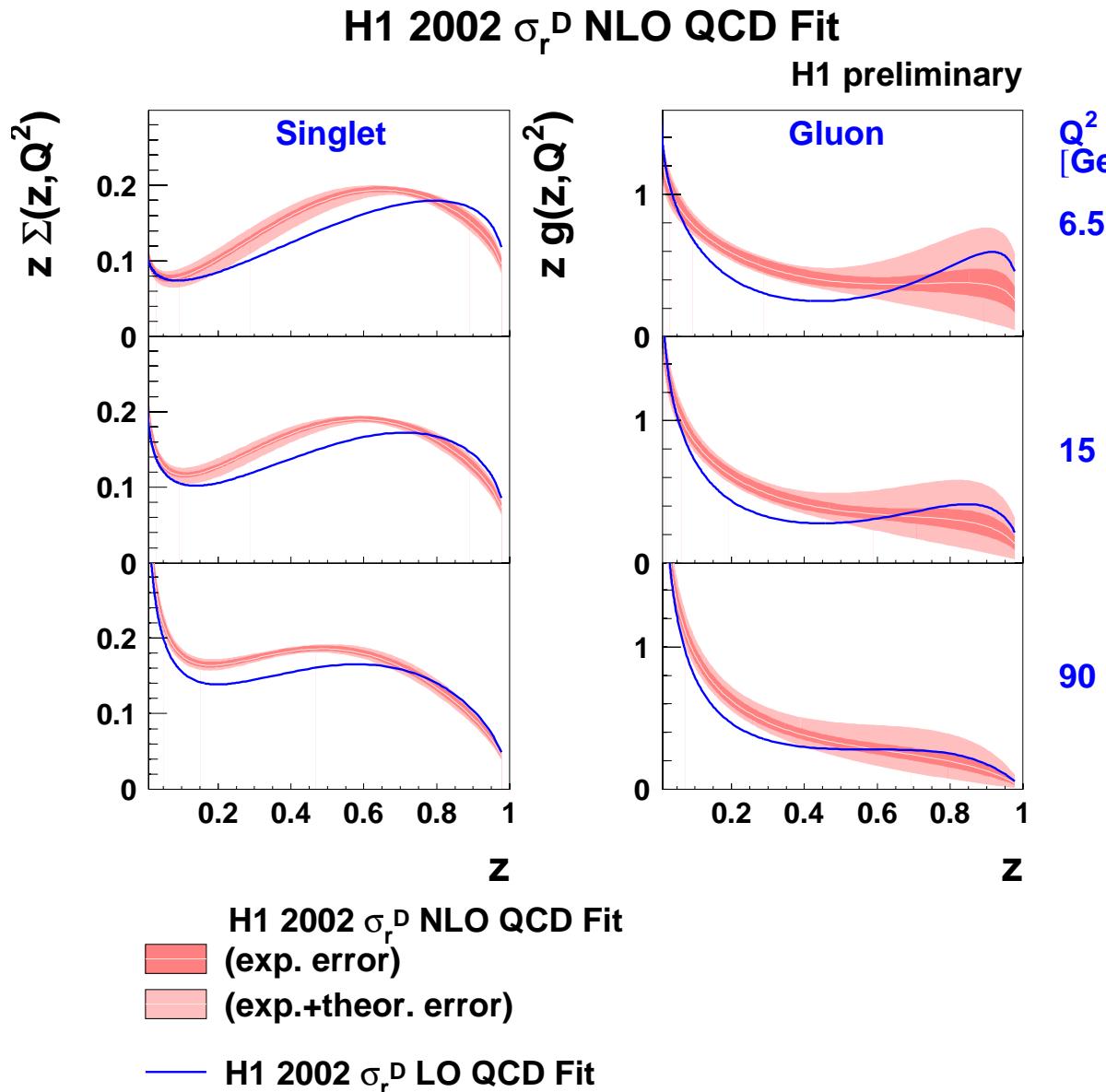
PDF parameterization:

- At starting scale $Q_0^2 = 3 \text{ GeV}^2$:
 - Singlet distribution $\Sigma(z, Q_0^2)$
 - Gluon distribution $g(z, Q_0^2)$
- Parameterization using unbiased, flexible functional form: Chebychev polynomials

Technique:

