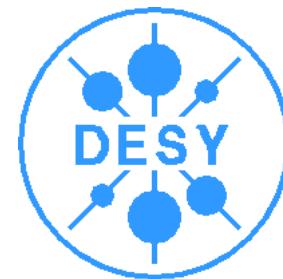


Hard diffraction at HERA: Recent Results and QCD Interpretation

Workshop on Diffraction and Glueballs at RHIC
May 2002, BNL / Upton (New York)

Frank-Peter Schilling (DESY)



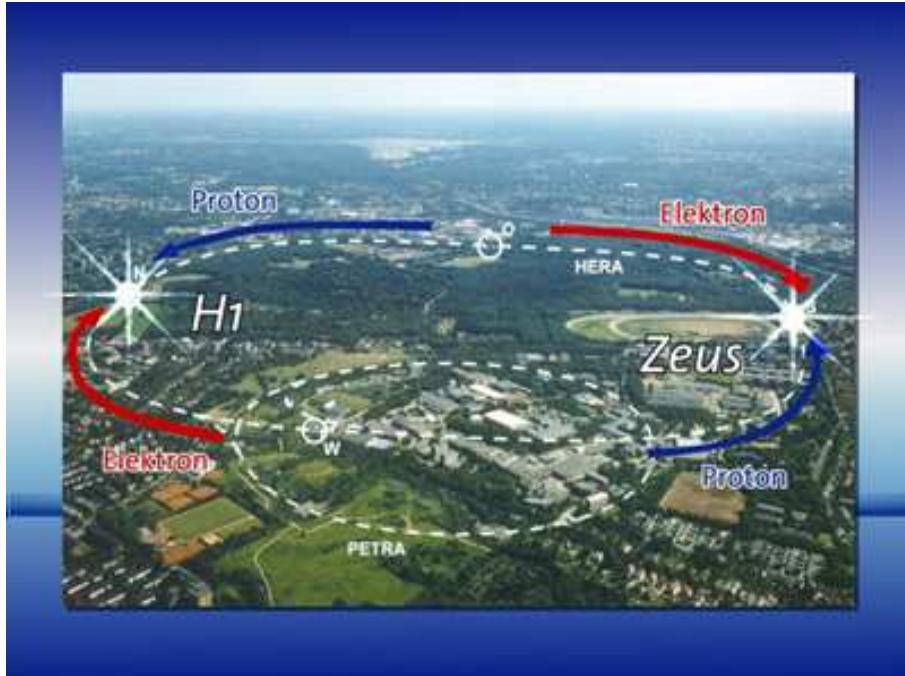
H1 Collaboration

- **Introduction**
- **New results in inclusive diffraction (F_2^D)**
- **New NLO DGLAP QCD fit and diffractive pdf's**
- **Diffractive final states (jets and charm)**
- **Dipole Models**
- **Summary and conclusions**

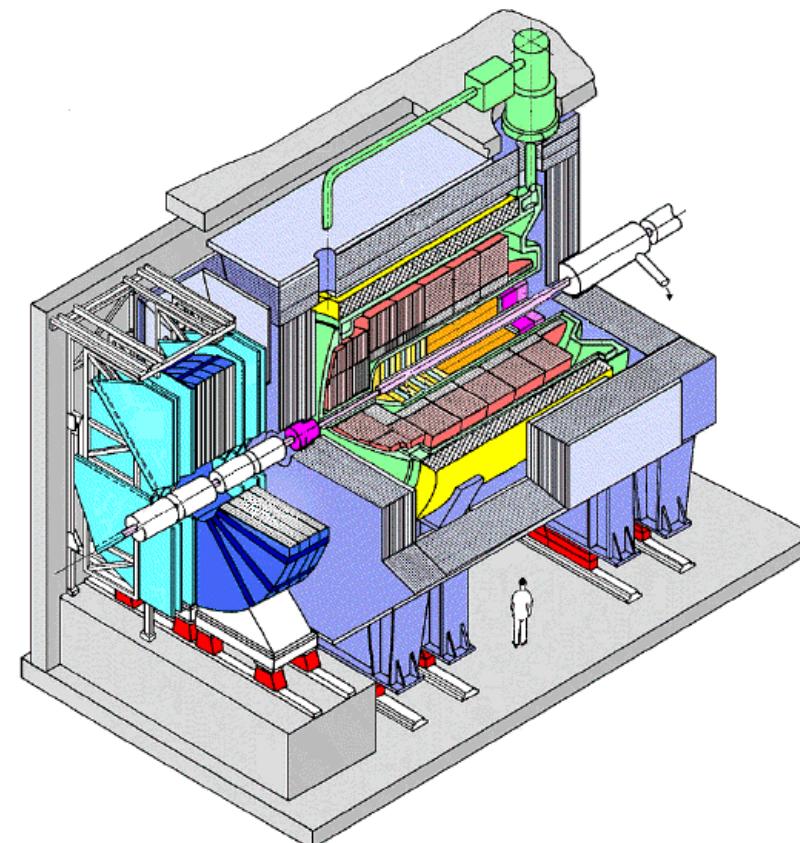
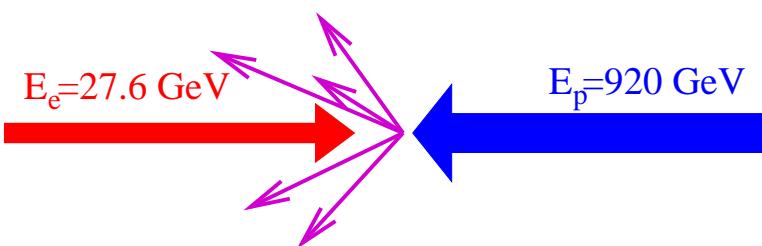
HERA and the H1 Detector

The HERA ep Collider at DESY / Hamburg: H1 and ZEUS:

ep collisions at $\sqrt{s} = 320$ GeV

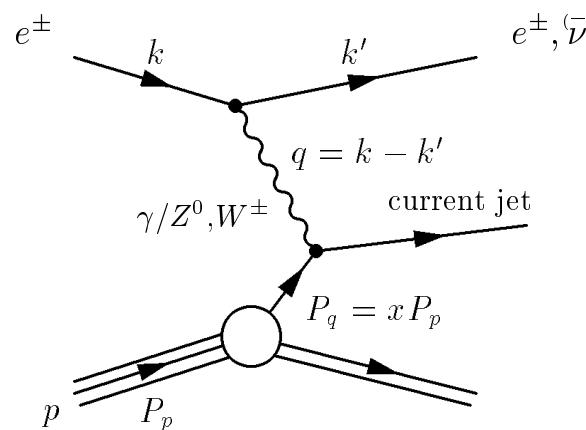


- HERA I: until 09/2000: $\mathcal{L} \sim 120 \text{ pb}^{-1}$
- Upgrade (lumi and dets.) just completed
- HERA II: $\mathcal{L} \sim 1 \text{ fb}^{-1}$ until 2006



~ 400 physicists

Preface: Deep-Inelastic Scattering (DIS) at HERA

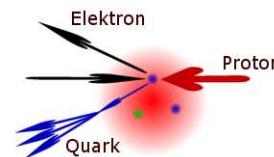
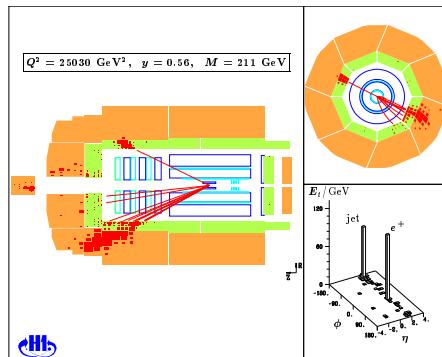


$$Q^2 = -q^2 = (k - k')^2$$

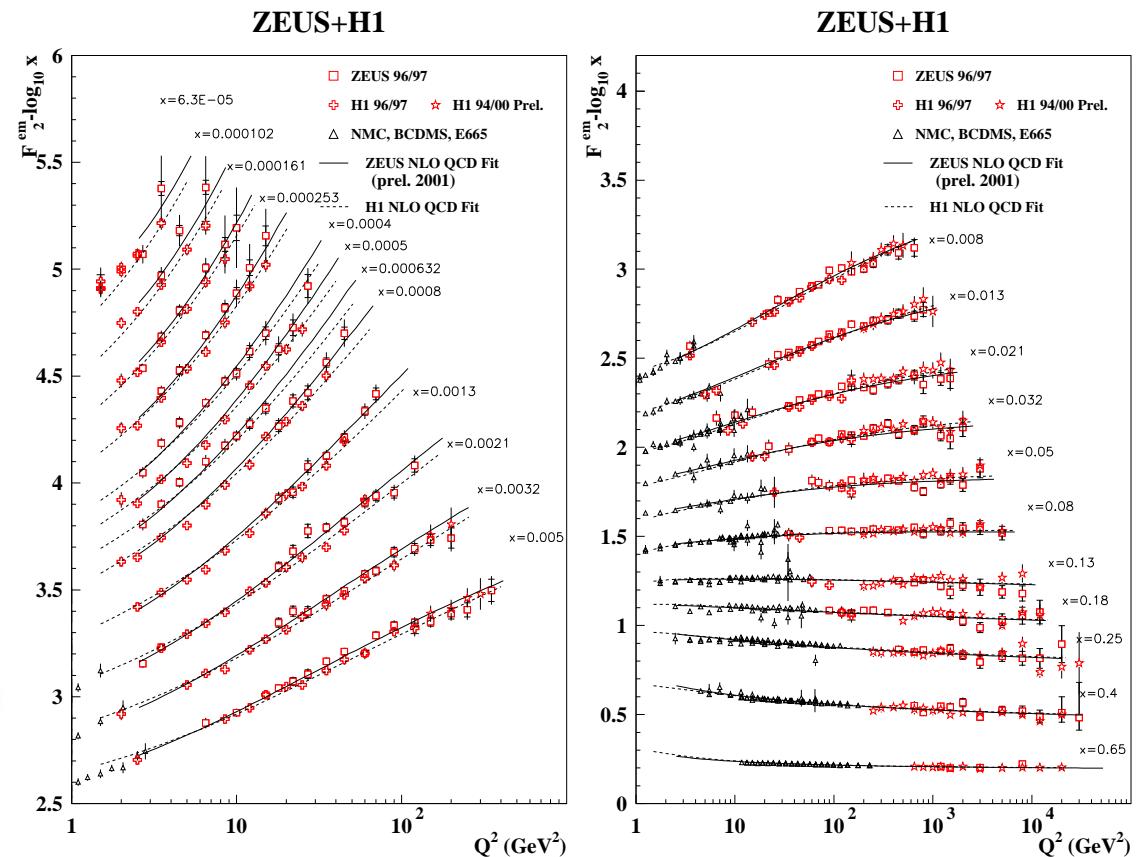
Photon virtuality

$$x = \frac{-q^2}{2P \cdot q} \quad (0 < x < 1)$$

Parton momentum frac. Bjorken-x



**Precise measurements of $F_2(x, Q^2)$:
Scaling violations**

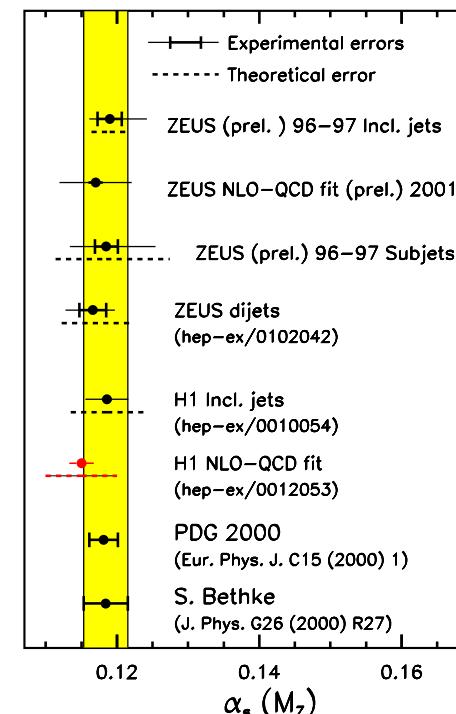
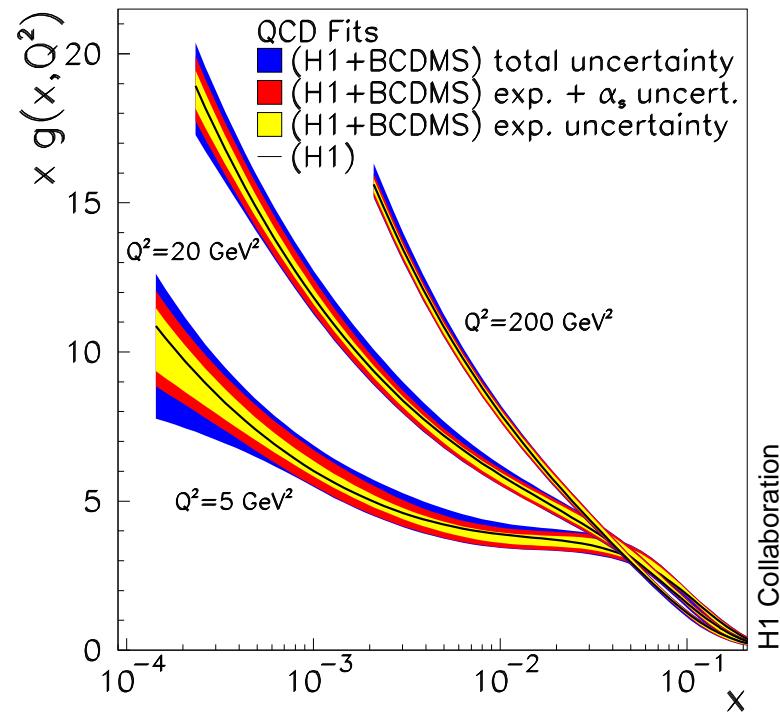


strong scaling violations at low x : $g \rightarrow q\bar{q}$

Preface: DIS at low x , gluons and α_s

QCD analysis of F_2 in framework of DGLAP:

$$\frac{dg(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[\sum_i q_i(z, Q^2) P_{gq} \left(\frac{x}{z} \right) + g(z, Q^2) P_{gg} \left(\frac{x}{z} \right) \right]$$



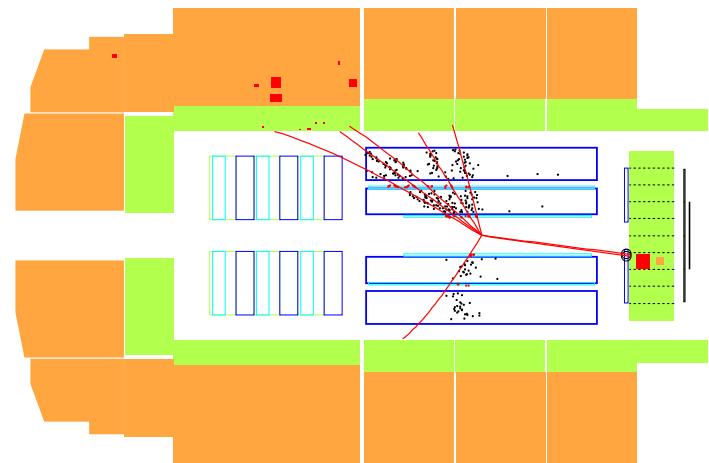
⇒ Precise determinations of the gluon distribution and α_s

Diffraction in DIS at HERA

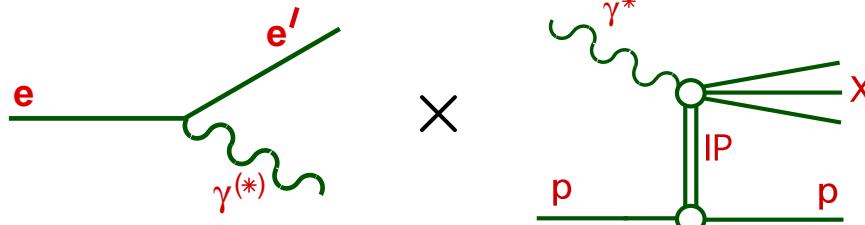
Early Observation at HERA:

10% of low- x DIS events are diffractive

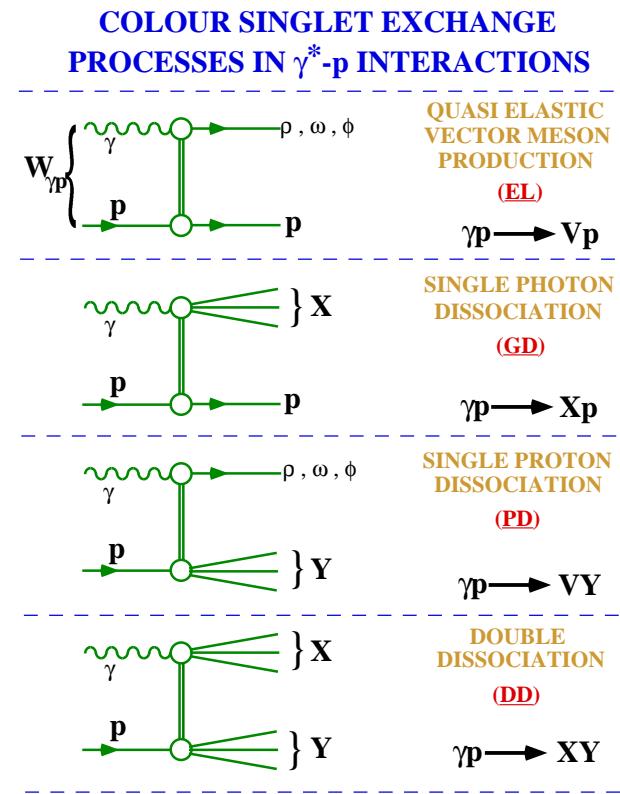
$$ep \rightarrow e' p' X$$



Can be viewed as diffractive $\gamma^* p$ interaction:



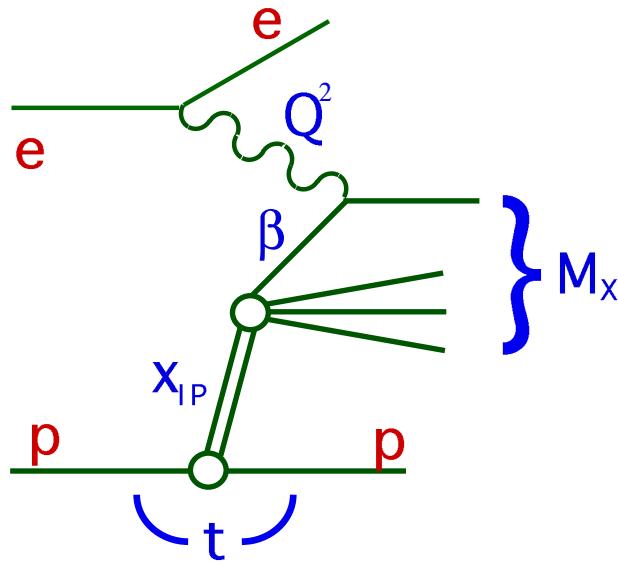
More generally: $\gamma^{(*)} p \rightarrow XY$



All can be measured by varying Q^2, W, t, M_X, M_Y

This talk mostly $\gamma^* p \rightarrow Xp$ (large Q^2 , small $|t|$)

Diffractive DIS



$$x_{IP} = \xi = \frac{Q^2 + M_X^2}{Q^2 + W^2} = x_{IP}/p$$

(momentum fraction of colour singlet exchange)

$$\beta = \frac{Q^2}{Q^2 + M_X^2} = x_q/IP$$

(fraction of exchange momentum carried by q coupling to γ^* , hence $x = x_{IP}\beta$)

$$t = (p - p')^2$$

(4-momentum transfer squared at p vertex)

Diffractive reduced cross section σ_r^D :

$$\frac{d^4\sigma}{dx_{IP} dt d\beta dQ^2} = \frac{4\pi\alpha^2}{\beta Q^4} \left(1 - y + \frac{y^2}{2}\right) \sigma_r^{D(4)}(x_{IP}, t, \beta, Q^2)$$

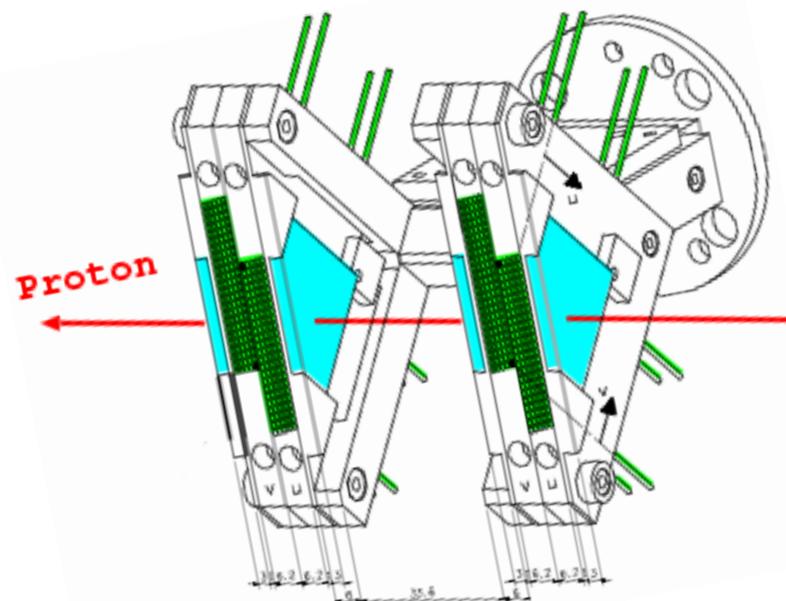
Structure functions F_2^D and F_L^D :

$$\sigma_r^{D(4)} = F_2^{D(4)} - \frac{y^2}{2(1-y+y^2/2)} F_L^{D(4)}$$

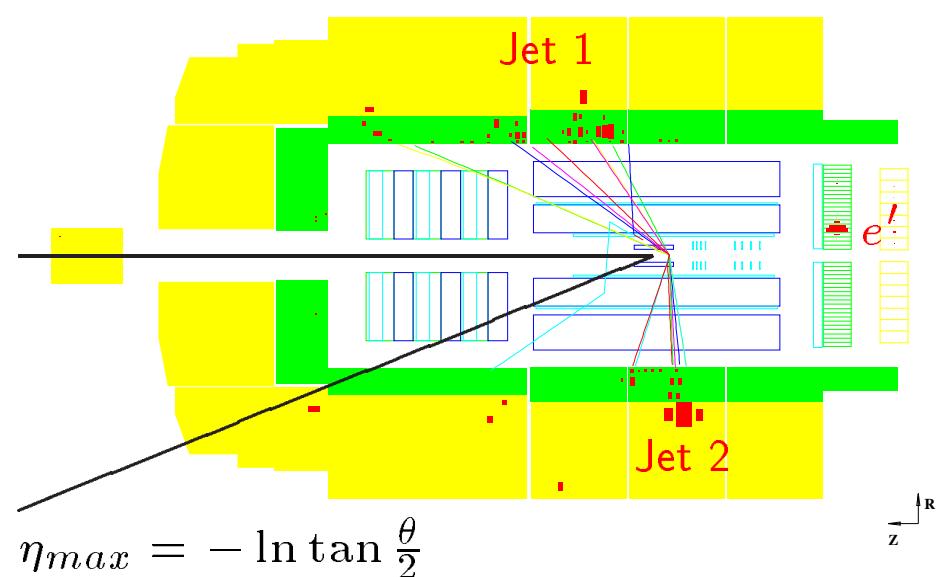
Integrated over t : $F_2^{D(3)} = \int dt F_2^{D(4)}$

Experimental Techniques

Forward Proton Spectrometer at $z = 65\ldots 90$ m



Rapidity Gap Selection in central detector



Measure leading proton

- Free of p dissociation bkgd.
- Measure t distribution
- low statistics (acceptance)

Require large rapidity gap

- $\Delta\eta$ large when $M_X \ll W$
- integrate over M_Y, t
- high statistics

Factorization Properties of F_2^D

QCD Factorization for diffractive DIS:

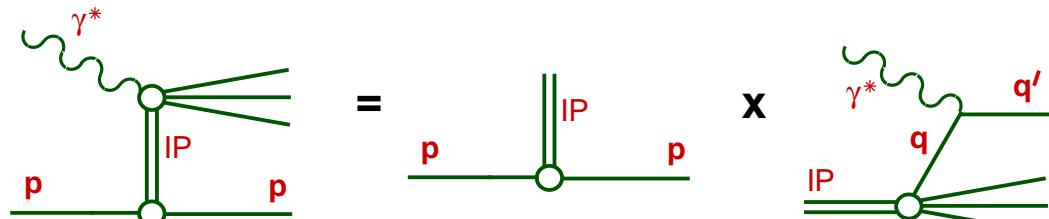
- **Diffractive parton distributions (Trentadue, Venezianio, Berera, Soper, Collins, ...):**

$$\frac{d^2\sigma(x, Q^2, x_{IP}, t) \gamma^* p \rightarrow p' X}{dx_{IP} dt} = \sum_i \int_x^{x_{IP}} d\xi \hat{\sigma}^{\gamma^* i}(x, Q^2, \xi) p_i^D(\xi, Q^2, x_{IP}, t)$$

- $\hat{\sigma}^{\gamma^* i}$ hard scattering part, as in incl. DIS
- p_i^D diffractive PDF's in proton, conditional probabilities, valid at fixed x_{IP}, t , obey DGLAP
- not proven for diffractive hadron-hadron scattering

Regge Factorization / resolved Pomeron model:

x_{IP}, t dependence factorizes out: Donnachie, Landshoff, Ingelman, Schlein, ...)