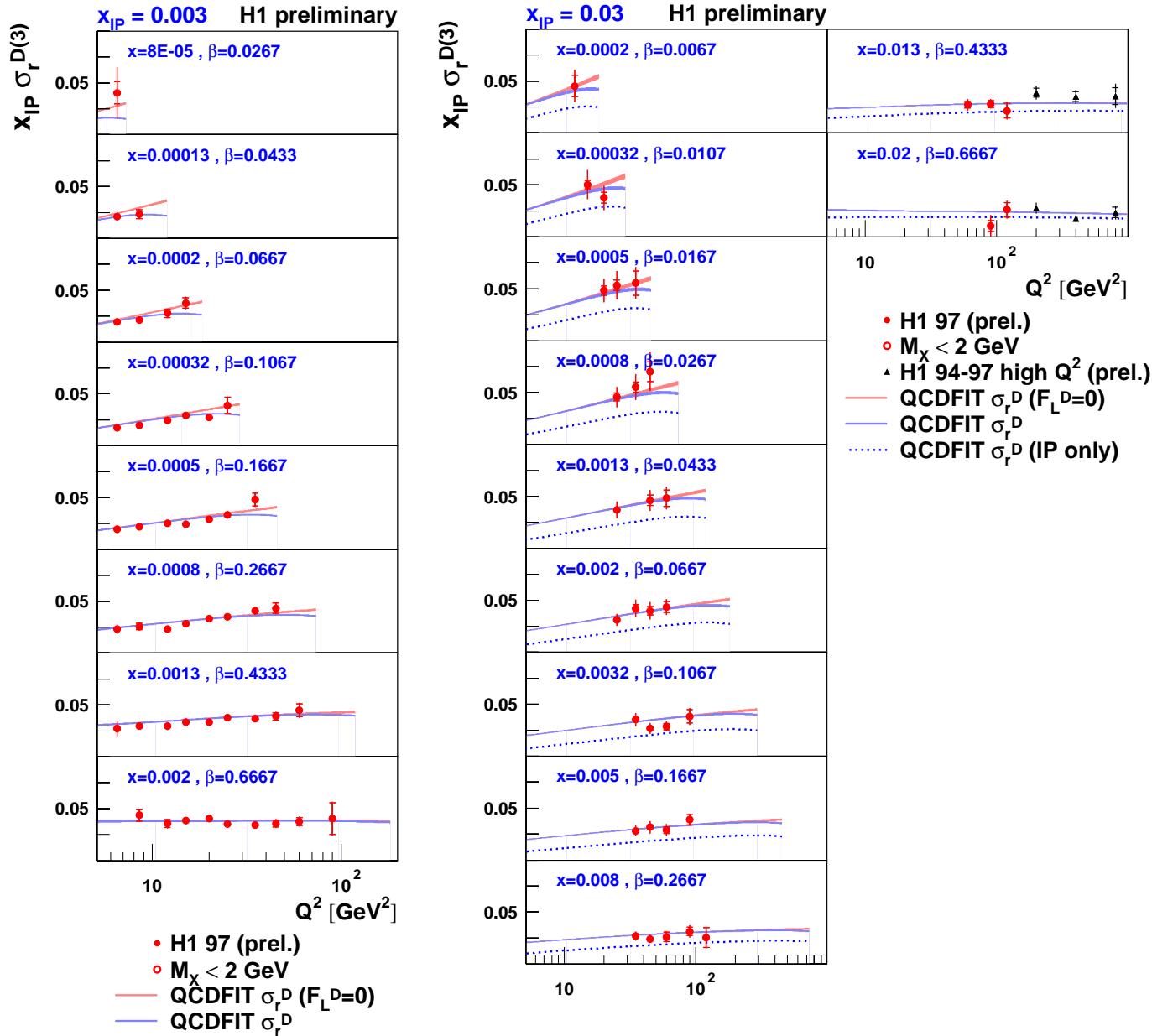
- Charm treatment in massive approach (BGF)
- Cut $M_X > 2 \text{ GeV}$ justifies leading twist analysis
- Full propagation of exp. and model systematic uncertainties !

Result of NLO fit



- pdfs extending to large fractional momenta z
- precise measurement of singlet distribution $\Sigma(z, Q^2)$
- hard gluon distribution, flat or rising towards $z \rightarrow 1$ (LO fit more peaked than central NLO fit)
- large uncertainty for $g(z, Q^2)$ at $z > 0.6$ (mainly related to model)

Comparison of NLO QCD fit with Data: Q^2 dep.