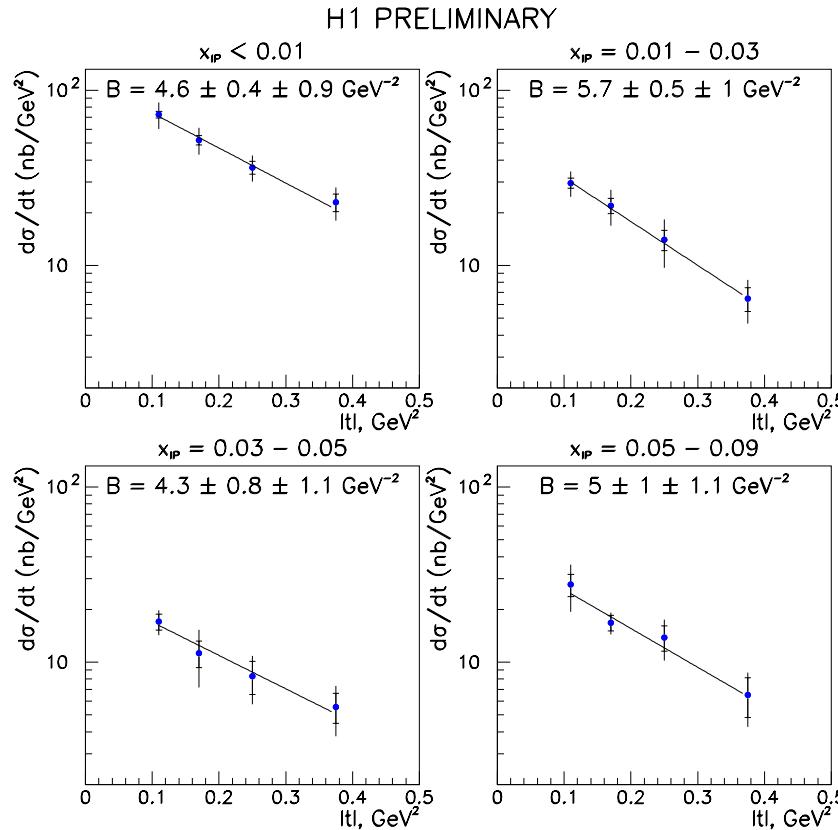


$$F_2^D(x_{IP}, t, \beta, Q^2) = f_{IP/p}(x_{IP}, t) F_2^{IP}(\beta, Q^2)$$

- additional assumption, no proof !
- consistent with present data if sub-leading IR included

New measurement of $F_2^{D(4)}$ using Roman Pots

Cross section differential in t at different x_{IP} :



Kinematic range:

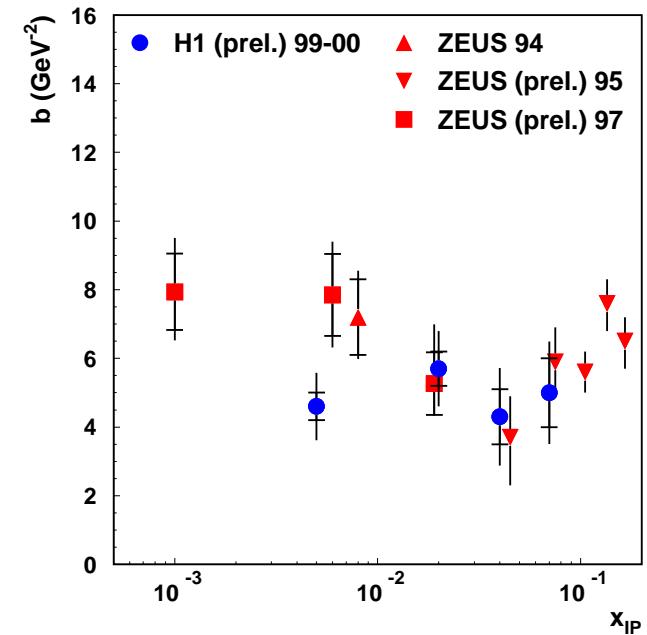
$$2 < Q^2 < 50 \text{ GeV}^2, 0.005 < \beta < 1$$

$$x_{IP} < 0.09, -0.45 < t < -0.08 \text{ GeV}^2$$

$$t \text{ dependence: } \frac{d\sigma}{dt} \sim e^{Bt}$$

In Regge theory expect shrinkage:

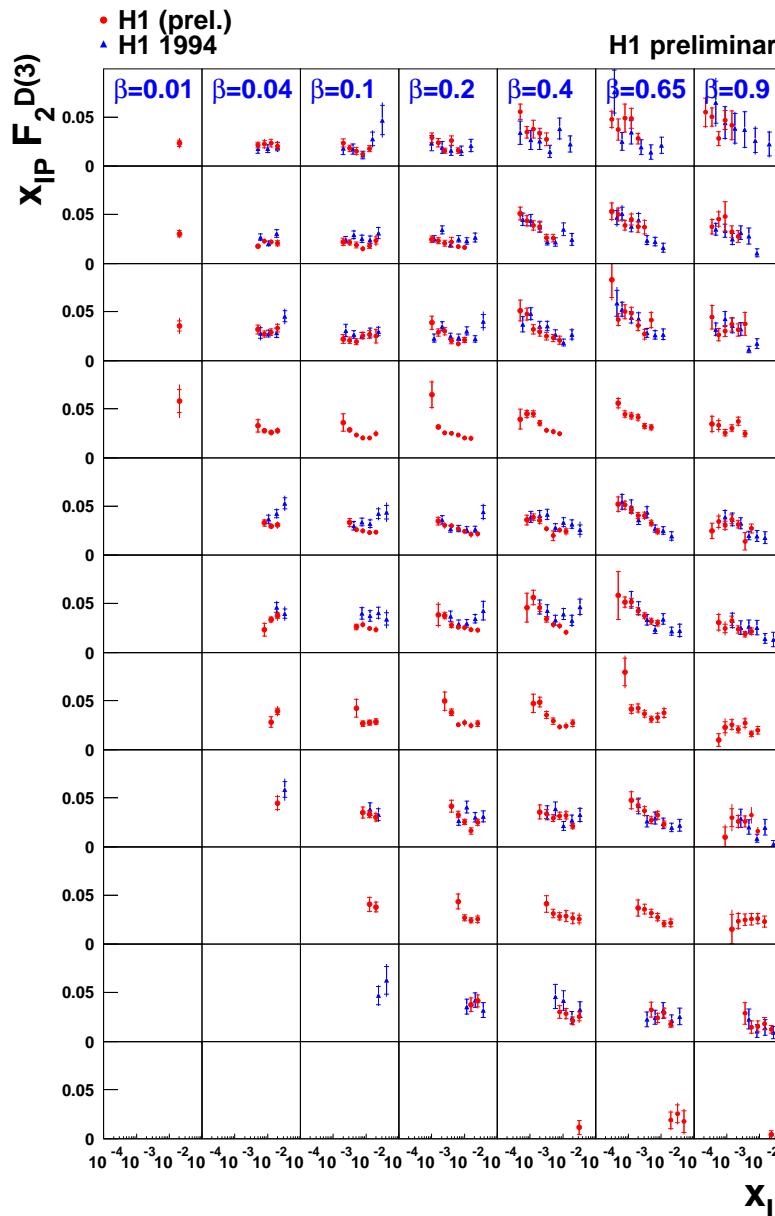
$$B = B_0 + 2\alpha' \ln \frac{1}{x_{IP}}$$



... data inconclusive so far

$$B = (5.0 \pm 0.3 \pm 0.8) \text{ GeV}^{-2}$$

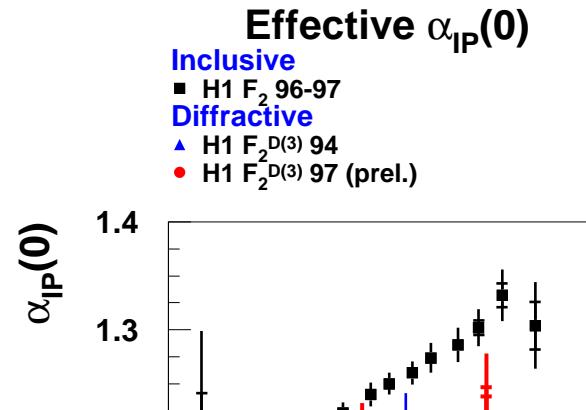
New $F_2^{D(3)}$ measurement (Rapidity gap)



5 times more statistics than previous data

Fit to x_{IP} dependence: effective $\alpha_{IP}(0)$

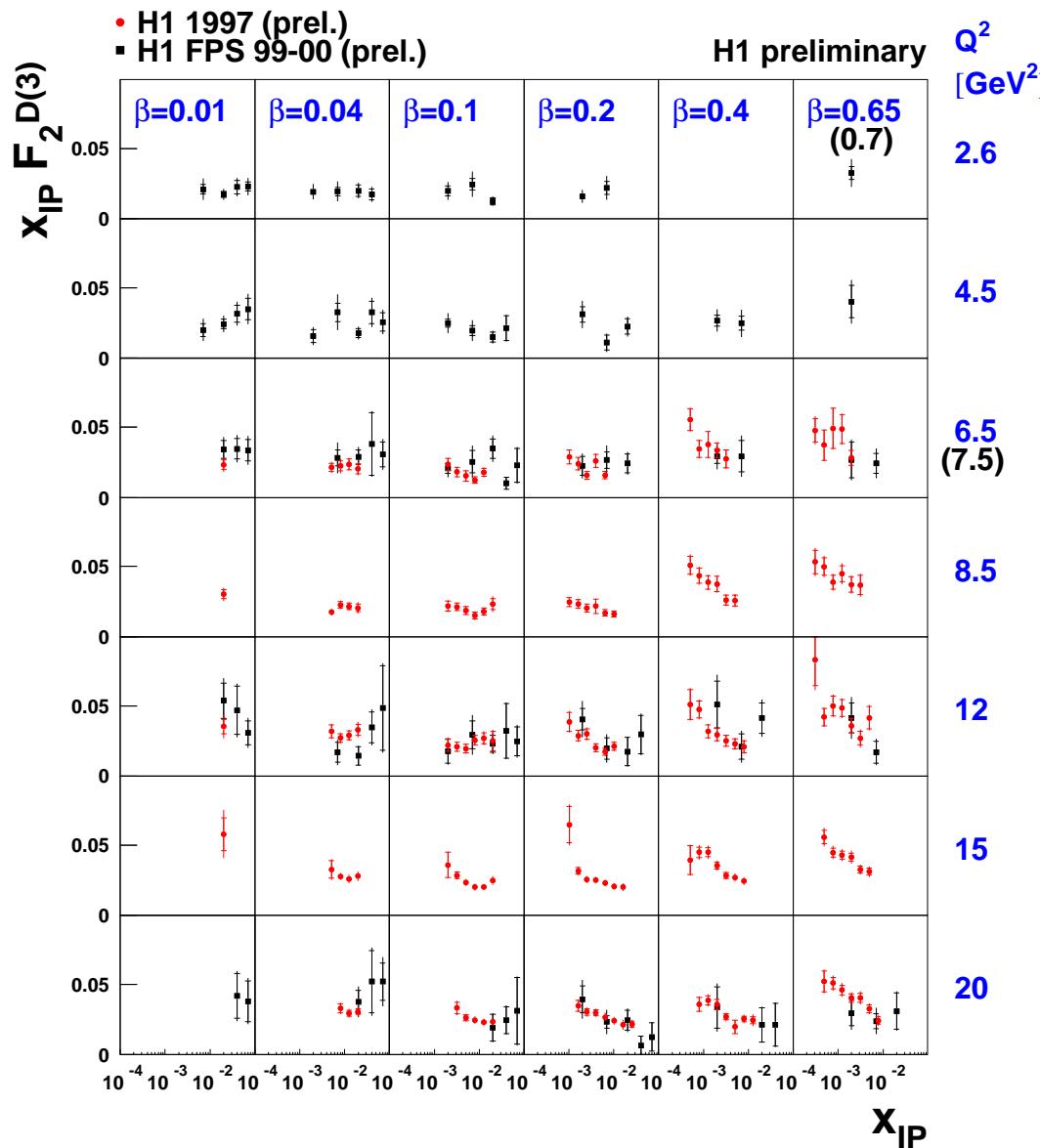
$$F_2^D(x_{IP}, \beta, Q^2) \sim B(\beta, Q^2) \left(\frac{1}{x_{IP}} \right)^{2\alpha_{IP}-1}$$



$$\alpha_{IP}(0) = 1.173 \pm 0.02 \pm 0.02^{+0.06}_{-0.03}$$

Growth with Q^2 slower in diffractive case?

Comparison Leading proton / rapidity gap data



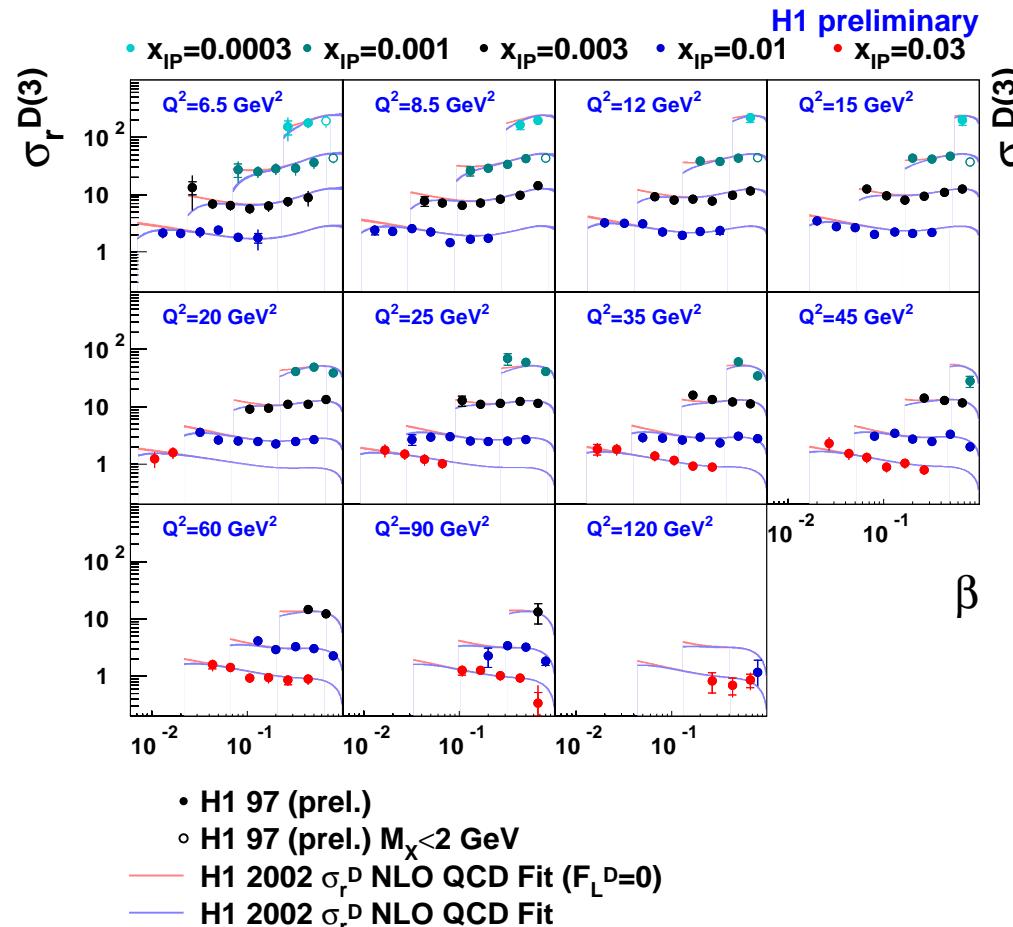
$$F_2^{D(3)} = \int dt F_2^{D(4)}$$

good agreement
between methods

justifies rapidity gap method!

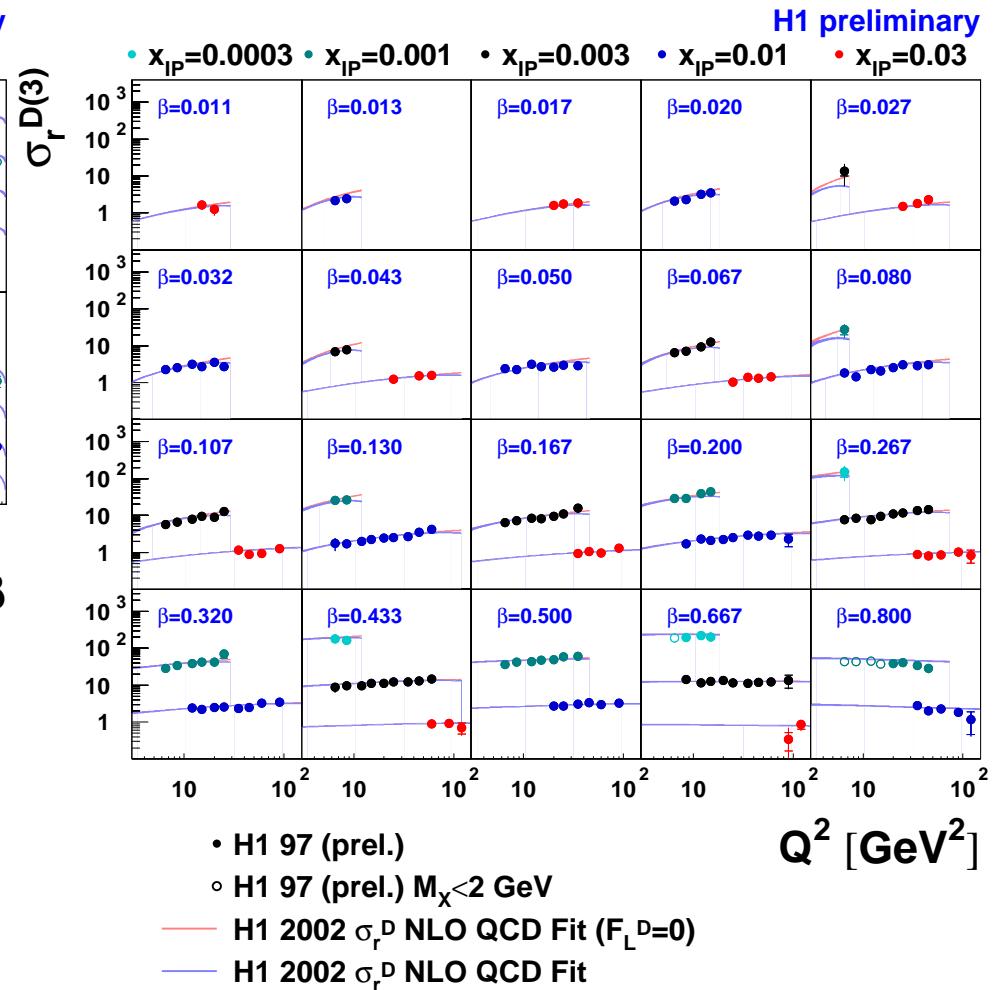
$F_2^{D(3)}$: β and Q^2 dependence overview

β dependence at fixed Q^2 :



sensitive to diffr. pdf's integrated over t

Q^2 dependence at fixed β :

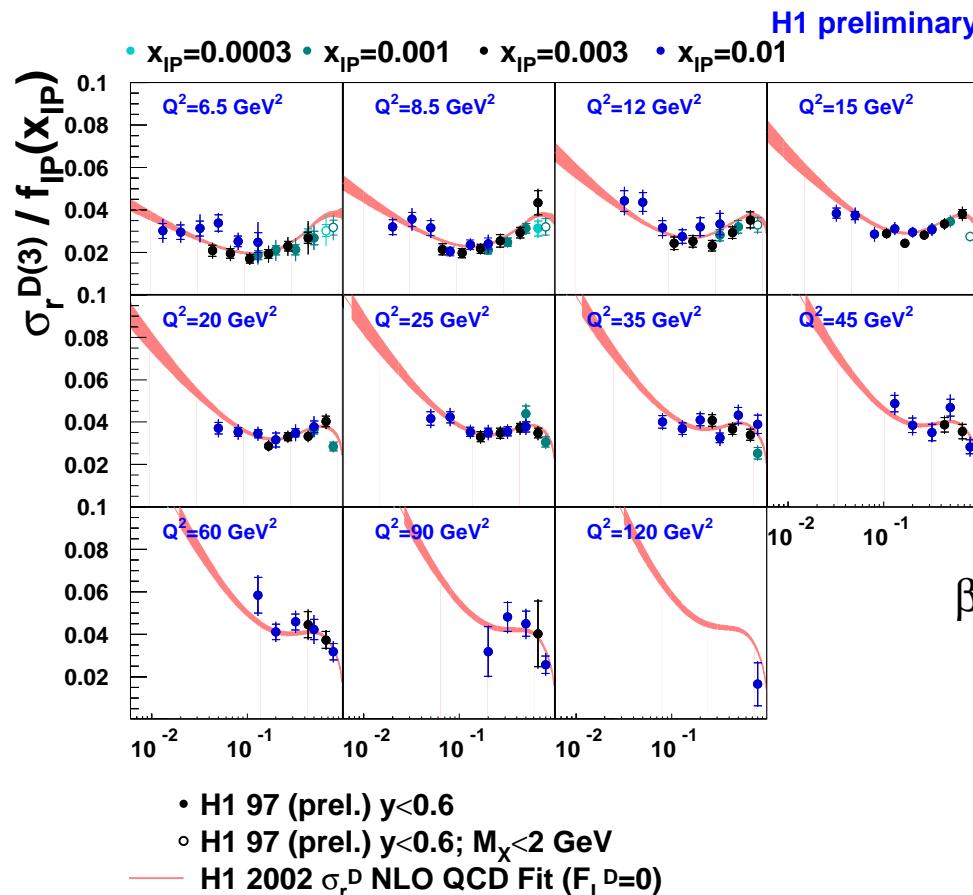


Different behaviour than for proton !

Taking out the x_{IP} dependence ...

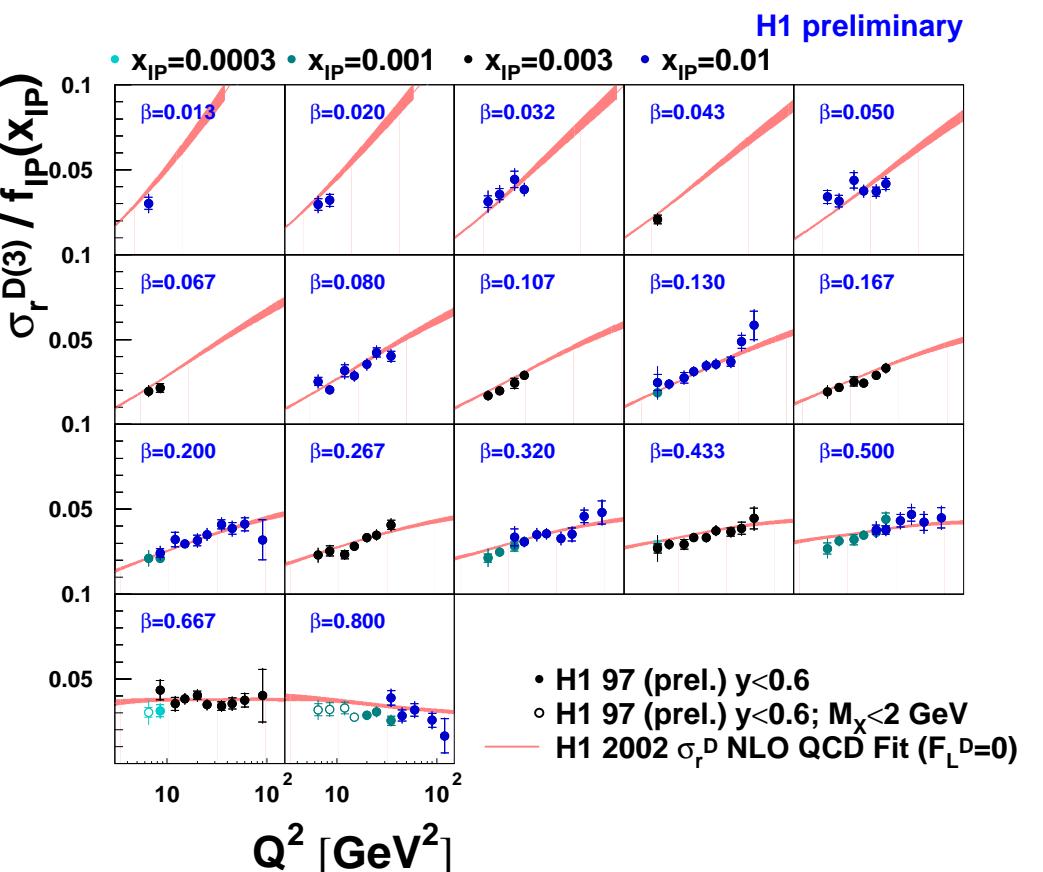
Data divided by flux factor $f_{IP}(x_{IP})$

β dependence at fixed Q^2 :



$$\beta \text{ dep.: } F_2^D = \sum e_i^2 (q_i + \bar{q}_i)$$

Q^2 dependence at fixed β :