Two example x_{IP} bins

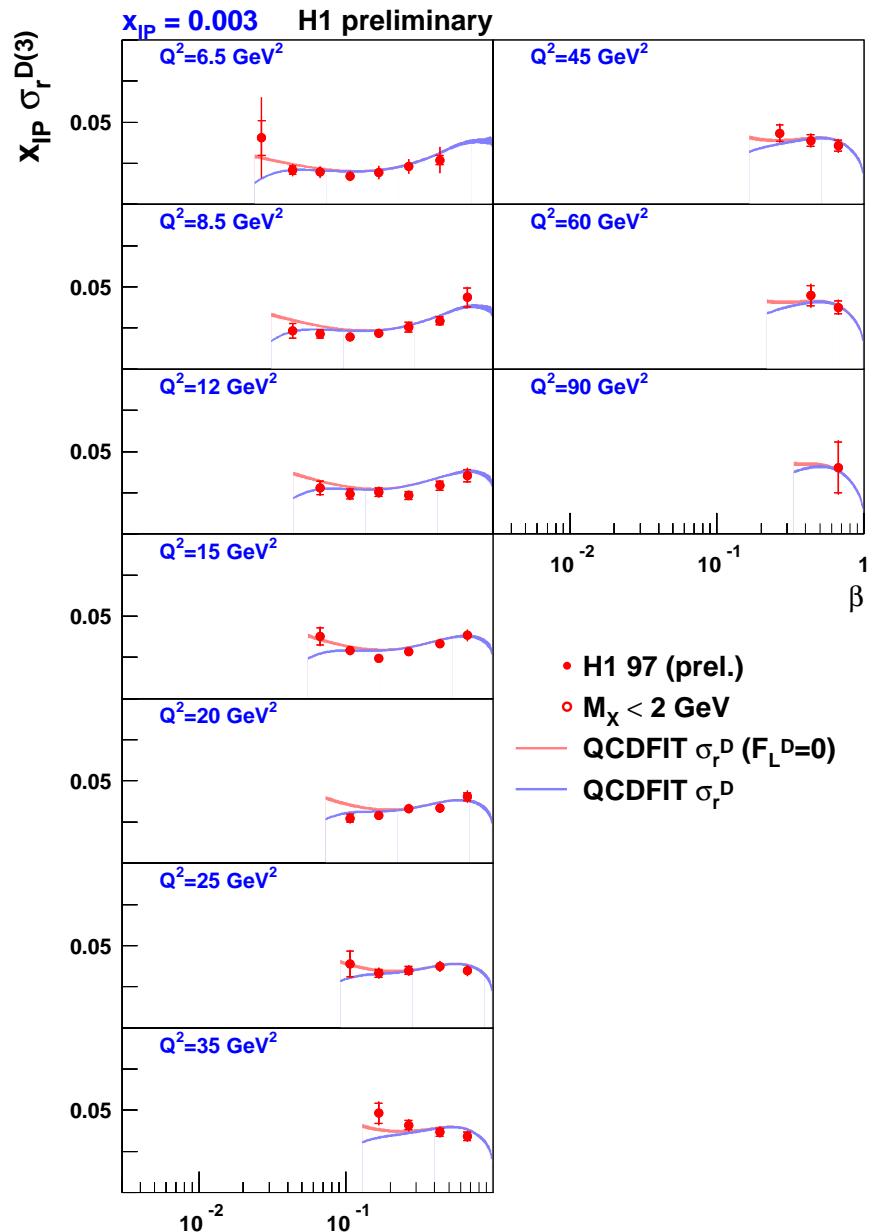
Q^2 scaling violations
well constrained
by data

Rising except at
highest β

Well reproduced
by QCD fit up to
 $Q^2 = 800 \text{ GeV}^2$

Sub-leading
contribution at
 $x_{IP} = 0.03$,
smaller than for
previous data

Comparison of NLO QCD fit with Data: β, x dep.



Example x_{IP} bin at 0.003:

Rising behaviour at $\beta \rightarrow 1$, low Q^2
reflected by $\Sigma(z, Q^2)$

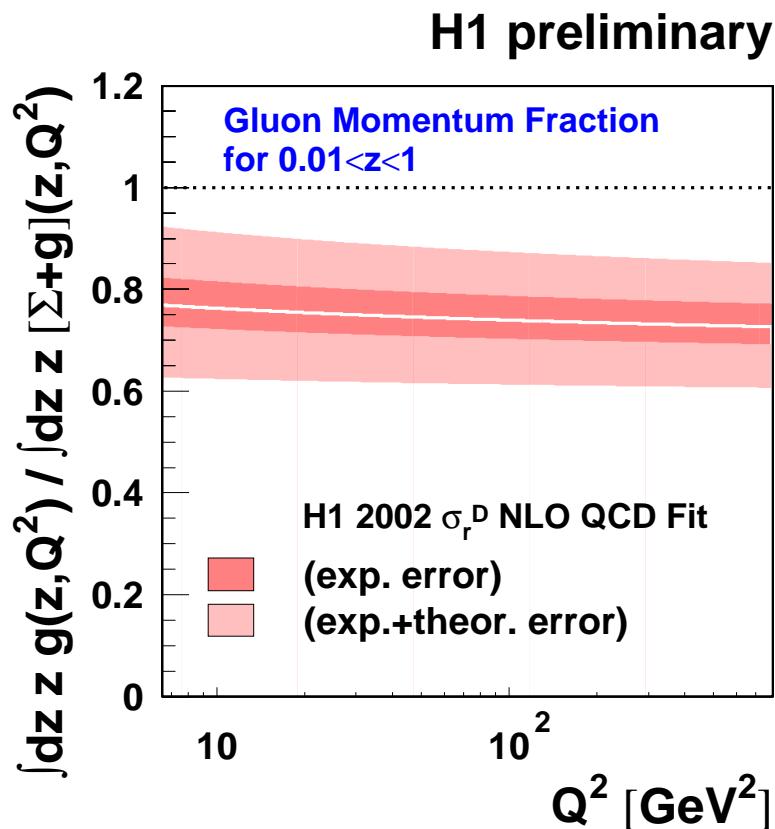
β dependence independent of x_{IP}

high $y \leftrightarrow$ low x or β at fixed x_{IP} :
Effect of F_L^D

presently no direct handle on
 F_L^D from data

Gluon Momentum Fraction

From NLO Fit:



- Integration of pdf's in measured range $0.01 < z < 1$
- Momentum fraction of colour singlet exchange carried by gluons **75%** for $6.5 < Q^2 < 800 \text{ GeV}^2$
- Fully consistent with results from previous H1 data

Comparison with H1 diffractive DIS final states

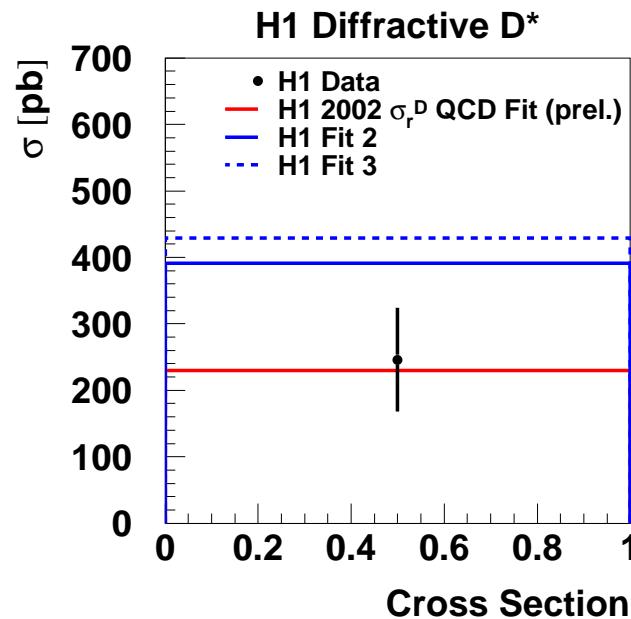
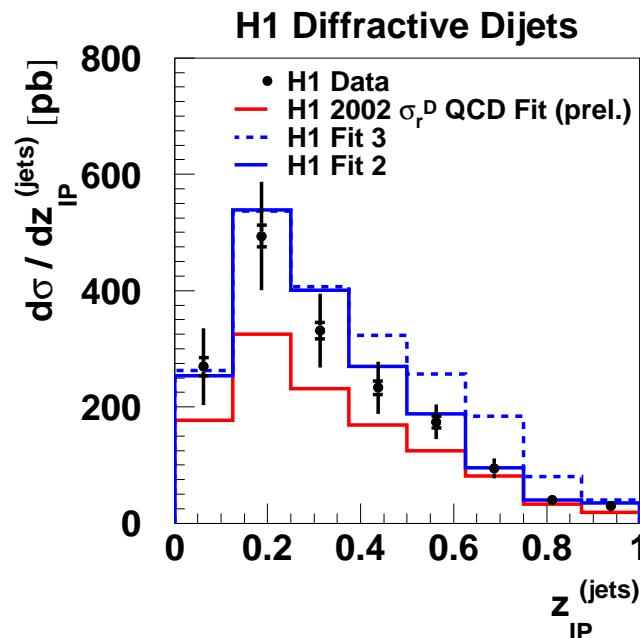
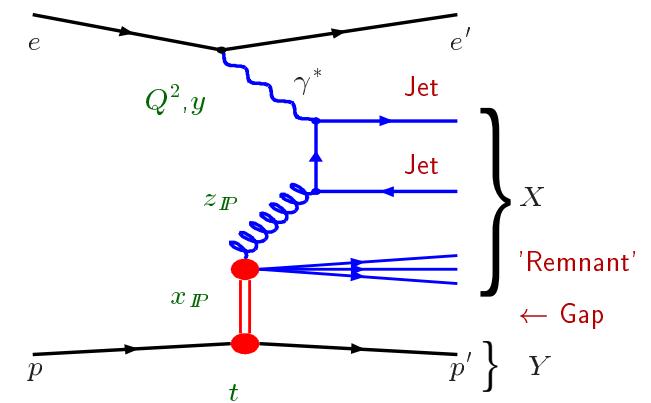
Use pdf's from LO fit to predict dijets / D^* cross sections in diffractive DIS as measured by H1:

Comparison based on MC model (RAPGAP)

$$\mu^2 = Q^2 + p_T^2 + m^2$$

Differential distributions remain well described

Normalization: pdf/NLO/scale uncertainty



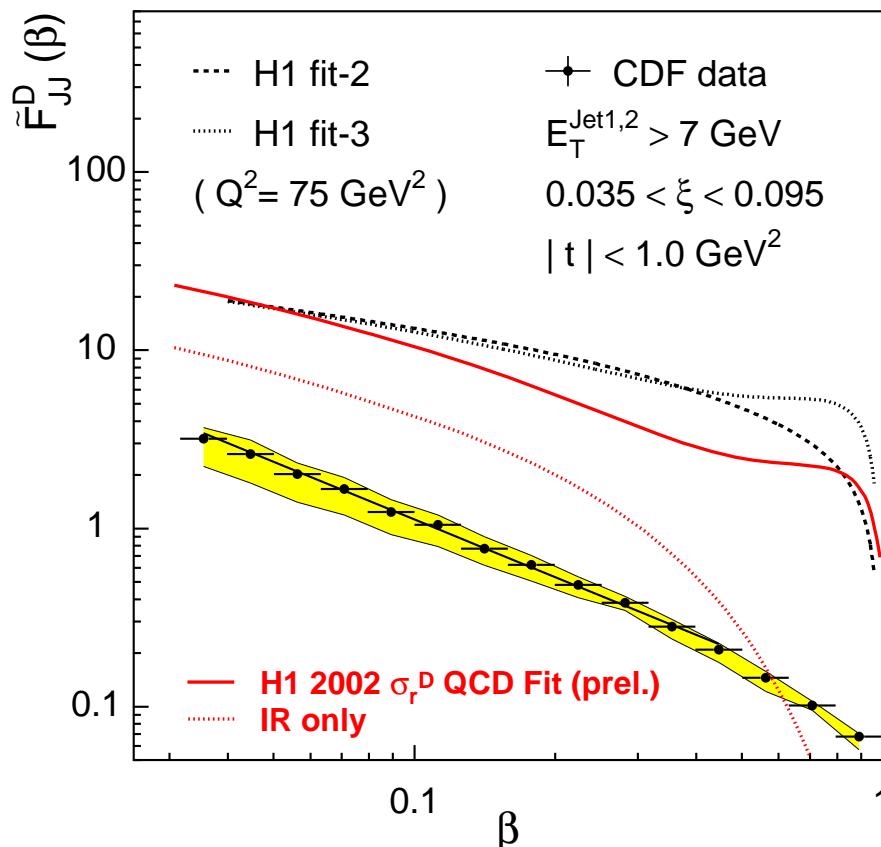
⇒ Consistent with QCD factorization !

Comparison with CDF diffractive Dijet cross sections

Dijet production with tagged leading anti-proton at TEVATRON:

Effective diffractive structure function \tilde{F}_{jj}^D :

$$\tilde{F}_{jj}^D(\beta) = \int dx_{IP} dt f(x_{IP}, t) \beta [g(\beta, Q^2) + \frac{4}{9}\Sigma(\beta, Q^2)] \quad (Q^2 = 75 \text{ GeV}^2)$$



- New fit confirms serious breakdown of factorization (gap survival, absorptive corrections)
- β dependence similar (except highest β)
- NOTE x_{IP} domain: 50% contribution from sub-leading exchange in this kinematic regime

Conclusions

- In diffractive DIS at HERA, the QCD structure of diffractive interactions (pomeron) can be probed using a virtual photon
- Experimental data have reached high precision
- Proof of QCD factorization in diffractive DIS provides firm theoretical basis
- Diffractive pdf's determined from F_2^D are dominated by large gluon contribution (75%) extending to large fractional momenta
- Comparisons with diffractive final states (dijets, charm): Consistent with QCD factorization
- Diffractive jets in $p\bar{p}$: Breakdown of factorization (gap survival, absorptive corrections)

Precision QCD in hard diffraction !