Scaling violations: gluon

Data consistent with Regge factorization

Modelling of $\sigma_r^{D(3)}$:

NLO DGLAP QCD Fit

- Shape of Q^2, β dep. of σ_r^D observed to be largely independent of x_{IP} :

$$\sigma_r^{D(4)}(x_{IP}, t, \beta, Q^2) = f_{IP}(x_{IP}, t) * \sigma_r^{D(2)}(\beta, Q^2)$$

- x_{IP} dependence conveniently parameterized as

$$f_{IP}(x_{IP}) = \int dt x_{IP}^{1-2\alpha_{IP}(t)} e^{Bt}$$

using $\alpha_{IP}(0) = 1.173 \pm 0.018$ (determined from data)

- Small contribution from sub-leading exchange at large $x_{IP} > 0.01$ required

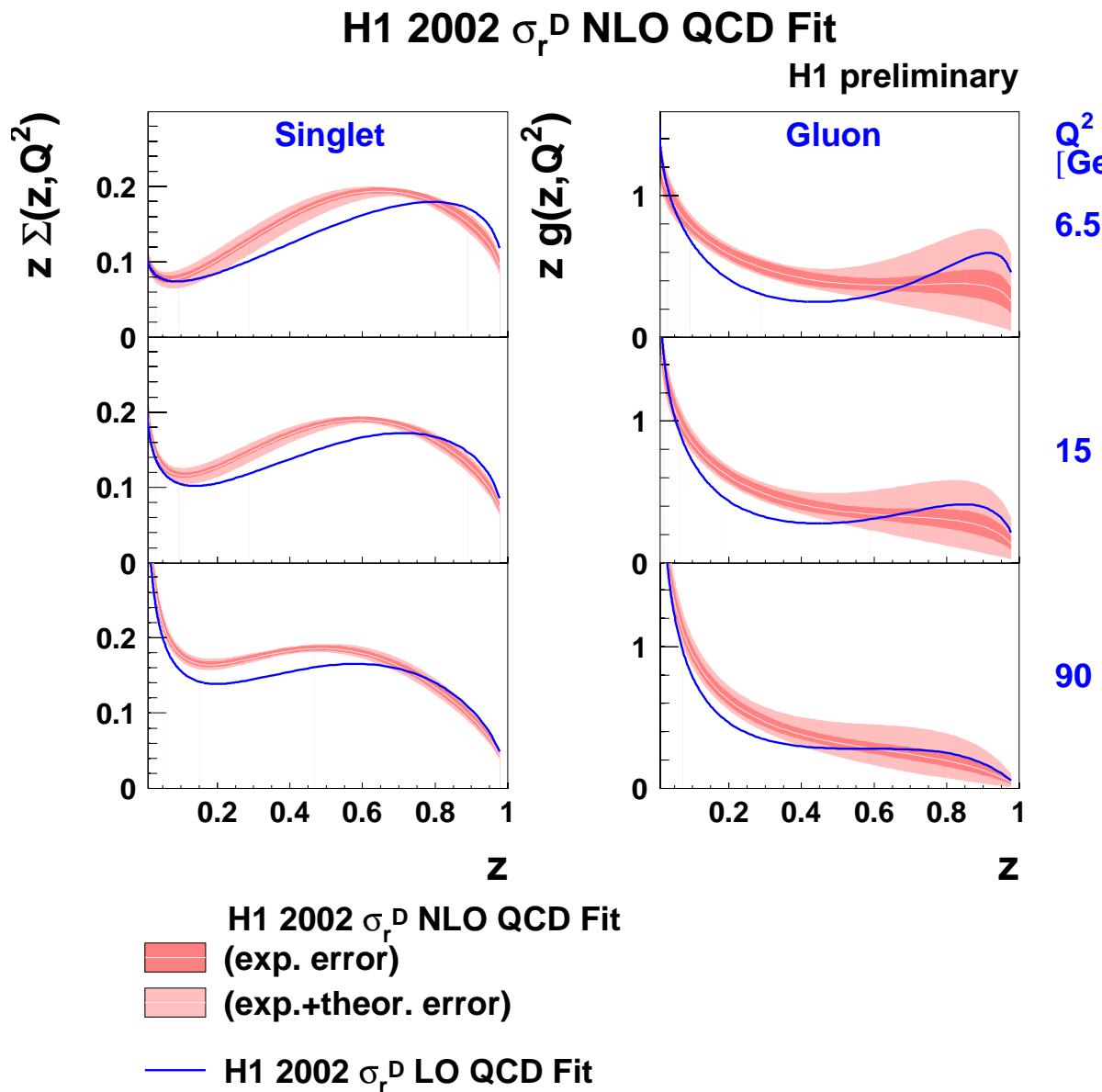
PDF parameterization:

- At starting scale $Q_0^2 = 3 \text{ GeV}^2$:
 - Singlet distribution $\Sigma(z, Q_0^2)$
 - Gluon distribution $g(z, Q_0^2)$
- Parameterization using unbiased, flexible functional form: Chebychev polynomials

Technique:

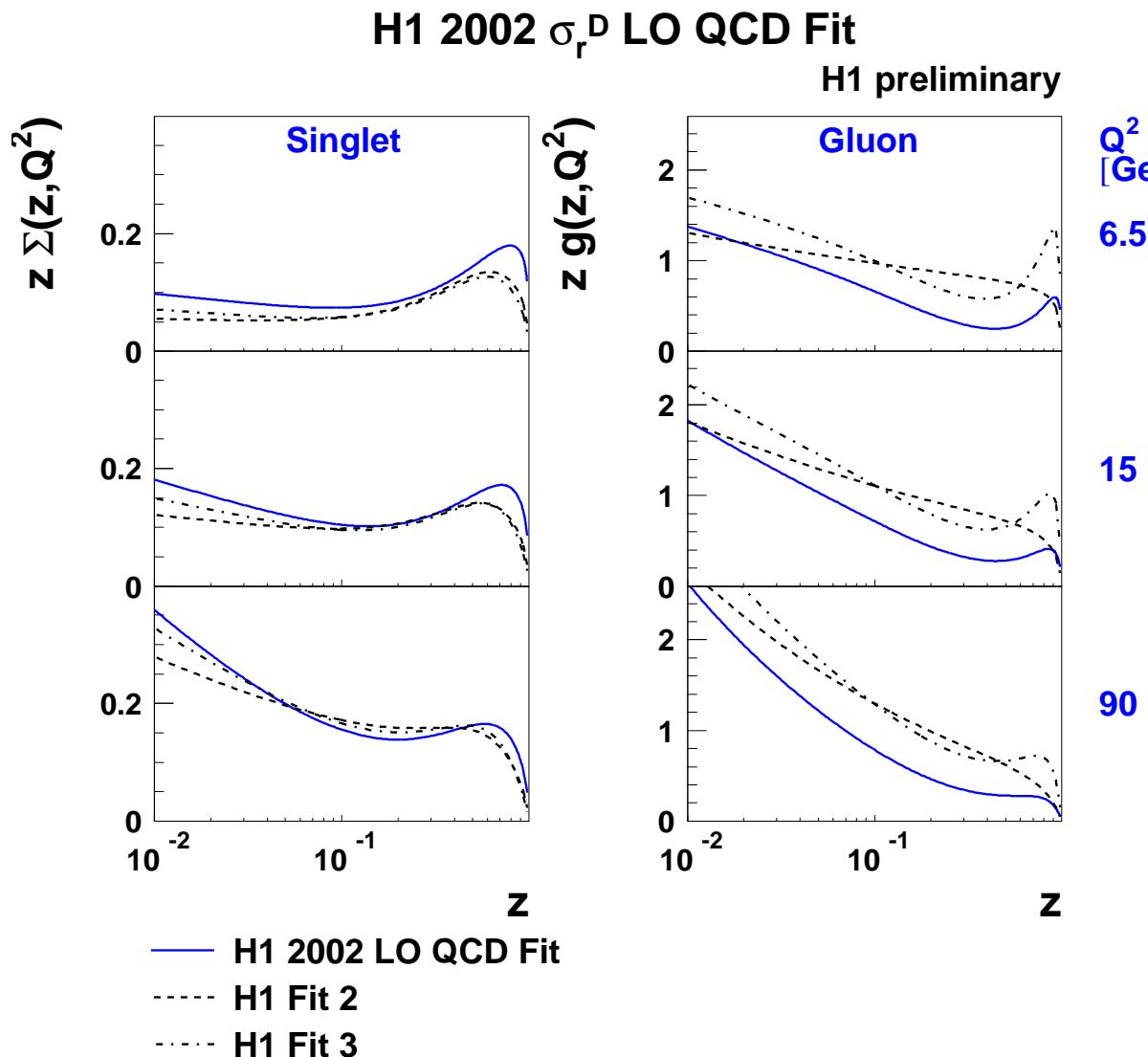
- Charm treatment in massive approach (BGF)
- Cut $M_X > 2 \text{ GeV}$ justifies leading twist analysis
- Full propagation of exp. and model systematic uncertainties !

Result of NLO fit



- pdfs extending to large fractional momenta z
- precise measurement of singlet distribution $\Sigma(z, Q^2)$
- hard gluon distribution, flat or rising towards $z \rightarrow 1$ (LO fit more peaked than central NLO fit)
- large uncertainty for $g(z, Q^2)$ at $z > 0.6$ (mainly related to model)

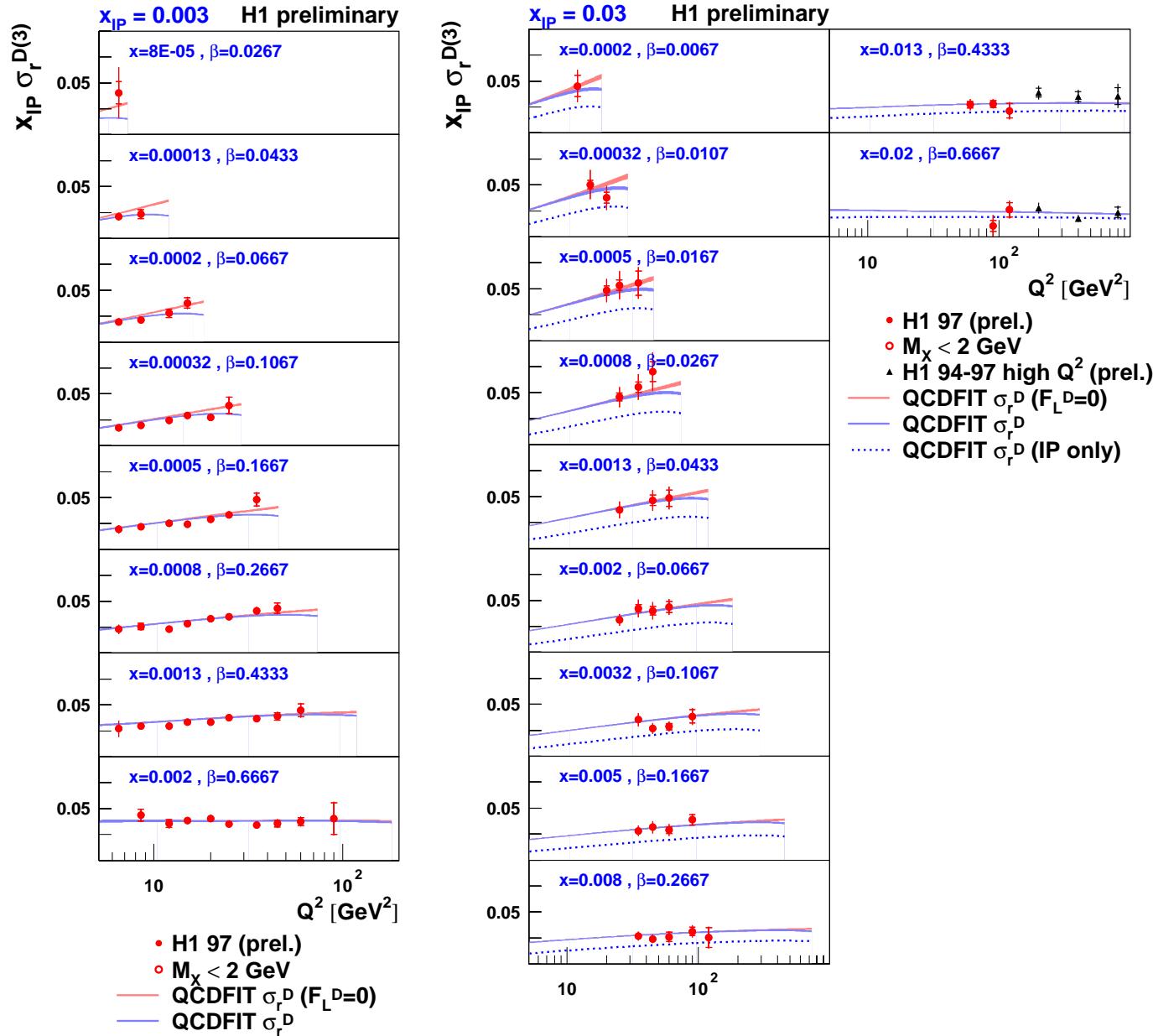
Leading order Fit and Comparison with previous H1 fits



- Comparisons with previous LO fits to 1994 data:
 - H1 Fit 2 (“flat gluon”)
 - H1 Fit 3 (“peaked gluon”)
- Reasonable agreement of $\Sigma(z, Q^2)$ for $z < 0.65$ (common fit range)
- Gluon normalization smaller by 20 – 30% at low z , 50% at high z

Agreement reasonable , taking errors of old and new fits into account

Comparison of NLO QCD fit with Data: Q^2 dep.



Two example x_{IP} bins

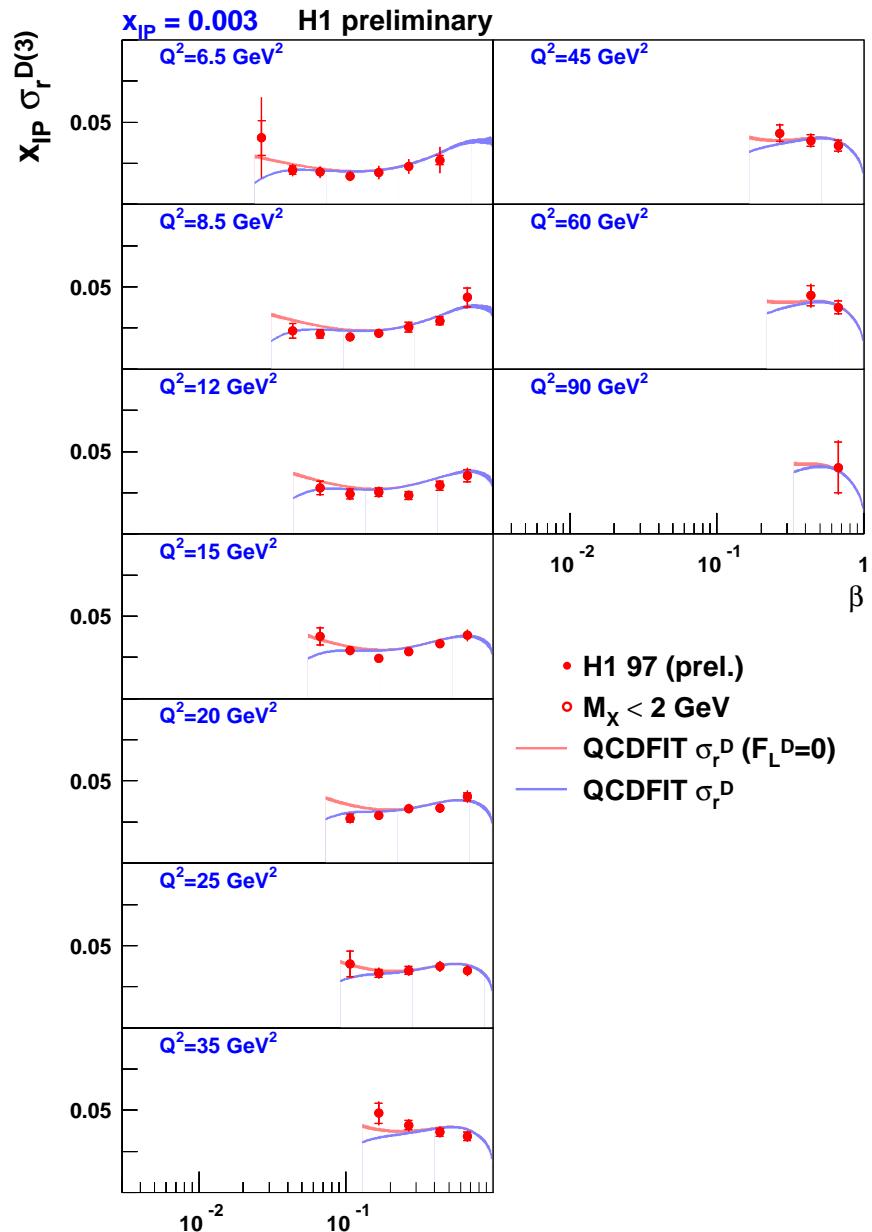
Q^2 scaling violations
well constrained
by data

Rising except at
highest β

Well reproduced
by QCD fit up to
 $Q^2 = 800 \text{ GeV}^2$

Sub-leading
contribution at
 $x_{IP} = 0.03$,
smaller than for
previous data

Comparison of NLO QCD fit with Data: β, x dep.



Example x_{IP} bin at 0.003:

Rising behaviour at $\beta \rightarrow 1$, low Q^2
reflected by $\Sigma(z, Q^2)$

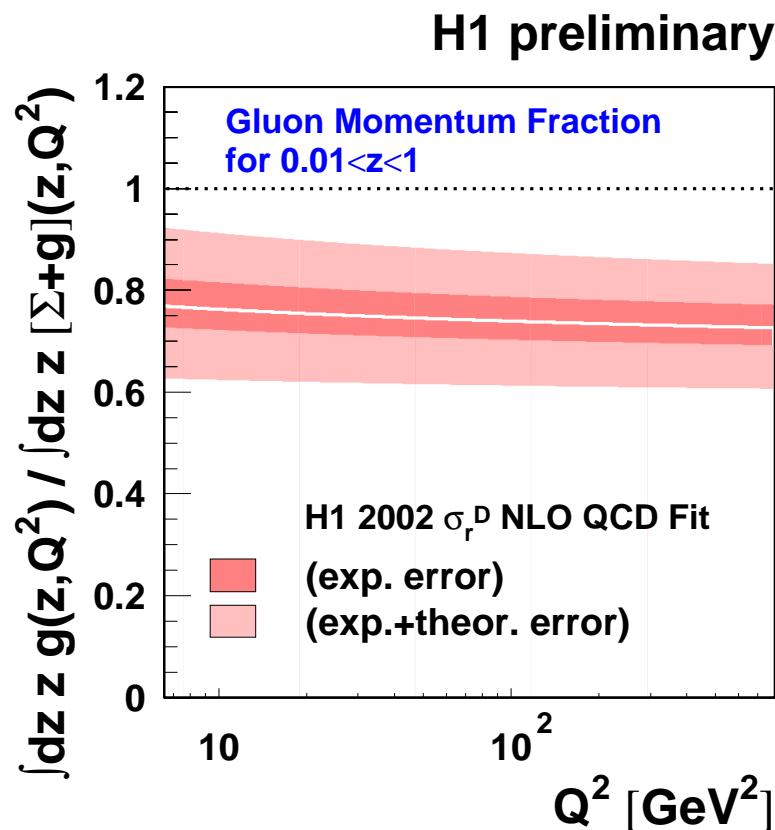
β dependence independent of x_{IP}

high $y \leftrightarrow$ low x or β at fixed x_{IP} :
Effect of F_L^D

presently no direct handle on
 F_L^D from data

Gluon Momentum Fraction

From NLO Fit:

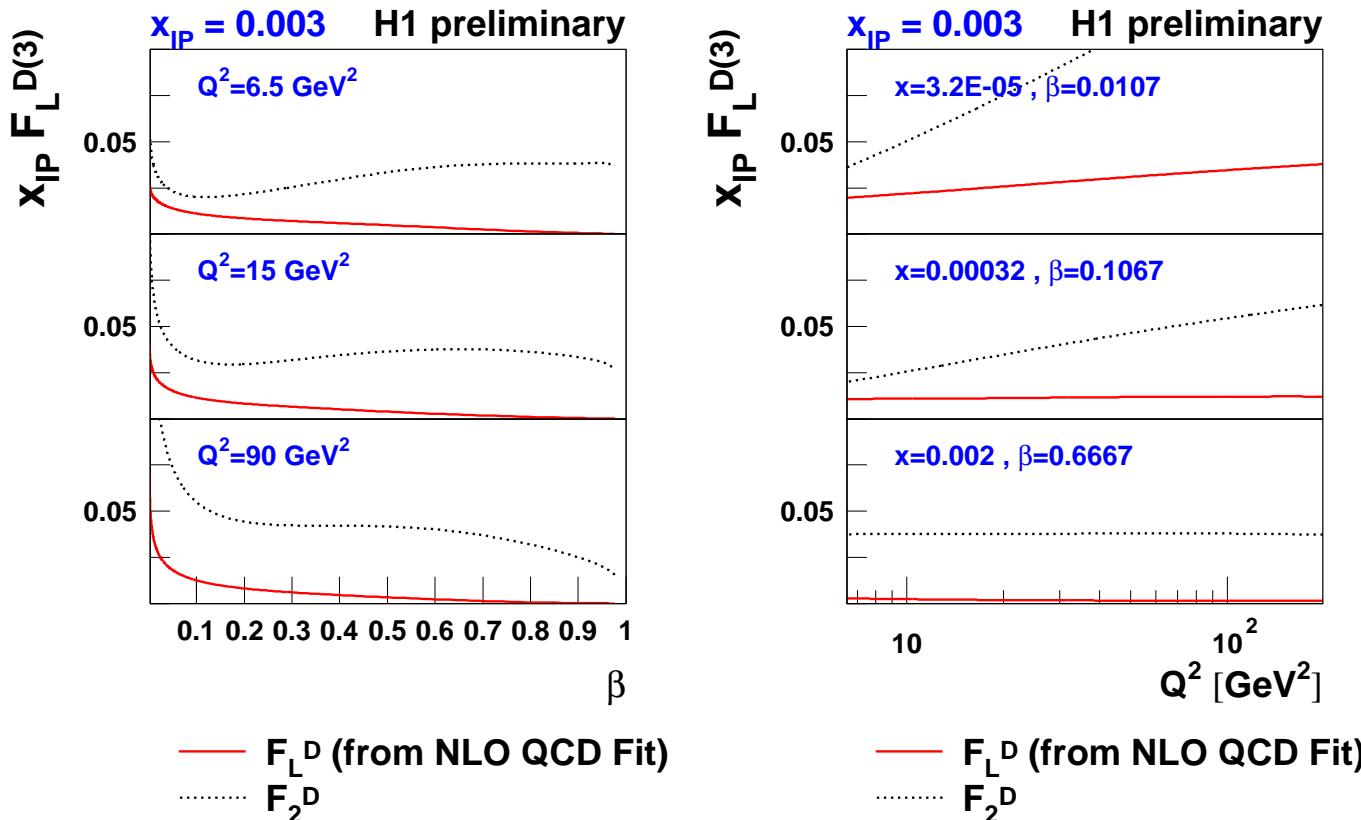


- Integration of pdf's in measured range $0.01 < z < 1$
- Momentum fraction of colour singlet exchange carried by gluons 75% for $6.5 < Q^2 < 800 \text{ GeV}^2$
- Fully consistent with results from previous H1 data

Longitudinal Structure Fraction F_L^D

At NLO QCD, the leading twist longitudinal structure function F_L^D is predicted:

$$F_L^D \sim \frac{\alpha_s}{2\pi} \left[C_q^L \otimes F_2^D + C_g^L \otimes \sum_i e_i^2 z g^D(z, Q^2) \right]$$



→ pert. F_L^D rel. large, in particular at low Q^2 , low β (due to large $g(z, Q^2)$)

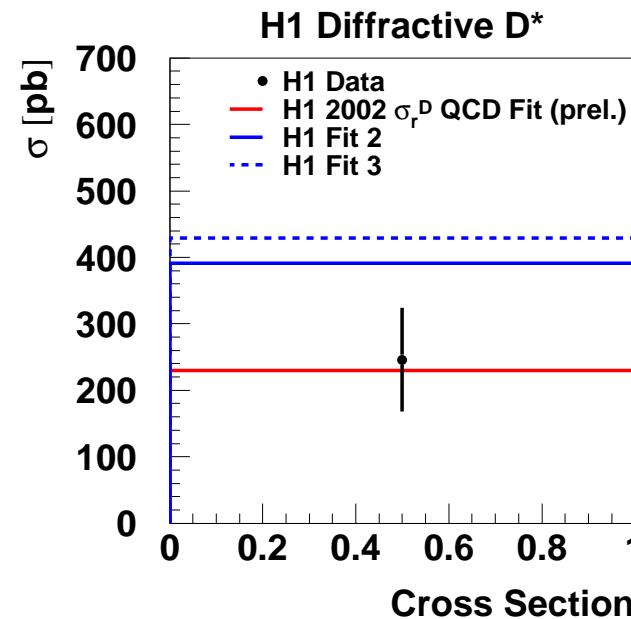
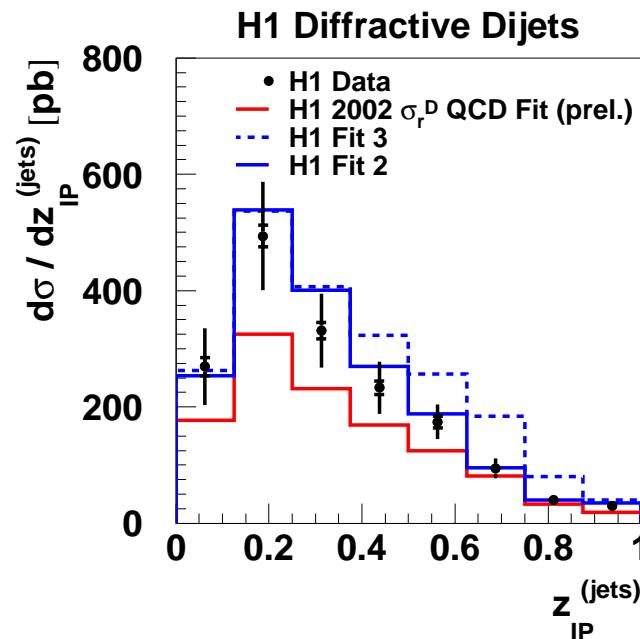
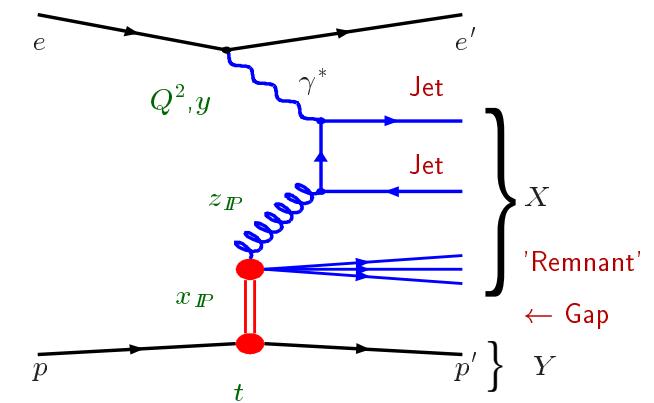
Comparison with H1 diffractive DIS final states

Use pdf's from LO fit to predict dijets / D^* cross sections
in diffractive DIS as measured by H1:

Comparison based on MC model (RAPGAP)

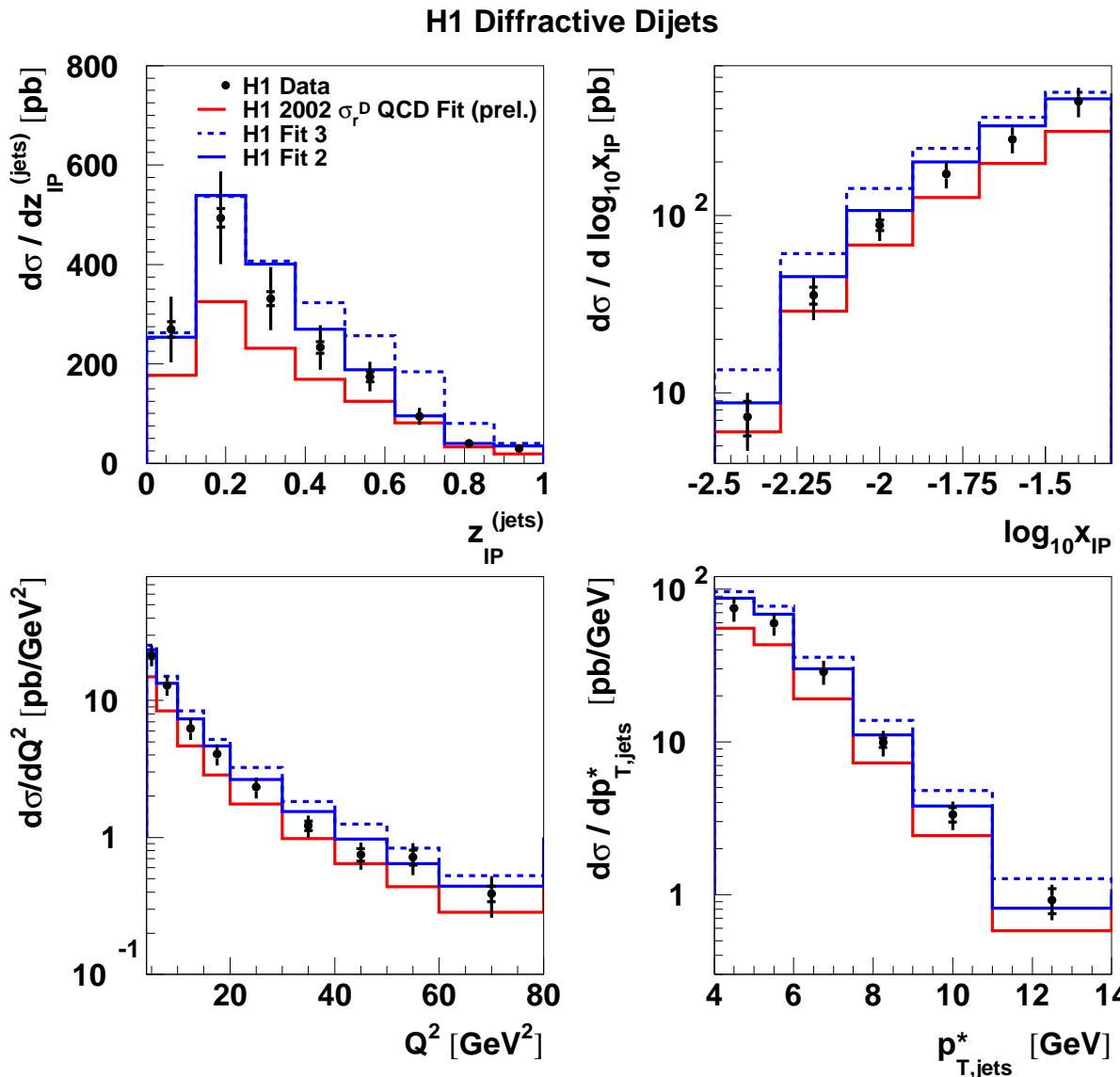
$$\mu^2 = Q^2 + p_T^2 + m^2$$

Differential distributions remain well described
Normalization: pdf/NLO/scale uncertainty



⇒ Consistent with QCD factorization !

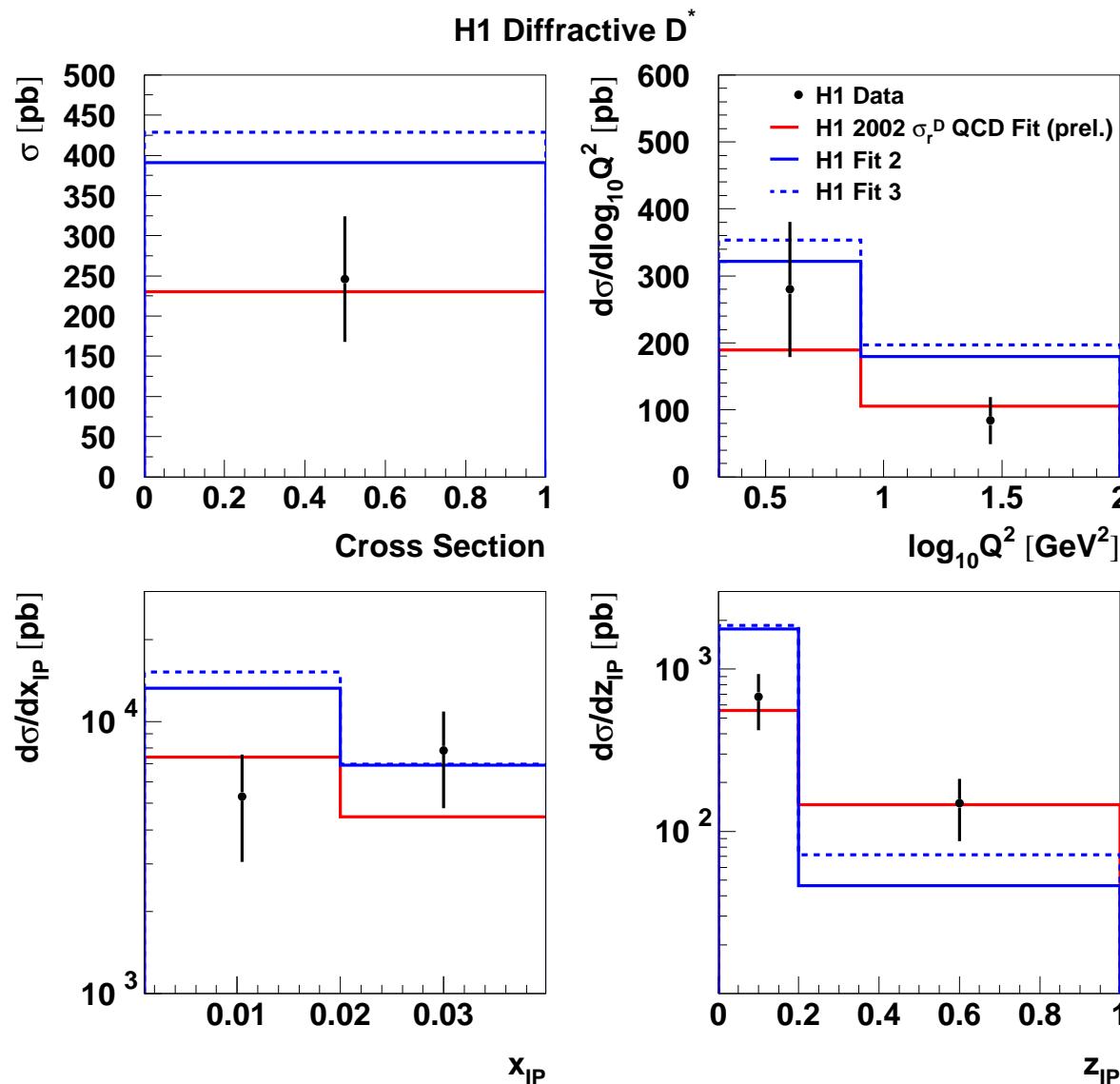
Diffractive Dijets



Good agreement, taking
NLO/pdf/scale
uncertainties
into account

Consistent with
QCD factorization

Diffractive Charm



Good agreement
within errors

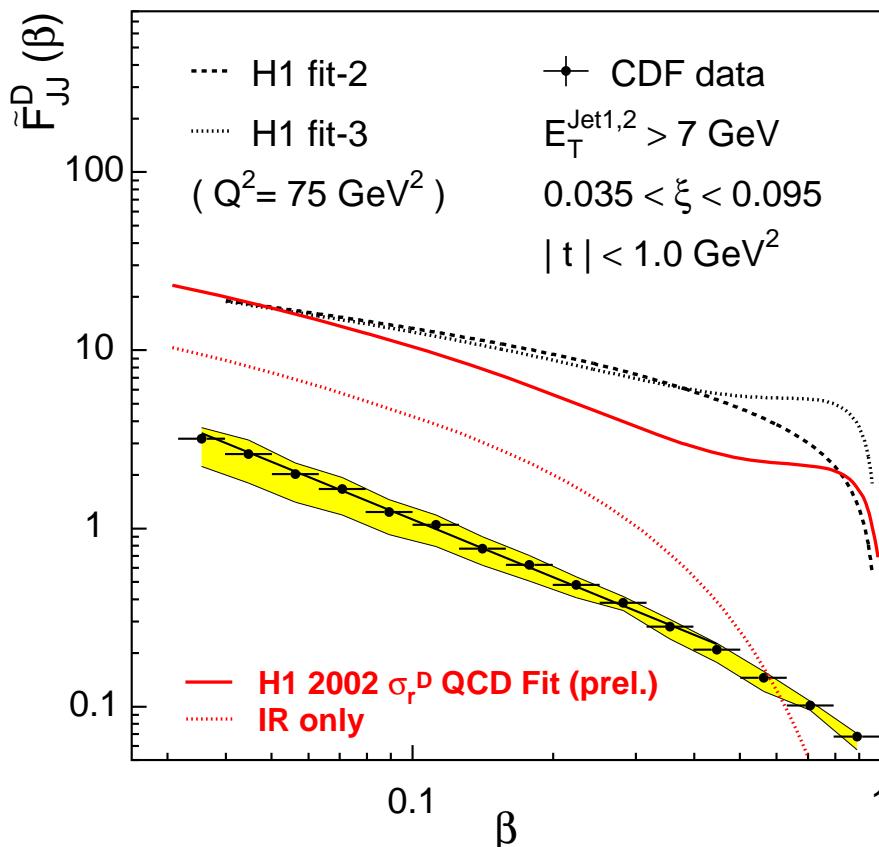
Consistent with
QCD factorization

Comparison with CDF diffractive Dijet cross sections

Dijet production with tagged leading anti-proton at TEVATRON:

Effective diffractive structure function \tilde{F}_{jj}^D :

$$\tilde{F}_{jj}^D(\beta) = \int dx_{IP} dt f(x_{IP}, t) \beta [g(\beta, Q^2) + \frac{4}{9}\Sigma(\beta, Q^2)] \quad (Q^2 = 75 \text{ GeV}^2)$$



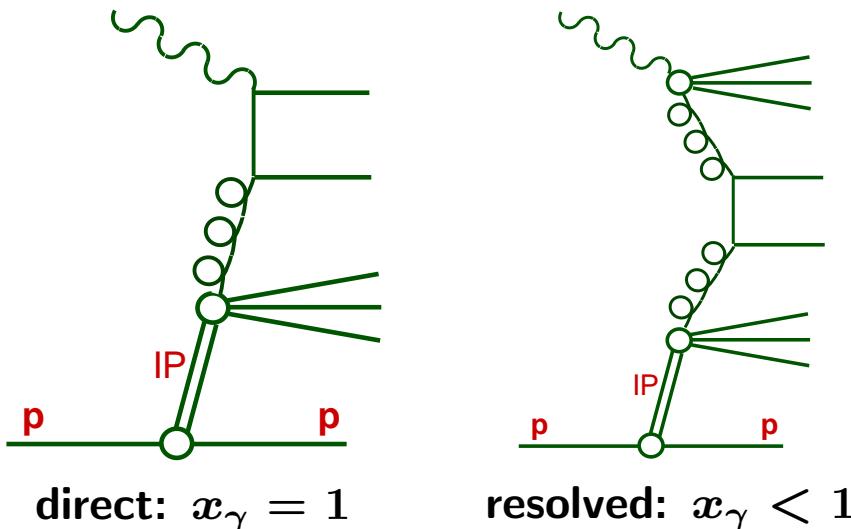
- New fit confirms serious breakdown of factorization (gap survival, absorptive corrections)
- β dependence similar (except highest β)
- **NOTE x_{IP} domain: 50% contribution from sub-leading exchange in this kinematic regime**

Control Experiment: Dijets in Photoproduction

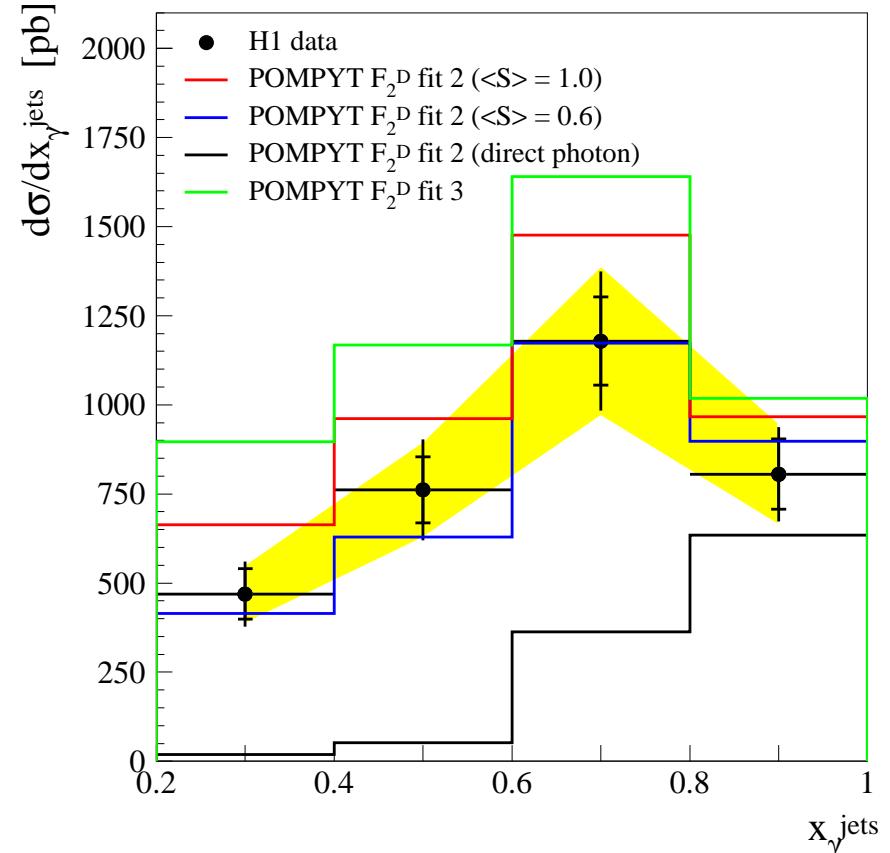
Hard diffractive photoproduction:

With and w/o remnants ...

x_γ : fraction of γ momentum entering hard scatter



Resolved γp resembles hadron-hadron!



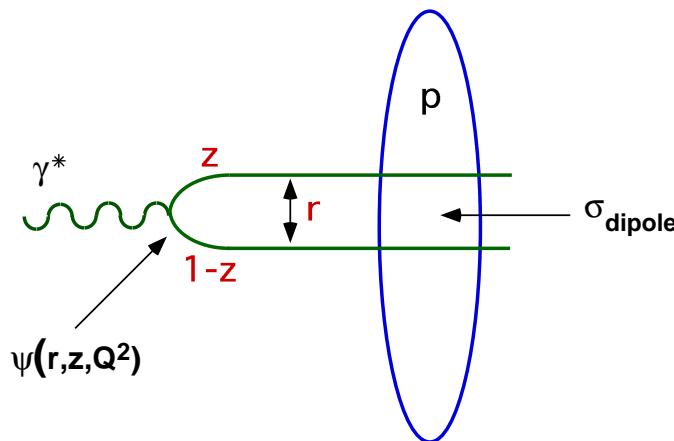
→ Description based on diffractive partons from F_2^D improved by suppressing resolved interactions by 'gap survival probability' of $\langle S \rangle = 0.6$

More data needed for firm conclusions ...

Colour Dipole and 2-gluon Exchange models

Consider low x in p rest frame:

$\gamma^* \rightarrow q\bar{q}$ well in advance of target
Interaction between photon and $q\bar{q}$ dipole



Diffractive if dipoles scatter elastically

Inclusive:

$$\sigma \sim |\Psi_{T,L}(r, z, Q^2)|^2 \otimes \sigma_{\text{dipole}}(x, r)$$

Diffractive:

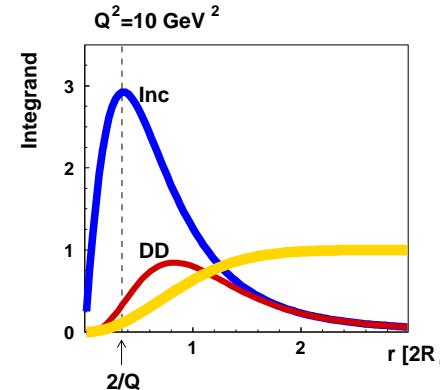
$$\sigma^D \sim |\Psi_{T,L}(r, z, Q^2)|^2 \otimes \sigma_{\text{dipole}}^2(x, r)$$

$\Psi_{T,L}(r, z, Q^2)$: Light cone wave functions

$$\begin{aligned}\gamma_T^* &\rightarrow q\bar{q} \\ \gamma_L^* &\rightarrow q\bar{q} \text{ (high } \beta) \\ \gamma_T^* &\rightarrow q\bar{q}g \text{ (low } \beta)\end{aligned}$$

Dipole cross section σ_{dipole} :

Model dependent, e.g. Golec-Biernat:



$\sim r^2$ as $r \rightarrow 0$
 $\rightarrow \text{const}$ as $r \rightarrow \infty$

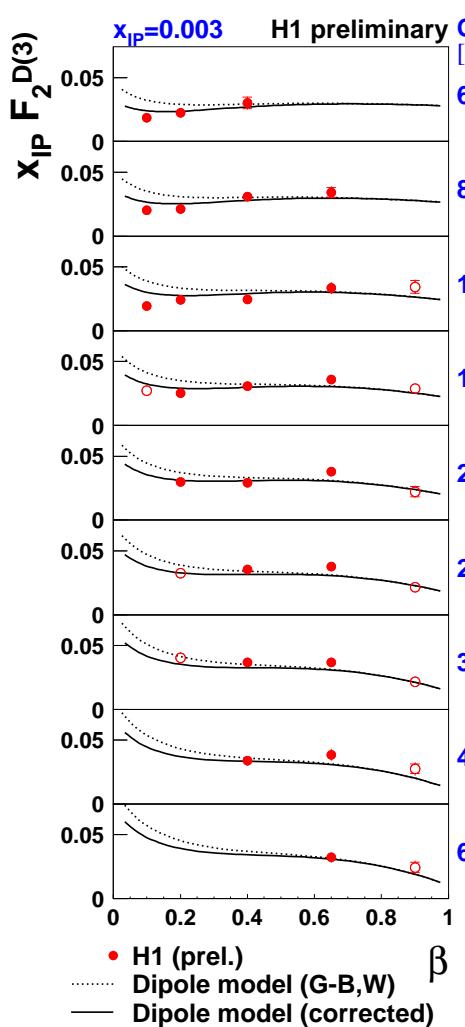
Sat. radius $R_0(x)$

Fix parameters by fit to inclusive $F_2(x, Q^2)$

Predict F_2^D at $t = 0$
Need t slope as input

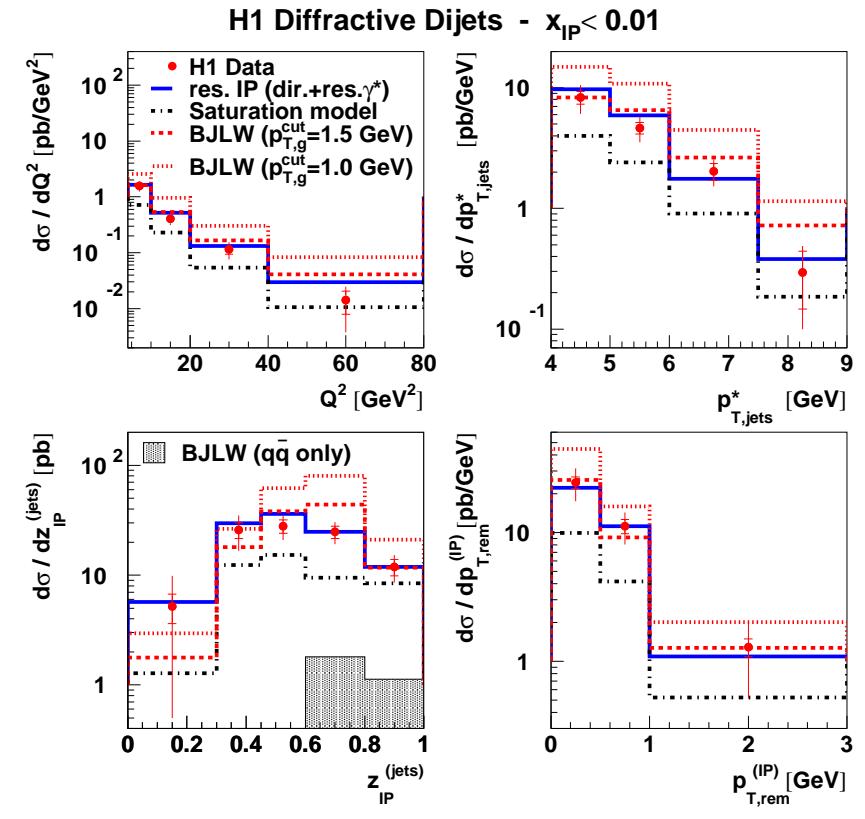
Dipole models: Comparison with Data

Inclusive: F_2^D



Agreement reasonable, except low Q^2, β

Diffr. Dijets



Saturation model underestimates normalization by factor 2

BJLW: Bartels et al. pQCD calculation for hard jets can describe data

Summary and Conclusions

Disclaimer:

- Impossible to cover everything in one talk. Here, the focus was on new results on inclusive hard diffraction (F_2^D) and QCD interpretation
- Many other interesting results (Vector mesons, DVCS, ...) could not be shown

Conclusions:

- Data on diffraction at HERA has reached **high precision !**
- Virtual photon in diffractive DIS enables to study quark/gluon (QCD) structure of diffraction
- New DGLAP NLO QCD fit to determine diffractive parton distributions (as for F_2) (justified by semi-inclusive QCD factorization proof)
- Large diffractive gluon distribution extending to large β
- Diffractive final states (Jets, charm):
QCD factorization works in ep , failure in pp confirmed!
- Real photoproduction: Bridge HERA - TEVATRON to understand factorization breaking ep vs pp ?